# Positive periodic solution for $p$-Laplacian neutral damped Duffing equation with strong singularities of attractive and repulsive type 

## Shaowen Yao' and Xiaozhong Zhang ${ }^{1 *}$

Correspondence
xiaozhong99@126.com
School of Mathematics and Information Science, Henan Polytechnic University, Jiaozuo, China


#### Abstract

This paper is devoted to investigating the following p-Laplacian neutral damped Duffing equation with singularity: $$
\left(\phi_{p}(u(t)-c u(t-\tau))^{\prime}\right)^{\prime}+P u^{\prime}(t)+g(u(t))=e(t)
$$ where $g$ has a singularity at $u=0$. Applying the Manásevich-Mawhin theorem on a continuous case of topological degree, we obtain the existence of a positive periodic solution for this equation.


MSC: 34B16; 34B18; 34C25
Keywords: Positive periodic solution; Neutral operator; Attractive and repulsive singularities; Duffing equation

## 1 Introduction

In recent years, there has been a fairly great amount of work on periodic solutions for Duffing equations with a singularity (see [1-5, 7-11] and the references cited therein), in account of its applications in applied sciences. For example, the Brillouin electron beam focusing problem $[2,8]$ is to prove the existence of a positive $\pi$-periodic solutions of

$$
u^{\prime \prime}+a(1+\cos 2 t) u=\frac{1}{u},
$$

where $a$ is a positive constant.
In 2012, Cheng and Ren [1] discussed the existence and multiplicity of positive periodic solutions for the following Duffing equation:

$$
\begin{equation*}
u^{\prime \prime}(t)+g(u(t))=p(t), \tag{1.1}
\end{equation*}
$$

where the nonlinear term $g$ has a strong singularity of repulsive type at $u=0$ and satisfies super-linearity condition at $u=+\infty$. It is concluded that there exist infinitely many positive periodic solutions for Eq. (1.1) by applications of the generalized Poincaré-Birkhoff twist theorem. Afterwards, Wang and Ma [11] investigated Eq. (1.1) with $g$ has a strong
singularity of repulsive type at $u=0$ and a semi-linearity condition at $x=+\infty$, using the Poincaré-Birkhoff theorem, the authors proved the existence of infinitely many $2 \pi$ periodic solutions.

Proceeding from [ $1,2,8,11$ ], in this paper, we further consider the following $p$-Laplacian neutral Duffing equation with singularity:

$$
\begin{equation*}
\left(\phi_{p}(u(t)-c u(t-\tau))^{\prime}\right)^{\prime}+P u^{\prime}(t)+g(u(t))=e(t), \tag{1.2}
\end{equation*}
$$

where $p \geq 2, \varphi_{p}(u)=|u|^{p-2} u$ for $u \neq 0$ and $\varphi_{p}(0)=0, c$ is a constant and $|c| \neq 1, P, \tau$ are constants and $0 \leq \tau<T ; e \in L^{2}(\mathbb{R})$ is a $T$-periodic function; $g:(0,+\infty) \rightarrow \mathbb{R}$ is a $L^{2}$ Carathéodory function, the nonlinear term $g$ of (1.2) has a singularity at $u=0$, i.e., we have a strong singularity of repulsive type $\left(g_{1}\right)$ :

$$
\lim _{u \rightarrow 0^{+}} g(u)=-\infty, \quad \text { and } \quad \lim _{u \rightarrow 0^{+}} \int_{1}^{u} g(s) d s=+\infty ;
$$

or a strong singularity of attractive type $\left(g_{2}\right)$ :

$$
\lim _{u \rightarrow 0^{+}} g(u)=+\infty, \quad \text { and } \quad \lim _{u \rightarrow 0^{+}} \int_{1}^{u} g(s) d s=-\infty
$$

Applying topological degree theory [6], we obtain the following conclusions.

Theorem 1.1 Assume that conditions $|c|<1$ and $\left(g_{1}\right)$ hold. Suppose the following conditions hold:
$\left(\mathrm{H}_{1}\right)$ There exist two positive constants $d_{1}, d_{2}$ with $d_{1}<d_{2}$ such that $g(u)-e(t)<0$ for $(t, u) \in[0, T] \times\left(0, d_{1}\right)$ and $g(u)-e(t)>0$ for $(t, u) \in[0, T] \times\left(d_{2},+\infty\right)$.
Then Eq. (1.2) has at least one positive T-periodic solution.

Theorem 1.2 Assume that the conditions $|c|>1,\left(g_{1}\right)$ and $\left(\mathrm{H}_{1}\right)$ hold. Suppose the following condition holds:
$\left(\mathrm{H}_{2}\right)$ There exist two positive constants $\alpha, \beta$ such that

$$
g(u) \leq \alpha u^{p-1}+\beta, \quad \text { for all } u>0 .
$$

Then Eq. (1.2) has at least one positive T-periodic solution if $\alpha^{\frac{1}{p}}(1+2|c|)^{\frac{1}{p}} T<|c|-1$.

Remark 1.3 It is worth mentioning that the condition of the nonlinear term $g$ is relatively weak in the case that $|c|<1$, i.e., the nonlinear term $g$ may satisfy the sub-linearity, semilinearity or super-linearity conditions at $x=\infty$. The nonlinear term $g$ in the case that $|c|>1$ only satisfies the semi-linearity condition at $u=\infty$. Obviously, our result can be more general.

