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On L^2 -boundedness of Fourier integral operators

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Abstract

Let $T_{a,\varphi}$ be a Fourier integral operator with symbol a and phase φ . In this paper, under the conditions $a(x, \xi) \in L^\infty S_\rho^{n(\rho-1)/2}(\omega)$ and $\varphi \in L^\infty \Phi^2$, the authors show that $T_{a,\varphi}$ is bounded from $L^2(\mathbb{R}^n)$ to $L^2(\mathbb{R}^n)$.

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1 Introduction and main results

Fourier integral operator on \mathbb{R}^n has been studied extensively and is related to many areas in analysis and PDEs. In [1], Sogge considered the Cauchy problem of the hyperbolic equations via the L^p -estimates theory of the Fourier integral operators (also see, for the local smoothing estimates of wave equations, e.g., [2, 3] and the references therein for some recent developments). For the Fourier integral operators with smooth amplitude, the L^2 -regularity theory is comparably more progress. In [4] and [5], Eskin and Hörmander found the local and global L^2 -regularity theory for Fourier integral operators, respectively. There are also some results for the L^p boundedness of Fourier integral operators with classical symbol and phase (see Littman [6], Miyachi [7], Peral [8], and Beals [9]).

Let \hat{f} be the Fourier transform of f . A Fourier integral operator T is a linear operator of the form

$$T_{a,\varphi}f = \frac{1}{(2\pi)^n} \int_{\mathbb{R}^n} e^{i\varphi(x,\xi)} a(x,\xi) \hat{f}(\xi) d\xi \quad (1.1)$$

with symbol $a(x, \xi)$ and phase $\varphi(x, \xi)$, respectively. In particular, for $\varphi(x, \xi) = \langle x, \xi \rangle$, the operator T_a is a so-called pseudo-differential operator. In [10], Hörmander showed that T_a is bounded in $L^2(\mathbb{R}^n)$, when $a \in S_{\rho,\delta}^m$, $\delta < 1$ and $m \leq n(\rho - \delta)/2$. For $a \in S_{1,1}^0$, Ching [11] proved that T_a is not bounded in $L^2(\mathbb{R}^n)$. Moreover, for $a \in S_{\rho,1}^m$, Rodino [12] showed that T_a is bounded in $L^2(\mathbb{R}^n)$ if and only if $m < n(\rho - 1)/2$. However, the operator T_a is not always L^2 -bounded for $a \in S_{\rho,1}^{n(\rho-1)/2}$; see, for example, [10–12]. The necessary and sufficient conditions of L^2 -boundedness of T_a were obtained by Higuchi [13] as $m = n(\rho - 1)/2$. It is natural to ask if the corresponding results hold for the Fourier integral operators. Recently, Kenig, David, Salvador, and Wolfgang [14–16] have studied the Fourier integral operators

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with rough symbol and rough phases, both of which behave in the spatial variable x like an L^∞ -function. More precisely, the symbol belongs to the class $L^\infty S_\rho^m$ whose constituent element a obeys

$$\|\partial_\xi^\alpha a(\cdot, \xi)\|_{L^\infty(\mathbb{R}^n)} \leq C_\alpha \langle \xi \rangle^{m-|\alpha|}.$$

Under this condition, for $m = \min\{0, \frac{n}{2}(\rho - \delta)\}$, $0 \leq \rho \leq 1$, $0 \leq \delta < 1$, and $a \in S_{\rho, \delta}^m$, Wolfgang [14] proved the global continuity on L^p -space with $p \in [1, \infty]$ of Fourier integral operators. A natural question is L^2 -boundedness of Fourier integral operators for $\delta = 1$ and $m = n(\rho - 1)/2$. In this paper, we answer the question and prove the results for the Fourier integral operators.

Our main result could be stated as follows.

Theorem 1.1 *Let $T_{a, \varphi}$ be a Fourier integral operator given by (1.1) with symbol $a(x, \xi) \in L^\infty S_\rho^{n(\rho-1)/2}(\omega)$ and phase function $\varphi \in L^\infty \Phi^2$ satisfying the Lipschitz rough non-degeneracy condition. Then, for $0 \leq \rho \leq 1$, there exists a positive constant C such that*

$$\|T_{a, \varphi} u\|_{L^2} \leq C \|u\|_{L^2}.$$

Here, the symbol class $L^\infty S_\rho^{n(\rho-1)/2}(\omega)$ is defined by Definition 2.2, the phase class $L^\infty \Phi^2$ is given by Definition 2.5, and the Lipschitz rough non-degeneracy condition is defined by Definition 2.6.

