Hindawi Publishing Corporation Journal of Inequalities and Applications Volume 2008, Article ID 961045, 14 pages doi:10.1155/2008/961045

Research Article

Global Existence and Extinction of Weak Solutions to a Class of Semiconductor Equations with Fast Diffusion Terms

Bin Wu

Department of Information and Computing Sciences, College of Mathematics and Physics, Nanjing University of Information Science and Technology, Nanjing 210044, China

Correspondence should be addressed to Bin Wu, wubing790831@126.com

Received 2 March 2008; Accepted 13 August 2008

Recommended by Y. Giga

We consider the transient drift-diffusion model with fast diffusion terms. This problem is not only degenerate but also singular. We first present existence result for general nonlinear diffusivities for the Dirichlet-Neumann mixed boundary value problem. Then, the extinction phenomenon of weak solutions for the homogeneous Dirichlet boundary problem is studied. Sufficient conditions on the extinction and decay estimates of solutions are obtained by using L^p -integral model estimate method.

Copyright © 2008 Bin Wu. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.

1. Introduction

The scaled semiconductor drift-diffusion model reads

$$-\nabla \cdot (\nabla \psi) = p - n + C(x), \tag{1.1}$$

$$n_t - \nabla \cdot J_n = r(n, p)(1 - np) + g, \quad J_n = (\nabla (n^m) - n\nabla \psi),$$
 (1.2)

$$p_t + \nabla \cdot J_p = r(n, p)(1 - np) + g, \quad -J_p = (\nabla(p^m) + p\nabla\psi), \tag{1.3}$$

with $x \in \Omega \subset R^N$, which denotes the bounded domain occupied by semiconductor crystal. Here, the unknowns ψ , n, and p denote the electrostatic potential, the electron density, and the hole density, respectively. The J_n represents the electron current, and J_p is the analogously defined physical quantity of the positively charged holes. Additionally, the given function C(x) denotes the doping profile (fixed charged background ions) characterizing the semiconductor under consideration, R(n,p) = r(n,p)(1-np) the net recombination-generation rate, and g the laser density. The constants m > 0 are the adiabatic or isothermal

(if m = 1) exponents. The regime m > 1 describes a slow diffusion process in the electron (hole) density, whereas 0 < m < 1 is related to fast diffusion.

We supplement these equations with physically motivated boundary conditions [1]:

$$(\psi, n, p) = (\psi_D, n_D, p_D), \qquad (x, t) \in \Sigma_D \equiv \Gamma_D \times (0, T), \tag{1.4}$$

$$\left(\frac{\partial \psi}{\partial \eta}, \frac{\partial n}{\partial \eta}, \frac{\partial p}{\partial \eta}\right) = (0, 0, 0), \qquad (x, t) \in \Sigma_N \equiv \Gamma_N \times (0, T), \tag{1.5}$$

$$(n,p) = (n_0, p_0), \qquad x \in \Omega, \ t = 0,$$
 (1.6)

where $\partial\Omega$ splits into two disjoint subsets Γ_N and Γ_D . Γ_N models the union of insulating boundary segments and Γ_D the union of ohmic contacts.

The standard drift-diffusion model corresponding to m=1 has been mathematically and numerically investigated in many papers (see [1–4]). Existence and uniqueness of weak solutions have been shown. Recently, the existence analysis of the bipolar drift-diffusion problem in the adiabatic case has been studied by many authors [5–9]. But, as far as we know, few works are concerned with the mixed boundary value problem of the drift-diffusion model with the fast diffusion terms.

The phenomenon of extinction is an important property of solutions for many evolutionary equations which have been studied extensively by many authors; see, for example, [10, 11]. Especially, there are some papers concerning the extinction for fast diffusive equations. For instance, in [12], the extinction and positivity for the evolutionary *p*-Laplacian equation without sources were studied by some authors, and Li and Wu [13] investigated the extinction for a fast diffusive filtration equation with nonlinear source terms. But, most of the work about extinction results is limited to single equation.

We prove first the global existence of weak solutions to problem (1.1)–(1.6) in the sense of the definition below. When $g \equiv 0$ and the boundary condition is homogeneous and of Dirichlet type, we show that the solutions n(x,t), p(x,t) vanish in finite time in the following two cases.