Remark 1.4 The condition $\left(g_{2}\right)$ contradicts the condition $\left(g_{1}\right)$. Therefore, the above methods of $[1,11]$ and condition of Theorems $1.1-1.2$ are no long applicable to prove the existence of a positive periodic solutions for Eq. (1.2) with strong singularity of attractive type. We need to give another method and conditions to get over this problem.

Theorem 1.5 Assume that conditions $|c|<1$ and $\left(g_{2}\right)$ hold. Furthermore, suppose the following condition holds:
$\left(\mathrm{H}_{3}\right)$ There exist two positive constants $d_{3}, d_{4}$ with $d_{3}<d_{4}$ such that $g(u)-e(t)>0$ for $(t, u) \in[0, T] \times\left(0, d_{3}\right)$ and $g(u)-e(t)<0$ for $(t, u) \in[0, T] \times\left(d_{4},+\infty\right)$.
Then Eq. (1.2) has at least one positive T-periodic solution.

Theorem 1.6 Assume that conditions $|c|>1,\left(g_{2}\right)$ and $\left(\mathrm{H}_{3}\right)$ hold. Suppose the following condition holds:
$\left(\mathrm{H}_{4}\right)$ There exist two positive constants $m, n$ such that

$$
-g(u) \leq m u^{p-1}+n, \quad \text { for all } u>0 .
$$

Then Eq. (1.2) has at least one positive T-periodic solution if $m^{\frac{1}{p}}(1+2|c|)^{\frac{1}{p}} T<|c|-1$.

## 2 Periodic solution for Eq. (1.2) with strong singularity of repulsive type

Firstly, let $A: C_{T} \rightarrow C_{T}$ be the operator on $C_{T}:=\{u \in C(\mathbb{R}, \mathbb{R}): u(t+T) \equiv u(t) \forall t \in \mathbb{R}\}$ given by

$$
(A u)(t):=u(t)-c u(t-\tau) \quad \forall u \in C_{T}, t \in \mathbb{R}
$$

Lemma 2.1 (see [12]) The operator $A$ has a continuous inverse $A^{-1}$ on $C_{T}$, satisfying
(1) $\left|\left[A^{-1} f\right](t)\right| \leq \frac{\|f\|}{|1-|c|,}, \forall f \in C_{T}$, where $\|f\|:=\max _{t \in \mathbb{R}}|f(t)|$.
(2) $\int_{0}^{T}\left|\left[A^{-1} f\right](t)\right| d t \leq \frac{1}{|1-|c||} \int_{0}^{T}|f(t)| d t, \forall f \in C_{T}$.

Secondly, we embed Eq. (1.2) into the following equation family with a parameter $\lambda \in$ $(0,1]$ :

$$
\begin{equation*}
\left(\phi_{p}(A u)^{\prime}(t)\right)^{\prime}+\lambda P u^{\prime}(t)+\lambda g(u(t))=\lambda e(t) . \tag{2.1}
\end{equation*}
$$

The following lemma is a consequence of Theorem 3.1 of [6].

Lemma 2.2 Assume that there exist positive constants $E_{1}, E_{2}, E_{3}$ and $E_{1}<E_{2}$ such that the following conditions hold:
(1) Each possible periodic solution $u$ to Eq. (2.1) such that $E_{1}<u(t)<E_{2}$, for all $t \in[0, T]$ and $\left\|u^{\prime}\right\|<E_{3}$.
(2) Each possible solution $C$ to the equation

$$
g(C)-\frac{1}{T} \int_{0}^{T} e(t) d t=0
$$

satisfies $C \in\left(E_{1}, E_{2}\right)$.
(3) We have

$$
\left(g\left(E_{1}\right)-\frac{1}{T} \int_{0}^{T} e(t) d t\right)\left(g\left(E_{2}\right)-\frac{1}{T} \int_{0}^{T} e(t) d t\right)<0
$$

Then Eq. (1.2) has at least one T-periodic solution.

### 2.1 Proof of Theorem 1.1

Proof of Theorem 1.1 Firstly, integrating both sides of Eq. (2.1) over [0, T], we get

$$
\begin{equation*}
\int_{0}^{T}[g(u(t))-e(t)] d t=0 \tag{2.2}
\end{equation*}
$$

In view of the mean value theorem of integrals, there exists a point $\xi \in(0, T)$ such that

$$
g(u(\xi))-e(\xi)=0 .
$$

From condition $\left(\mathrm{H}_{1}\right)$ and $u(t)$ being continuous, we have

$$
\begin{equation*}
d_{1} \leq u(\xi) \leq d_{2} \tag{2.3}
\end{equation*}
$$

Multiplying both sides of Eq. (2.1) by $(A u)^{\prime}(t)$ and integrating from 0 to $T$, we deduce

$$
\begin{align*}
& \int_{0}^{T}\left(\phi_{p}\left((A u)^{\prime}(t)\right)\right)^{\prime}(A u)^{\prime}(t) d t+\lambda P \int_{0}^{T} u^{\prime}(t)(A u)^{\prime}(t) d t+\lambda \int_{0}^{T} g(u(t))(A u)^{\prime}(t) d t \\
& \quad=\lambda \int_{0}^{T} e(t)(A u)^{\prime}(t) d t \tag{2.4}
\end{align*}
$$