Remark 1.1 Here we remark that, for $a \in S_{\rho, 1}^{n(\rho-1)/2}$, Higuchi and Nagase [13] pointed out that the boundedness of the pseudo-differential operator T_a from $L^2(\mathbb{R}^n)$ to $L^2(\mathbb{R}^n)$ is not always true. As the main result in this paper, we give an answer for this problem for the Fourier integral operator $T_{a, \varphi}$. The main idea of our approach is treating the symbol class $L^\infty S_\rho^m(\omega)$, where $m = n(\rho - 1)/2$. In particular, our results of $L^2(\mathbb{R}^n)$ -boundedness for $T_{a, \varphi}$ are also the best as far as we know. We also remark that our methods are different from the previous methods; see, for example, [13].

Finally, we make some conventions on notation. Throughout this article, we denote by C a positive constant which is independent of the main parameters, but it may vary from line to line. We sometimes write $A \lesssim B$ as shorthand for $A \leq CB$. Let \mathbb{R}^n be an n -dimensional Euclidean space, $x = (x_1, \dots, x_n)$ be a point in \mathbb{R}^n , $\mathbb{R}_*^n = \mathbb{R}^n \setminus \{0\}$, $\mathbb{N} = \{1, 2, \dots\}$, $\mathbb{Z}_+ = \mathbb{N} \cup \{0\}$, and $\mathbb{Z}_+^n = (\mathbb{Z}_+)^n$. For any multi-index $\alpha = (\alpha_1, \dots, \alpha_n)$ and $\beta = (\beta_1, \dots, \beta_n) \in \mathbb{Z}_+^n$, we let

$$|\alpha| = \sum_{j=1}^n \alpha_j, \quad \alpha + \beta = (\alpha_1 + \beta_1, \dots, \alpha_n + \beta_n), \quad \partial_x^\alpha = \frac{\partial^\alpha}{\partial_{x_1}^{\alpha_1} \dots \partial_{x_n}^{\alpha_n}},$$

and $\nabla_\xi = (\partial_{\xi_1}, \dots, \partial_{\xi_n})$. Also, in the sequel we use the notation

$$|\xi| = \left(\sum_{j=1}^n \xi_j^2 \right)^{1/2} \quad \text{and} \quad \langle \xi \rangle = (1 + |\xi|^2)^{1/2}.$$

2 Definitions, notations, and preliminaries

The following definition is just [17].

Definition 2.1 Let $m \in \mathbb{R}$ and $0 \leq \delta, \rho \leq 1$. For any two multi-indices α and β , we assume that the function $a(x, \xi)$ satisfies the following condition:

$$|\partial_x^\beta \partial_\xi^\alpha a(x, \xi)| \leq C_{\alpha, \beta} \langle \xi \rangle^{m - \rho|\alpha| + \delta|\beta|}, \tag{2.1}$$

where $C_{\alpha\beta}$ is a positive constant only dependent on α and β . Let the smooth amplitude $S_{\rho, \delta}^m$ be the set of all smooth functions $a(x, \xi)$ satisfying condition as in (2.1). Then the pseudo-differential operator T_a with the symbol $a(x, \xi) \in S_{\rho, \delta}^m$ is given formally by

$$(T_a f)(x) = \int_{\mathbb{R}^n} e^{ix \cdot \xi} a(x, \xi) \hat{f}(\xi) d\xi.$$

The following definition for the class $L^\infty S_\rho^m(\omega)$ plays an important role in our setting.

Definition 2.2 Let m be a real number. A function $a(x, \xi)$, which is smooth in the frequency variable ξ and bounded measurable in the spatial variable x , belongs to the symbol class $L^\infty S_\rho^m(\omega)$ if, for all multi-indices α , it satisfies

$$\|\partial_\xi^\alpha a(x, \xi)\|_{L^\infty(\mathbb{R}^n)} \leq C_\alpha \langle \xi \rangle^{m - \rho|\alpha|} \omega(\langle \xi \rangle),$$

where $\omega(t)$ satisfies

$$\int_1^\infty \frac{\omega(t)^2}{t} dt < \infty, \tag{2.2}$$

and $\omega(t)$ is a nonnegative and decreasing function on $[1, \infty)$.