Moreover,

$$
\begin{equation*}
\int_{0}^{T}\left(\phi_{p}(A u)^{\prime}(t)\right)^{\prime}(A u)^{\prime}(t) d t=\int_{0}^{T}(A u)^{\prime}(t) d \phi_{p}(A u)^{\prime}(t)=0 \tag{2.5}
\end{equation*}
$$

and

$$
\begin{align*}
\int_{0}^{T} g(u(t))(A u)^{\prime}(t) d t & =\int_{0}^{T} g(u(t))\left(A u^{\prime}\right)(t) d t \\
& =\int_{0}^{T} g(u(t))\left(u^{\prime}(t)-c u^{\prime}(t-\tau)\right) d t \\
& =\int_{0}^{T} g(u(t)) d u(t)-c \int_{0}^{T} g(u(t)) d u(t-\tau) \\
& =-c \int_{0}^{T} g(u(t)) d u(t) \\
& =0, \tag{2.6}
\end{align*}
$$

since $(A u)^{\prime}(t)=\left(A u^{\prime}\right)(t)$ and $d u(t)=\frac{d u(t-\tau)}{d(t-\tau)} d t=d u(t-\tau)$.
Substituting Eqs. (2.5) and (2.6) into (2.4), we obtain

$$
\begin{equation*}
P \int_{0}^{T} u^{\prime}(t)\left(A u^{\prime}\right)(t) d t=\int_{0}^{T} e(t)\left(A u^{\prime}\right)(t) d t \tag{2.7}
\end{equation*}
$$

From Eq. (2.7), we arrive at

$$
\left.\left|P \int_{0}^{T}\right| u^{\prime}(t)\right|^{2} d t\left|=\left|P c \int_{0}^{T} u^{\prime}(t) u^{\prime}(t-\tau) d t+\int_{0}^{T} e(t)\left(A u^{\prime}\right)(t) d t\right|\right.
$$

From the Hölder inequality and $\int_{0}^{T}\left|u^{\prime}(t-\tau)\right|^{2} d t=\int_{0}^{T}\left|u^{\prime}(t)\right|^{2} d t$, we see that

$$
\begin{aligned}
|P| \int_{0}^{T}\left|u^{\prime}(t)\right|^{2} d t \leq & |P||c| \int_{0}^{T}\left|u^{\prime}(t)\right|\left|u^{\prime}(t-\tau)\right| d t+(1+|c|) \int_{0}^{T}|e(t)|\left|u^{\prime}(t)\right| d t \\
\leq & |P||c|\left(\int_{0}^{T}\left|u^{\prime}(t)\right|^{2} d t\right)^{\frac{1}{2}}\left(\int_{0}^{T}\left|u^{\prime}(t-\tau)\right|^{2} d t\right)^{\frac{1}{2}} \\
& +(1+|c|)\left(\int_{0}^{T}|e(t)|^{2} d t\right)^{\frac{1}{2}}\left(\int_{0}^{T}\left|u^{\prime}(t)\right| d t\right)^{\frac{1}{2}} \\
= & |P||c| \int_{0}^{T}\left|u^{\prime}(t)\right|^{2} d t+(1+|c|)\|e\|_{2}\left(\int_{0}^{T}\left|u^{\prime}(t)\right|^{2} d t\right)^{\frac{1}{2}},
\end{aligned}
$$

where $\|e\|_{2}:=\left(\int_{0}^{T}|e(t)|^{2} d t\right)^{\frac{1}{2}}$. Obviously, $|P|-|P||c|>0$, since $|c|<1$. Hence, we deduce

$$
\begin{equation*}
\left(\int_{0}^{T}\left|u^{\prime}(t)\right|^{2} d t\right)^{\frac{1}{2}} \leq \frac{(1+|c|)\|e\|_{2}}{|P|-|P||c|}:=M_{1}^{\prime} . \tag{2.8}
\end{equation*}
$$

From Eqs. (2.3), (2.8) and the Hölder inequality, we see that

$$
\begin{align*}
u(t) & =u(\xi)+\int_{\xi}^{t} u^{\prime}(s) d s \leq d_{2}+\sqrt{T}\left(\int_{0}^{T}\left|u^{\prime}(t)\right|^{2} d t\right)^{\frac{1}{2}} \\
& \leq d_{2}+\sqrt{T} \frac{(1+|c|)\|e\|_{2}}{|P|-|P||c|}:=M_{1} . \tag{2.9}
\end{align*}
$$

On the other hand, from Eq. (2.2), it is clear that

$$
\begin{align*}
\int_{0}^{T}|g(u(t))| d t & =\int_{g(u(t)) \geq 0} g(u(t)) d t-\int_{g(u(t)) \leq 0} g(u(t)) d t \\
& =2 \int_{g(u(t)) \geq 0} g(u(t)) d t-\int_{0}^{T} e(t) d t . \tag{2.10}
\end{align*}
$$