Remark 2.1 If $\omega(t)$ satisfies (2.1), then $\sum_{j=0}^\infty \omega^2(2^j) < \infty$.

David and Wong [14] gave the class Φ^k as follows.

Definition 2.3 ([14], Φ^k) A real-valued function $\varphi(x, \xi)$ belongs to the class Φ^k if $\varphi(x, \xi) \in C^\infty(\mathbb{R}^n \times \mathbb{R}_*^n)$ is positively homogeneous of degree 1 in the frequency variable ξ and satisfies the following condition: for any pair of multi-indices α and β , satisfying $|\alpha| + |\beta| \geq k$, there exists a positive constant $C_{\alpha, \beta}$ such that

$$\sup_{(x, \xi) \in \mathbb{R}^n \times \mathbb{R}_*^n} |\xi|^{-1 + \alpha} |\partial_\xi^\alpha \partial_x^\beta \varphi(x, \xi)| \leq C_{\alpha, \beta}.$$

In connection to the problem of local boundedness of Fourier integral operators, one considers phase functions $\varphi(x, \xi)$ that are positively homogeneous of degree 1 in the frequency variable ξ for which

$$\left| \det \frac{\partial^2 \varphi(x, \xi)}{\partial x_j \partial \xi_k} \right| \neq 0.$$

The latter is referred to as the non-degeneracy condition. However, for the purpose of proving global regularity results, we require a stronger condition than the non-degeneracy condition above.

Definition 2.4 ([14], The strong non-degeneracy condition) A real-valued function $\varphi(x, \xi) \in C^2(\mathbb{R}^n \times \mathbb{R}_*^n)$ satisfies strong non-degeneracy condition if there exists a positive constant c such that

$$\left| \det \frac{\partial^2 \varphi(x, \xi)}{\partial x_j \partial \xi_k} \right| \geq c$$

for all $(x, \xi) \in \mathbb{R}^n \times \mathbb{R}_*^n$.

Remark 2.2 The phases in class Φ^2 satisfying the strong non-degeneracy condition arise naturally in the study of the equations of hyperbolic type, namely

$$\varphi(x, \xi) = |\xi| + \langle x, \xi \rangle$$

belongs to the class Φ^2 and satisfies the strong non-degeneracy condition.

In [14], they introduced the nonsmooth version of the class Φ^k which will be used in our setting.

Definition 2.5 ([14], $L^\infty \Phi^k$) A real-valued function $\varphi(x, \xi)$ belongs to the phase class $L^\infty \Phi^k$ if it is positively homogeneous of degree 1 and smooth on \mathbb{R}_*^n in the frequency variable ξ , bounded measurable in the spatial variable x , and if for all multi-indices $|\alpha| \geq k$ it satisfies

$$\sup_{(x, \xi) \in \mathbb{R}^n \times \mathbb{R}_*^n} |\xi|^{-1+\alpha} |\partial_\xi^\alpha \varphi(x, \xi)| \leq C_\alpha.$$

Motivated by [14], we also need a Lipschitz rough non-degeneracy condition as follows.

Definition 2.6 (The Lipschitz rough non-degeneracy condition) A real-valued function satisfies Lipschitz rough non-degeneracy condition if it is C^∞ on \mathbb{R}_*^n in the frequency variable ξ , bounded measurable in the spatial variable x , and there exist positive constants C_1 and C_2 such that, for all $x, y \in \mathbb{R}^n$ and $\xi \in \mathbb{R}_*^n$,

$$\begin{aligned} |\partial_\xi \varphi(x, \xi) - \partial_\xi \varphi(y, \xi)| &\geq C_1 |x - y|, \\ |\partial_\xi^\alpha \varphi(x, \xi) - \partial_\xi^\alpha \varphi(y, \xi)| &\leq C_2 |x - y| \quad \text{for } |\alpha| \geq 2. \end{aligned}$$

3 Proof of the main result

In this section, we shall prove the main result, i.e., Theorem 1.1.