Case (I). If $\bar{e}:=\frac{1}{T} \int_{0}^{T} e(t) d t \leq 0$, from Eq. (2.10), we have

$$
\int_{0}^{T}|g(u(t))| d t \leq 2 \int_{0}^{T}\left(g^{+}(u(t))-e(t)\right) d t
$$

where $g^{+}(u):=\max \{g(u), 0\}$. Since $g^{+}(u(t))-e(t) \geq 0$, from condition $\left(\mathrm{H}_{1}\right)$, we know $u(t) \geq$ $d_{2}$. Then we deduce

$$
\begin{align*}
\int_{0}^{T}|g(u(t))| d t & \leq 2 \int_{0}^{T} g^{+}(u(t)) d t+\int_{0}^{T}|e(t)| d t \\
& \leq 2 T\left\|g_{M_{1}}^{+}\right\|+T^{\frac{1}{2}}\|e\|_{2} \tag{2.11}
\end{align*}
$$

where $\left\|g_{M_{1}}^{+}\right\|:=\max _{d_{2} \leq u \leq M_{1}} g^{+}(u)$. As $(A u)(0)=(A u)(T)$, there exists a point $t_{1} \in[0, T]$ such that $(A u)^{\prime}\left(t_{1}\right)=0$, from Eqs. (2.2), (2.8) and (2.11), we have

$$
\begin{align*}
\left\|\phi_{p}\left((A u)^{\prime}\right)\right\| & =\max _{t \in\left[t_{1}, t_{1}+T\right]}\left\{\left|\int_{t_{1}}^{t}\left(\phi_{p}\left((A u)^{\prime}(s)\right)\right)^{\prime} d s\right|\right\} \\
& \leq|P| \int_{0}^{T}\left|u^{\prime}(t)\right| d t+\int_{0}^{T}|g(u(t))| d t+\int_{0}^{T}|e(t)| d t \\
& \leq|P| T^{\frac{1}{2}}\left(\int_{0}^{T}\left|u^{\prime}(t)\right|^{2} d t\right)^{\frac{1}{2}}+\int_{0}^{T}|g(u(t))| d t+\int_{0}^{T}|e(t)| d t \\
& \leq|P| T^{\frac{1}{2}} M_{1}^{\prime \frac{1}{2}}+2 T\left\|g_{M_{1}}^{+}\right\|+2 T^{\frac{1}{2}}\|e\|_{2}:=M_{2}^{\prime} . \tag{2.12}
\end{align*}
$$

We claim that there exists a positive constant $M_{2}^{* *}>M_{2}^{\prime}+1$ such that, for all $t \in \mathbb{R}$,

$$
\begin{equation*}
\left\|(A u)^{\prime}\right\| \leq M_{2}^{* *} \tag{2.13}
\end{equation*}
$$

In fact, if $x^{\prime}$ is not bounded, there exists a positive constant $M_{2}^{\prime \prime}$ such that $\left\|u^{\prime}\right\|>M_{2}^{\prime \prime}$ for some $u^{\prime} \in \mathbb{R}$. Therefor, it is clear that

$$
\left\|\phi_{p}(A u)^{\prime}\right\|=\left\|\phi_{p}\left(A u^{\prime}\right)\right\|=\left\|A u^{\prime}\right\|^{p-1}=(1+|c|)^{p-1}\left\|u^{\prime}\right\|^{p-1} \geq(1+|c|)^{p-1} M_{2}^{\prime \prime p-1}:=M_{2}^{*}
$$

Then we have a contradiction. So, Eq. (2.13) holds. By Lemma 2.1 and Eq. (2.13), we get

$$
\begin{align*}
\left\|u^{\prime}\right\| & =\left\|A^{-1} A u^{\prime}\right\|=\left\|A^{-1}(A u)^{\prime}\right\| \\
& \leq \frac{\left\|(A u)^{\prime}\right\|}{1-|c|} \\
& \leq \frac{M_{2}^{* *}}{1-|c|}:=M_{2} \tag{2.14}
\end{align*}
$$

since $|c|<1$.
Case (II). If $\bar{e}>0$, from Eq. (2.10), we obtain

$$
\int_{0}^{T}|g(u(t))| d t \leq 2 \int_{0}^{T} g^{+}(u(t)) d t
$$

Since $g^{+}(u(t)) \geq 0$, from condition $\left(\mathrm{H}_{1}\right)$, we know that there exists a positive constant $d_{2}^{*}$ such that $u(t) \geq d_{2}^{*}$. Therefore, we see that

$$
\begin{aligned}
\int_{0}^{T}|g(u(t))| d t & \leq 2 \int_{0}^{T} g^{+}(u(t)) d t \\
& \leq 2 T\left\|g_{M}^{+}\right\|
\end{aligned}
$$

where $\left\|g_{M}^{+}\right\|:=\max _{d_{2}^{*} \leq u \leq M_{1}} g^{+}(x)$. Similarly, we deduce $\left\|u^{\prime}\right\| \leq M_{2}$.

Multiplying both sides of Eq. (2.1) by $u^{\prime}(t)$ and integrating on the interval $[\xi, t]$, where $\xi \in[0, T]$ is defined in Eq. (2.3), we get

$$
\begin{align*}
\lambda \int_{u(\xi)}^{u(t)} g(v) d v= & \lambda \int_{\xi}^{t} g(u(s)) u^{\prime}(s) d s \\
= & -\int_{\xi}^{t}\left(\phi_{p}\left((A u)^{\prime}(s)\right)\right)^{\prime} u^{\prime}(s) d s-\lambda P \int_{\xi}^{t}\left|u^{\prime}(s)\right|^{2} d s \\
& +\lambda \int_{\xi}^{t} e(s) u^{\prime}(s) d s . \tag{2.15}
\end{align*}
$$