First we need a dyadic partition of unity. Let A be the annulus $A = \{\xi \in \mathbb{R}^n; \frac{1}{2} \leq |\xi| \leq 2\}$ and

$$\chi_0(\xi) + \sum_{j=1}^\infty \chi_j(\xi) = 1 \quad \text{for all } \xi \in \mathbb{R}^n,$$

where $\chi_0(\xi) \in C_0^\infty(B(0, 2))$ and $\chi_j(\xi) = \chi(2^{-j}\xi)$ when $j \geq 1$ with $\chi(\xi) \in C_0^\infty(A)$. Now we decompose the operator $T_{a,\varphi}$ as follows:

$$T_{a,\varphi} = T_{\chi_0}(D) + \sum_{j=1}^\infty T_{\chi_j}(D) = T_0(D) + \sum_{j=1}^\infty T_j(D). \tag{3.1}$$

The first term in (3.1) is bounded on $L^2(\mathbb{R}^n)$ from Theorem 1.1.8 in [14]. After a change of variables, we have

$$\begin{aligned} T_j(D) &= \frac{1}{(2\pi)^n} \int_{\mathbb{R}^n} e^{i\varphi(x,\xi)} \chi_j(\xi) a(x, \xi) \hat{u}(\xi) d\xi \\ &= \frac{2^{j\varrho n}}{(2\pi)^n} \int_{\mathbb{R}^n} e^{i \cdot 2^{j\varrho} \varphi(x,\xi)} \chi_j(2^{j\varrho} \xi) a(x, 2^{j\varrho} \xi) \hat{u}(2^{j\varrho} \xi) d\xi \\ &= \frac{2^{j\varrho n}}{(2\pi)^n} \int_{\mathbb{R}^n} e^{i \cdot 2^{j\varrho} \varphi(x,\xi)} \chi_j(2^{j\varrho} \xi) a(x, 2^{j\varrho} \xi) \int_{\mathbb{R}^n} e^{-i2^{j\varrho} \xi \cdot y} u(y) dy d\xi \\ &= \frac{2^{j\varrho n}}{(2\pi)^n} \int_{\mathbb{R}^n \times \mathbb{R}^n} e^{i \cdot 2^{j\varrho} (\varphi(x,\xi) - y \cdot \xi)} \chi_j(2^{j\varrho} \xi) a(x, 2^{j\varrho} \xi) u(y) d\xi dy. \end{aligned}$$

The kernel of the operator $T_j(D)$ is given by

$$T_j(x, y) = \frac{2^{j\varrho n}}{(2\pi)^n} \int_{\mathbb{R}^n} e^{i \cdot 2^{j\varrho} (\varphi(x,\xi) - y \cdot \xi)} \chi_j(2^{j\varrho} \xi) a(x, 2^{j\varrho} \xi) d\xi.$$

Let

$$a_j(x, \xi) = \chi(2^{j(1-\varrho)} \xi) a(x, 2^{j\varrho} \xi).$$

Then

$$A_j = \underset{\xi}{\text{Supp}} a_j \subset \{ \xi; 2^{-1} 2^{j(1-\varrho)} < |\xi| < 2 \cdot 2^{j(1-\varrho)} \}$$

and it satisfies

$$|\partial_\xi^\alpha a_j(x, \xi)| \leq C_\alpha \cdot 2^{jn(\rho-1)/2}. \tag{3.2}$$

We can confine ourselves to dealing with the high frequency component T_j of $T_{a,\varphi}$. Here we shall use a $S_j = T_j T_j^*$ argument, and therefore,

$$\begin{aligned} S_j u(x) &= \frac{1}{(2\pi)^n} \int_{\mathbb{R}^n \times \mathbb{R}^n} e^{i(\varphi(x,\xi) - \varphi(y,\xi))} \chi_j^2(\xi) a(x, \xi) \overline{a(y, \xi)} u(y) dy d\xi \\ &= \frac{2^{j\varrho n}}{(2\pi)^n} \int_{\mathbb{R}^n \times \mathbb{R}^n} e^{i2^{j\varrho} (\varphi(x,\xi) - \varphi(y,\xi))} \\ &\quad + \chi_j^2(2^{j\varrho} \xi) a(x, 2^{j\varrho} \xi) \overline{a(y, 2^{j\varrho} \xi)} u(y) d\xi dy. \end{aligned}$$

The kernel of the operator $S_j = T_j T_j^*$ reads

$$S_j(x, y) = \frac{2^{j\varrho n}}{(2\pi)^n} \int_{\mathbb{R}^n} e^{i2^{j\varrho} (\varphi(x,\xi) - \varphi(y,\xi))} \chi_j^2(2^{j\varrho} \xi) a(x, 2^{j\varrho} \xi) \overline{a(y, 2^{j\varrho} \xi)} d\xi.$$

Let $b_j(x, y, \xi) = \chi_j^2(2^{j\varrho} \xi) a(x, 2^{j\varrho} \xi) \overline{a(y, 2^{j\varrho} \xi)}$. Then

$$\text{Supp } b_j \subset \left\{ \xi : \frac{2^{j(1-\varrho)}}{2} < |\xi| < 2 \cdot 2^{j(1-\varrho)} \right\}.$$

We claim that

$$|\partial_\xi^\alpha b_j(x, y, \xi)| \leq C_\alpha 2^{jn(\rho-1)} \omega^2(2^j).$$

In fact,

$$\begin{aligned} |\partial_\xi^\alpha b_j(x, y, \xi)| &= |\partial_\xi^\alpha [\chi_j^2(2^{j\varrho} \xi) a(x, 2^{j\varrho} \xi) \overline{a(y, 2^{j\varrho} \xi)}]| \\ &= \sum_{\alpha_1 + \alpha_2 = \alpha} |\partial_\xi^{\alpha_1} [a(x, 2^{j\varrho} \xi) \overline{a(y, 2^{j\varrho} \xi)}]| |\partial_\xi^{\alpha_2} \chi^2(2^{-j(1-\varrho)} \xi)| \\ &\lesssim \sum_{\alpha_1 + \alpha_2 = \alpha} (2^{j\varrho})^{|\alpha_1|} |(\partial_\xi^{\alpha_1} (a \cdot \bar{a}))(x, 2^{j\varrho} \xi)| \omega^2(2^j) \\ &\quad \times 2^{-j(1-\varrho)|\alpha_2|} |(\partial_\xi^{\alpha_2} \chi)(2^{-j(1-\varrho)} \xi)| \\ &\lesssim \sum_{\alpha_1 + \alpha_2 = \alpha} 2^{j\varrho|\alpha_1|} |2^{j\varrho} \xi|^{n(\rho-1)-\varrho|\alpha_1|} 2^{-j(1-\varrho)|\alpha_2|} \omega^2(2^j) \\ &\lesssim \sum_{\alpha_1 + \alpha_2 = \alpha} 2^{j\varrho|\alpha_1|} 2^{jn(\rho-1)-\varrho|\alpha_1|} 2^{-j(1-\varrho)|\alpha_2|} \omega^2(2^j) \\ &= \sum_{\alpha_1 + \alpha_2 = \alpha} 2^{jn(\rho-1)-j(1-\varrho)|\alpha_2|} \omega^2(2^j) \\ &= 2^{jn(\rho-1)} \sum_{\alpha_2} 2^{-j(1-\varrho)|\alpha_2|} \omega^2(2^j) \\ &\lesssim 2^{jn(\rho-1)} \omega^2(2^j). \end{aligned} \tag{3.3}$$

Next we consider the following differential operators for $j \in \mathbb{N}$:

$$L_j(x, y, D) = \frac{\nabla_\xi \Phi \nabla_\xi}{i2^{j\varrho} |\nabla_\xi \Phi|^2}, \tag{3.4}$$

where $\Phi(x, y, \xi) = \varphi(x, \xi) - \varphi(y, \xi)$. So $L_j^N(x, y, D)e^{i2^{j\varrho} \Phi} = e^{i2^{j\varrho} \Phi}$ and

$$L_j^*(x, y, D) = -\nabla_\xi \frac{\nabla_\xi \Phi}{i2^{j\varrho} |\nabla_\xi \Phi|^2}. \tag{3.5}$$

From this and (3.4), it follows that

$$\begin{aligned} S_j(x, y) &= \frac{2^{j\varrho n}}{(2\pi)^n} \int_{\mathbb{R}^n} e^{i2^{j\varrho}(\varphi(x, \xi) - \varphi(y, \xi))} \chi_j^2(2^{j\varrho} \xi) a(x, 2^{j\varrho} \xi) \overline{a(y, 2^{j\varrho} \xi)} d\xi \\ &= \frac{2^{j\varrho n}}{(2\pi)^n} \int_{\mathbb{R}^n} (L_j^N e^{i2^{j\varrho} \Phi}) b_j(x, y, \xi) d\xi \\ &= \frac{2^{j\varrho n}}{(2\pi)^n} \int_{\mathbb{R}^n} e^{i2^{j\varrho} \Phi} (L_j^*)^N b_j d\xi. \end{aligned}$$