Furthermore, from Eqs. (2.12), (2.8) and (2.14), we get

$$
\begin{aligned}
\lambda\left|\int_{u(\xi)}^{u(t)} g(v) d v\right| \leq & \int_{\xi}^{t}\left|\left(\phi_{p}\left((A u)^{\prime}(s)\right)\right)^{\prime}\right|\left|u^{\prime}(s)\right| d s+\lambda P \int_{\xi}^{t}\left|u^{\prime}(s)\right|^{2} d s \\
& +\lambda \int_{\xi}^{t}|e(s)|\left|u^{\prime}(s)\right| d s \\
\leq & \lambda M_{2} M_{2}^{\prime}+\lambda|P|\left(M_{1}^{\prime}\right)^{2}+\lambda M_{2} T^{\frac{1}{2}}\|e\|_{2}:=\lambda M_{3}^{\prime} .
\end{aligned}
$$

From the repulsive condition $\left(g_{1}\right)$, we know that there exists a constant $M_{3}>0$ such that

$$
\begin{equation*}
u(t) \geq M_{3}, \quad \forall t \in[\xi, T] . \tag{2.16}
\end{equation*}
$$

Similarly, we can discuss $t \in[0, \xi]$.
From Eqs. (2.3), (2.9), (2.14) and (2.16), it is obvious that a periodic solution $u$ to Eq. (2.1) satisfies

$$
E_{1}<u(t)<E_{2}, \quad\left\|u^{\prime}\right\|<E_{3} .
$$

Then the condition (1) of Lemma 2.2 is satisfied. For a possible solution $C$ to the equation

$$
g(C)-\frac{1}{T} \int_{0}^{T} e(t) d t=0
$$

we have $C \in\left(E_{1}, E_{2}\right)$. Hence, the condition (2) of Lemma 2.2 holds. Finally, it is clear that the condition (3) of Lemma 2.2 is also satisfied. In fact, from condition $\left(\mathrm{H}_{1}\right)$, we can get

$$
g\left(E_{1}\right)-\frac{1}{T} \int_{0}^{T} e(t) d t<0
$$

and

$$
g\left(E_{2}\right)-\frac{1}{T} \int_{0}^{T} e(t) d t>0
$$

Using Lemma 2.2, it is concluded that Eq. (1.2) has at least one positive periodic solution.

### 2.2 Proof of Theorem 1.2

Proof of Theorem 1.2 The same strategy and notation are followed as in the proof of Theorem 1.1. Then we see that

$$
\begin{equation*}
u(t) \leq d_{2}+\int_{0}^{T}\left|u^{\prime}(t)\right| d t \tag{2.17}
\end{equation*}
$$

Multiplying both sides of Eq. (2.1) by $(A u)(t)$ and integrating on the interval [ $0, T$ ], it is clear that

$$
\begin{align*}
& \int_{0}^{T}\left(\phi_{p}\left((A u)^{\prime}(t)\right)\right)^{\prime}(A u)(t) d t+\lambda P \int_{0}^{T} u^{\prime}(t)(A u)(t) d t+\lambda \int_{0}^{T} g(u(t))(A u)(t) d t \\
& \quad=\lambda \int_{0}^{T} e(t)(A u)(t) d t \tag{2.18}
\end{align*}
$$

Substituting $\int_{0}^{T}\left(\phi_{p}(A u)^{\prime}(t)\right)^{\prime}(A u)(t) d t=-\int_{0}^{T}\left|(A u)^{\prime}(t)\right|^{p} d t$ and $P \int_{0}^{T} u^{\prime}(t) u(t) d t=0$ into Eq. (2.18), we have

$$
\begin{align*}
\int_{0}^{T}\left|(A u)^{\prime}(t)\right|^{p} d t= & -\lambda P c \int_{0}^{T} u^{\prime}(t) u(t-\tau) d t+\lambda \int_{0}^{T} g(u(t))(u(t)-c u(t-\tau)) d t \\
& -\lambda \int_{0}^{T} e(t)(u(t)-c u(t-\tau)) d t \tag{2.19}
\end{align*}
$$

Furthermore, we deduce

$$
\begin{equation*}
\int_{0}^{T} u^{\prime}(t) u(t-\tau) d t=\int_{0}^{T} u(t-\tau) d u(t)=\int_{0}^{T} u(t-\tau) d u(t-\tau)=0 \tag{2.20}
\end{equation*}
$$

From condition $\left(\mathrm{H}_{2}\right)$ and $u(t)>0$, we see that

$$
\begin{equation*}
\int_{0}^{T} g(u(t)) u(t) d t \leq \alpha \int_{0}^{T}(u(t))^{p} d t+\beta \int_{0}^{T} u(t) d t \tag{2.21}
\end{equation*}
$$

Substituting Eqs. (2.20) and (2.21) into (2.19), applying the Hölder inequality, we obtain

$$
\begin{align*}
\int_{0}^{T}\left|(A u)^{\prime}(t)\right|^{p} d t \leq & \alpha \int_{0}^{T}(u(t))^{p} d t+\beta \int_{0}^{T} u(t) d t+|c| \int_{0}^{T}|g(u(t))||u(t-\tau)| d t \\
& +(1+|c|)\|u\| \int_{0}^{T}|e(t)| d t \\
\leq & \alpha T\|u\|^{p}+\beta T\|u\|+|c|\|u\| \int_{0}^{T}|g(u(t))| d t \\
& +(1+|c|)\|u\| T^{\frac{1}{2}}\|e\|_{2} . \tag{2.22}
\end{align*}
$$