Moreover, by (3.5), we see that

$$\begin{aligned} \partial_{\xi_{\mu_1}} \left[\frac{\nabla_{\xi} \Phi}{|\nabla_{\xi} \Phi|^2} b_j \right] &= \left(\partial_{\xi_{\mu_1}} \left[\frac{\nabla_{\xi} \Phi b_j}{|\nabla_{\xi} \Phi|^2} \right] \right) \\ &= \left[\frac{\partial_{\xi_1} \Phi}{|\nabla_{\xi} \Phi|^2} \right] \partial_{\xi_{\mu_1}} b_j + \partial_{\xi_{\mu_1}} \left[\frac{\partial_{\xi_1} \Phi}{|\nabla_{\xi} \Phi|^2} \right] b_j, \dots, \left[\frac{\partial_{\xi_n} \Phi}{|\nabla_{\xi} \Phi|^2} \right] \partial_{\xi_{\mu_1}} b_j \\ &\quad + \partial_{\xi_{\mu_1}} \left[\frac{\partial_{\xi_n} \Phi}{|\nabla_{\xi} \Phi|^2} \right] b_j, \end{aligned}$$

which implies that

$$\begin{aligned} L_j^* b_j &= \nabla_{\xi} \left[\frac{\nabla_{\xi} \Phi}{-i2^j e |\nabla_{\xi} \Phi|^2} b_j \right] \\ &= \frac{1}{-i2^j e} \sum_{l=1}^n \left\{ \left[\frac{\partial_{\xi_l} \Phi}{|\nabla_{\xi} \Phi|^2} \right] \partial_{\xi_l} b_j + \partial_{\xi_l} \left[\frac{\partial_{\xi_l} \Phi}{|\nabla_{\xi} \Phi|^2} \right] b_j \right\}. \end{aligned}$$

Thus

$$\begin{aligned} (L_j^*)^N b_j &= \frac{1}{(-i2^j e)^N} \nabla_{\xi} \underbrace{\left\{ \frac{\nabla_{\xi} \Phi}{|\nabla_{\xi} \Phi|^2} \cdots \left[\nabla_{\xi} \left(\frac{\nabla_{\xi} \Phi}{|\nabla_{\xi} \Phi|^2} b_j \right) \right] \right\}}_N \\ &= \frac{1}{(-i2^j e)^N} \sum_{\alpha_1 + \dots + \alpha_N + \beta = N} \partial^{\alpha_1} \left(\frac{\Phi_{\mu_1}}{|\nabla \Phi|^2} \right) \cdots \partial^{\alpha_N} \left(\frac{\Phi_{\mu_N}}{|\nabla \Phi|^2} \right) \partial^{\beta} b_j, \end{aligned} \tag{3.6}$$

where $\Phi_{\mu_k} = \partial_{\xi_{\mu_k}} \Phi$. Because of the following equation

$$\partial_{\xi}^k \left(\frac{\Phi_{\mu}}{|\nabla \Phi|^2} \right) = \sum_{k_0 + k_1 + \dots + k_j = k} \frac{C_{k_0, \dots, k_j} \partial_{\xi}^{k_0} \Phi_{\mu} \partial_{\xi}^{k_1} |\nabla \Phi|^2 \cdots \partial_{\xi}^{k_j} |\nabla \Phi|^2}{|\nabla \Phi|^{2+2j}},$$

by Definition 2.6, we see that

$$|\partial_{\xi}^{k_0} \Phi_{\mu}| \lesssim |x - y| \quad \text{and} \quad |\partial_{\xi}^k |\nabla \Phi|^2| \lesssim |x - y|^2.$$

From this, together with (3.3) and (3.6), we further obtain

$$\begin{aligned} |(L_j^*)^N b_j| &\lesssim \frac{1}{2^j e^N} \cdot \frac{1}{|x - y|^N} 2^{jn(\rho-1)} \omega^2(2^j). \\ &= \frac{1}{[2^j e |x - y]|^N} 2^{jn(\rho-1)} \omega^2(2^j). \end{aligned}$$