From Eq. (2.10) and condition $\left(\mathrm{H}_{2}\right)$, we get

$$
\begin{align*}
\int_{0}^{T}|g(u(t))| d t & =2 \int_{g(u(t)) \geq 0} g(u(t)) d t-\int_{0}^{T} e(t) d t \\
& \leq 2 \alpha\|u\|^{p-1} T+2 \beta T+\|e\|_{2} T^{\frac{1}{2}} \tag{2.23}
\end{align*}
$$

Substituting Eqs. (2.17) and (2.23) into (2.22), we see that

$$
\begin{align*}
& \int_{0}^{T}\left|(A u)^{\prime}(t)\right|^{p} d t \\
& \quad \leq \alpha(1+2|c|) T\|u\|^{p}+(1+2|c|)\left(\beta T+T^{\frac{1}{2}}\|e\|_{2}\right)\|u\| \\
& \quad \leq \alpha(1+2|c|) T\left(d_{2}+\int_{0}^{T}\left|u^{\prime}(t)\right| d t\right)^{p}+N_{1}\left(d_{2}+\int_{0}^{T}\left|u^{\prime}(t)\right| d t\right) \\
& \quad=\alpha(1+2|c|) T\left(\left(\int_{0}^{T}\left|u^{\prime}(t)\right| d t\right)^{p}+p d_{2}\left(\int_{0}^{T}\left|u^{\prime}(t)\right| d t\right)^{p-1}+\cdots+d_{2}^{p}\right) \\
& \quad+N_{1}\left(d_{2}+\int_{0}^{T}\left|u^{\prime}(t)\right| d t\right) \tag{2.24}
\end{align*}
$$

where $N_{1}:=(1+2|c|)\left(\beta T+T^{\frac{1}{2}}\|e\|_{2}\right)$. Applying Lemma 2.1, we have

$$
\begin{align*}
\int_{0}^{T}\left|u^{\prime}(t)\right| d t & =\int_{0}^{T}\left|\left(A^{-1} A u^{\prime}\right)(t)\right| d t \\
& \leq \frac{\int_{0}^{T}\left|(A u)^{\prime}(t)\right| d t}{|c|-1} \\
& \leq \frac{T^{\frac{1}{q}}\left(\int_{0}^{T}\left|(A u)^{\prime}(t)\right|^{p} d t\right)^{\frac{1}{p}}}{|c|-1} \tag{2.25}
\end{align*}
$$

since $\frac{1}{p}+\frac{1}{q}=1$ and $|c|>1$. We apply the inequality

$$
(x+y)^{k} \leq x^{k}+y^{k}, \quad \text { for } x, y>0,0<k<1 .
$$

Substituting Eqs. (2.24) into (2.25), we have

$$
\begin{aligned}
\int_{0}^{T}\left|u^{\prime}(t)\right| d t \leq & \frac{\alpha^{\frac{1}{p}}(1+2|c|)^{\frac{1}{p}} T \int_{0}^{T}\left|u^{\prime}(t)\right| d t}{|c|-1} \\
& +\frac{\alpha^{\frac{1}{p}}(1+2|c|)^{\frac{1}{p}}\left(p d_{2}\right)^{\frac{1}{p}} T\left(\int_{0}^{T}\left|u^{\prime}(t)\right| d t\right)^{\frac{p-1}{p}}}{|c|-1} \\
& +\cdots+\frac{T^{\frac{1}{q}}\left(N_{1}^{\frac{1}{p}}+\left(p d_{2}^{p-1}\right)^{\frac{1}{p}}\right)\left(\int_{0}^{T}\left|u^{\prime}(t)\right| d t\right)^{\frac{1}{p}}+T^{\frac{1}{p}}\left(d_{2}+\left(N_{1} d_{2}\right)^{\frac{1}{p}}\right)}{|c|-1}
\end{aligned}
$$

Since $\alpha^{\frac{1}{p}}(1+2|c|)^{\frac{1}{p}} T<|c|-1$, we know that there exists a positive constant $M_{1}^{\prime}$ such that

$$
\begin{equation*}
\int_{0}^{T}\left|u^{\prime}(t)\right| d t \leq M_{1}^{\prime} \tag{2.26}
\end{equation*}
$$

From Eqs. (2.17) and (2.26), we have

$$
\|u\| \leq d_{2}+\frac{1}{2} \int_{0}^{T}\left|u^{\prime}(t)\right| d t \leq d_{2}+\frac{M_{1}^{\prime}}{2}:=M_{1}
$$

The proof is left for the reader, being the same as that of Theorem 1.1.

## 3 Periodic solution for Eq. (1.2) with strong singularity of attractive type

### 3.1 Proof of Theorem 1.5

Proof of Theorem 1.5 We follow the same strategy and notation as in the proof of Theorem 1.1. We can get

$$
u(t) \leq M_{1} .
$$

From Eq. (2.10), we have

$$
\begin{align*}
\int_{0}^{T}|g(u(t))| d t & =\int_{g(u(t)) \geq 0} g(u(t)) d t-\int_{g(u(t)) \leq 0} g(u(t)) d t \\
& =-2 \int_{g(u(t)) \leq 0} g^{-}(u(t)) d t+\int_{0}^{T} e(t) d t, \tag{3.1}
\end{align*}
$$

where $g^{-}:=\min \{g(u), 0\}$.
Case (I). If $\bar{e} \geq 0$, from Eq. (3.1), we obtain

$$
\int_{0}^{T}|g(u(t))| d t \leq-2 \int_{0}^{T}\left(g^{-}(u(t))-e(t)\right) d t
$$