Integration by parts yields

$$\begin{aligned} S_j(x, y) &= \frac{2^j e^n}{(2\pi)^n} \int_{\mathbb{R}^n} e^{i2^j e \Phi} (L_j^*)^N b_j d\xi \\ &\lesssim \frac{2^j e^n}{[2^j e |x - y]|^N} 2^{jn(\rho-1)} \omega^2(2^j) 2^{j(1-e)n} \\ &= \frac{2^j e^n}{[2^j e |x - y]|^N} \omega^2(2^j). \end{aligned}$$

Thus,

$$|S_j(x, y)| \sum_{l=0}^N (2^{le} |x - y|)^l \lesssim 2^{je n} \omega^2(2^j),$$

which implies that

$$|S_j(x, y)| \lesssim \frac{2^{je n}}{(1 + 2^{je} |x - y|)^N} \omega^2(2^j).$$

This further gives

$$\begin{aligned} \sup_x \int_{\mathbb{R}^n} |S_j(x, y)| dy &\lesssim 2^{je n} \omega^2(2^j) \int_{\mathbb{R}^n} \frac{1}{(1 + 2^{je} |x - y|)^N} dy \\ &\lesssim \omega^2(2^j) \int_{\mathbb{R}^n} \frac{1}{(1 + |z|)^N} dz \lesssim \omega^2(2^j). \end{aligned}$$

By Young’s inequality, we obtain

$$\|S_j u(x)\|_{L^2} \lesssim \omega^2(2^j) \|u(x)\|_{L^2}.$$

Therefore, we have

$$\begin{aligned} \|T_j^* u\|_{L^2}^2 &= \langle T_j^* u, T_j^* u \rangle = \langle u, T_j T_j^* u \rangle \\ &\leq \|u\|_{L^2} \|S_j u\|_{L^2} \\ &\leq C \omega^2(2^j) \|u\|_{L^2}^2. \end{aligned}$$

Namely,

$$\|T_j u\|_{L^2} \leq C \omega(2^j) \|u\|_{L^2}. \tag{3.7}$$

Next we need a Littlewood–Paley decomposition. Let $\psi_0 : \mathbb{R}^n \rightarrow \mathbb{R}$ be a smooth radial function which is equal to one on the unit ball centric at the origin and supported on its concentric double. Set $\psi(\xi) = \psi_0(\xi) - \psi_0(2\xi)$ and $\psi_k(\xi) = \psi(2^{-k}\xi)$. Then

$$\psi_0(\xi) + \sum_{k=1}^{\infty} \psi_k(\xi) = 1 \quad \text{for all } \xi \in \mathbb{R}^n,$$

and $\text{supp } \psi_k(\xi) \subset \{\xi : 2^{k-1} \leq |\xi| \leq 2^{k+1}\}$ for $k \geq 1$. And we further have

$$\hat{u}(\xi) = \sum_{k=0}^{\infty} \hat{u}(\xi) \psi_k(\xi) = \sum_{k=0}^{\infty} \hat{u}_k(\xi).$$

Then

$$T_{a,\varphi} u = \sum_{j=1}^{\infty} T_j u = \sum_{k=0}^{\infty} \sum_{j=1}^{\infty} \int_{\mathbb{R}^n} e^{i\varphi(x,\xi)} a_j(x, \xi) \hat{u}_k(\xi) d\xi.$$

For simplicity of notation, we write

$$\sum_{k=0}^{\infty} \sum_{j=1}^{\infty} \int_{\mathbb{R}^n} e^{i\varphi(x,\xi)} a_j(x, \xi) \hat{u}_k(\xi) d\xi = \sum_{j=1}^{\infty} T_j u_j,$$

where

$$u_j = \sum_{k=0}^{\infty} a_j(x, \xi) \hat{u}_k(\xi).$$

From this, (3.7), Cauchy–Schwartz’s inequality, and Remark 2.1, it follows that

$$\begin{aligned} \|T_{a,\varphi} u\|_{L^2} &= \left\| \sum_{j=0}^{\infty} T_j u_j \right\|_{L^2} \\ &\lesssim \sum_{j=0}^{\infty} \omega(2^j) \|u_j\|_{L^2} \\ &\lesssim \left(\sum_{j=0}^{\infty} \omega^2(2^j) \right)^{1/2} \left(\sum_{j=0}^{\infty} \|u_j\|_{L^2}^2 \right)^{1/2} \\ &\lesssim \|u\|_{L^2}. \end{aligned}$$

This finishes the proof of Theorem 1.1.

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There do not exist any competing interests regarding this article.

Authors’ contributions

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