Since $g^{-}(u(t))-e(t) \leq 0$, from condition $\left(\mathrm{H}_{3}\right)$, we know that $u(t) \geq d_{4}$. Then we deduce

$$
\begin{align*}
\int_{0}^{T}|g(u(t))| d t & \leq-2 \int_{0}^{T} g^{-}(u(t)) d t+\int_{0}^{T}|e(t)| d t \\
& \leq 2 T\left\|g_{M_{1}}^{-}\right\|+T^{\frac{1}{2}}\|e\|_{2} \tag{3.2}
\end{align*}
$$

where $\left\|g_{M_{1}}^{-}\right\|:=\max _{d_{4} \leq u \leq M_{1}}\left(-g^{-}(u)\right)$. From Eqs. (2.12) and (3.2), we see that

$$
\begin{align*}
\left\|\phi_{p}\left((A u)^{\prime}\right)\right\| & \leq|p| \int_{0}^{T}\left|u^{\prime}(t)\right| d t+\int_{0}^{T}|g(u(t))| d t+\int_{0}^{T}|e(t)| d t \\
& \leq|p| T^{\frac{1}{2}} M_{1}^{\prime \frac{1}{2}}+2 T\left\|g_{M_{1}}^{-}\right\|+2 T^{\frac{1}{2}}\|e\|_{2}:=M_{2}^{\prime} \tag{3.3}
\end{align*}
$$

Case (II). If $\bar{e}<0$, from Eq. (3.1), we arrive at

$$
\int_{0}^{T}|g(u(t))| d t \leq-2 \int_{0}^{T} g^{-}(u(t)) d t
$$

Since $g^{-}(u(t)) \leq 0$, from condition $\left(\mathrm{H}_{3}\right)$, we know that there exists a positive constant $d_{4}^{*}$ such that $u(t) \geq d_{4}^{*}$. Therefore, we have

$$
\begin{aligned}
\int_{0}^{T}|g(u(t))| d t & \leq-2 \int_{0}^{T} g^{-}(u(t)) d t \\
& \leq 2 T\left\|g_{M}^{-}\right\|
\end{aligned}
$$

where $\left\|g_{M}^{-}\right\|:=\max _{d_{4}^{*} \leq u \leq M_{1}}\left(-g^{-}(u)\right)$. Similarly, we can get $\left|\phi_{p}\left((A u)^{\prime}(t)\right)\right| \leq M_{2}^{\prime}$.
The proof left for the reader, being the same as that of Theorem 1.1.

### 3.2 Proof of Theorem 1.6

Proof of Theorem 1.6 The same strategy and notation are followed as in the proof of Theorem 1.2. We only consider $\int_{0}^{T}|g(u(t))| d t$. From Eqs. (2.23), (3.1) and condition $\left(\mathrm{H}_{4}\right)$, we deduce

$$
\begin{align*}
\int_{0}^{T}|g(u(t))| d t & =-2 \int_{g(u) \leq 0} g^{-1}(u(t)) d t+\int_{0}^{T} e(t) d t \\
& \leq 2 m\|u\|^{p-1} T+2 n T+T^{\frac{1}{2}}\|e\|_{2} . \tag{3.4}
\end{align*}
$$

From Eqs. (2.24) and (3.1), we obtain

$$
\begin{aligned}
\int_{0}^{T}\left|(A u)^{\prime}(t)\right|^{p} d t \leq & m(1+2|c|) T\left(\left(\int_{0}^{T}\left|u^{\prime}(t)\right| d t\right)^{p}+\cdots+d_{4}^{p}\right) \\
& +N_{1}\left(d_{4}+\int_{0}^{T}\left|u^{\prime}(t)\right| d t\right) .
\end{aligned}
$$

The proof left for the reader, being the same as that of Theorem 1.1.

## 4 Examples

Example 4.1 Consider the following neutral Duffing equation with strong singularity of repulsive type:

$$
\begin{equation*}
\left(\phi_{p}\left(u(t)-\frac{1}{2} u(t-\tau)\right)^{\prime}\right)^{\prime}+P u^{\prime}(t)+\sum_{i=1}^{n} u^{2 i}(t)-\frac{1}{u^{\mu}}=e^{\cos t} \tag{4.1}
\end{equation*}
$$

where $\tau$ is a positive constant and $0<\tau<T, P$ and $\mu$ are constants and $\mu \geq 1, n$ is an integer.

Comparing Eqs. (4.1) to (1.2), we know that $g(u)=\sum_{i=1}^{n} u^{2 i}(t)-\frac{1}{u^{\mu}}, T=2 \pi, c=\frac{1}{2}<1$. Obviously, there exist constants $d_{1}=0.1$ and $d_{2}=1$ such that condition $\left(\mathrm{H}_{1}\right)$ holds. In fact, $\lim _{u \rightarrow 0^{+}} \int_{1}^{u} g(s) d s=\lim _{u \rightarrow 0^{+}} \int_{1}^{u}\left(\sum_{i=1}^{n} u^{2 i}(t)-\frac{1}{u^{\mu}}\right) d s=+\infty$, thus, the condition $\left(\mathrm{H}_{2}\right)$ holds. Therefore, applying Theorem 1.1, we know that Eq. (4.1) has at least one positive $2 \pi$-periodic solution.

Example 4.2 Consider the following neutral Duffing equation with strong singularity of attractive type:

$$
\begin{equation*}
\left(\phi_{p}(u(t)-40 u(t-\tau))^{\prime}\right)^{\prime}+8 u^{\prime}(t)-16 u^{3}+\frac{1}{u}=e^{\sin 2 t} \tag{4.2}
\end{equation*}
$$

where $p=4, \tau$ is a constant and $0 \leq \tau<T$.
It is clear that $T=\pi, P=8, c=40>1, g(u)=16 u^{3}-\frac{1}{u}, e(t)=e^{\sin 2 t}$. Obviously, it is easy to see that there exist constants $d_{3}=\frac{1}{4}$ and $d_{4}=\frac{3}{5}$ such that condition $\left(\mathrm{H}_{3}\right)$ holds. $\lim _{u \rightarrow 0^{+}} \int_{1}^{u} g(s) d s=\lim _{u \rightarrow 0^{+}} \int_{1}^{u}\left(-16 s^{3}+\frac{1}{s}\right) d s=-\infty$, thus, the condition $\left(\mathrm{g}_{2}\right)$ holds. Consider $-g(u) \leq 16 u^{3}+1$, where $m=16, n=1$. So, condition $\left(\mathrm{H}_{4}\right)$ is satisfied. Next, it is verified that

$$
\frac{m^{\frac{1}{p}}(1+2|c|)^{\frac{1}{p}} T}{|c|-1}=\frac{2 \times 3 \times \pi}{39}<1
$$

Therefore, using Theorem 1.6, it is concluded that Eq. (4.2) has at least one positive $\pi$ periodic solution.

## 5 Conclusion

In this paper, using the Manásevich-Mawhin theorem on the continuous case of topological degree, we discuss the existence of a positive periodic solution for the $p$-Laplacian neutral Duffing equation (1.2). The nonlinear term $g$ satisfies the strong singularity of attractive and repulsive type at $u=0$ and may obey the sub-linearity, semi-linearity or superlinearity conditions at $u=\infty$. First, we obtain the existence of a positive periodic solution for Eq. (1.2) with a strong singularity of repulsive type. Afterwards, we prove the existence of a positive periodic solution for Eq. (1.2) with a strong singularity of attractive type. Our results improve and extend the results in $[1,2,8,11]$.

## Acknowledgements

SWY and XZZ are grateful to anonymous referees for their constructive comments and suggestions, which have greatly improved this paper.

## Funding

This work was supported by National Natural Science Foundation of China (Nos. 71601072, 11501170)

## Abbreviations

Not applicable

Availability of data and materials
Not applicable.

## Ethics approval and consent to participate

SWY and XZZ contributed to each part of this study equally and declare that they have no competing interests.

## Competing interests

SWY and XZZ declare that they have no competing interests.

## Consent for publication

SWY and XZZ read and approved the final version of the manuscript

## Authors' contributions

SWY and XZZ contributed equally and significantly in writing this article. Both authors read and approved the final manuscript.

## Publisher's Note

Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.
Received: 2 December 2018 Accepted: 8 April 2019 Published online: 15 April 2019

## References

1. Cheng, Z., Ren, J.: Periodic and subharmonic solutions for Duffing equation with a singularity. Discrete Contin. Dyn. Syst., Ser. A 32(5), 1557-1574 (2013)
2. Cheng, Z., Yao, S.: New results for Brillouin electron beam focusing system. Bound. Value Probl. 2017, 70 (2017)
3. Chu, J., Wang, F.: Prevalence of stable periodic solutions for Duffing equations. J. Differ. Equ. 260(11), 7800-7820 (2016)
4. Ding, T., Zanolin, F.: Periodic solutions of Duffing's equations with superquadratic potential. J. Differ. Equ. 97(2), 328-378 (1992)
5. Fonda, A., Manásevich, R., Zanolin, F.: Subharmonic solutions for some second-order differential equations with singularities. Proc. R. Soc. Edinb. 120(3-4), 231-243 (1992)
6. Lu, S.: Periodic solutions to a second order p-Laplacian neutral functional differential system. Nonlinear Anal. TMA 69(11), 4215-4229 (2008)
7. Mishra, L., Agarwal, R., Sen, M.: Solvability and asymptotic behavior for some nonlinear quadratic integral equation involving Erdelyi-Kober fractional integrals on the unbounded interval. Prog. Fract. Differ. Appl. 2(3), 153-168 (2016)
8. Ren, J., Cheng, Z., Siegmund, S.: Positive periodic solution for Brillouin electron beam focusing system. Discrete Contin. Dyn. Syst., Ser. B 16(1), 385-392 (2011)
9. Torres, P.: Weak singularities may help periodic solutions to exist. J. Differ. Equ. 232(1), 277-284 (2007)
10. Wang, H.: Positive periodic solutions of singular systems with a parameter. J. Differ. Equ. 249(12), 2986-3002 (2010)
11. Wang, Z., Ma, T.: Existence and multiplicity of periodic solutions of semilinear resonant Duffing equations with singularities. Nonlinearity 25(2), 279-307 (2012)
12. Zhang, M.: Periodic solutions of linear and quasilinear neutral functional differential equations. J. Math. Anal. Appl. 189(3), 378-392 (1995)

Submit your manuscript to a SpringerOpen ${ }^{\circ}$ journal and benefit from:

- Convenient online submission
- Rigorous peer review
- Open access: articles freely available online
- High visibility within the field
- Retaining the copyright to your article