- (i) The initial data are appropriately small.
- (ii) The doping profile is sufficiently small and generation rate grows very slowly even if the initial data are very large.

Simultaneously, the decay estimates of solutions are obtained. For the proof of our result, we employ L^p -integral model estimate method and a crucial lemma on differential inequality. This technique has been successfully applied to the porous medium equation [14].

We make the following assumptions.

- (H1) $\Omega \subset \mathbb{R}^N$ (N = 1, 2, 3) is bounded and $\partial \Omega \in \mathbb{C}^{0,1}$, whose outward normal vector is η .
- (H2) $C(x), g(x) \in L^{\infty}(\Omega)$ and $g(x) \ge 0$ for a.e. $x \in \Omega$.
- (H3) r(n,p) is a locally Lipschitz continuous function defined for (n,p) and $0 \le r(n,p) \le \overline{r} < \infty$.
- (H4) $n_D^m, p_D^m, \psi_D \in H^1(\Omega) \cap L^{\infty}(\Omega)$, and $n_D, p_D \ge 0$ in Ω .
- (H5) $n_0, p_0 \in L^{\infty}(\Omega)$ and a.e. $n_0, p_0 \ge 0$ in Ω .

Remark 1.1. It is obvious that (H4) implies that $n_D, p_D \in H^1(\Omega) \cap L^{\infty}(\Omega)$. Let

$$Y = \{ \omega \in H^1(\Omega) \mid \omega|_{\Gamma_D} = 0 \}$$
(1.7)

and let Y^* denote the dual space of Y.

Definition 1.2. (ψ, n, p) is called the weak solution to problem (1.1)–(1.6) if $n \in n_D + L^2(0, T; Y)$, $p \in p_D + L^2(0, T; Y)$, $\psi \in \psi_D + L^2(0, T; Y)$, $n_t, p_t \in L^2(0, T; Y^*)$, $n|_{t=0} = n_0$, $p|_{t=0} = p_0$, and there hold

$$\int_{\Omega} \nabla \psi \cdot \nabla \phi \, dx = \int_{\Omega} (p - n + C) \phi \, dx, \quad \forall t \in (0, T), \ \forall \phi \in Y,$$

$$\langle n_{t}, v \rangle + \int_{0}^{T} \int_{\Omega} (\nabla (n^{m}) - n \nabla \psi) \cdot \nabla v \, dx \, dt$$

$$= \int_{0}^{T} \int_{\Omega} [r(n, p)(1 - np) + g] v \, dx \, dt, \quad \forall v \in L^{2}(0, T; Y),$$

$$\langle p_{t}, v \rangle + \int_{0}^{T} \int_{\Omega} (\nabla (p^{m}) + p \nabla \psi) \cdot \nabla v \, dx \, dt$$

$$= \int_{0}^{T} \int_{\Omega} [r(n, p)(1 - np) + g] v \, dx \, dt, \quad \forall v \in L^{2}(0, T; Y).$$
(1.8)

2. Existence

This section is devoted to the proof of global existence of weak solutions to problem (1.1)–(1.6). We will prove the following existence theorem.

Theorem 2.1. *Under hypotheses* (H1)–(H4), there exists at least one weak solution of problem (1.1)–(1.6).

The following lemma will be used in the proof of existence result.

Lemma 2.2 (see [9]). Let (H1) hold, $\alpha, \nu, F \in L^{\infty}(Q_T)$, and $F, \nu \geq 0$ for a.e. $(x,t) \in Q_T$, $\theta_0 \in L^{\infty}_+(\Omega)$, $\overline{\theta} \in H^1(\Omega) \cap L^{\infty}(\Omega)$, $\gamma \in (L^2(Q_T))^N$, $\gamma \cdot \eta = 0$, div $\gamma = i \in L^{\infty}(Q_T)$ weakly, and $\alpha \geq \varepsilon > 0$. Then, there exists a unique solution θ of the following problem:

$$\theta_{t} - \nabla \cdot (\alpha \nabla \theta - \gamma \theta) + \nu \theta = F, \quad (x, t) \in Q_{T},$$

$$\theta|_{\Sigma_{D}} = \overline{\theta}, \quad \frac{\partial \theta}{\partial \eta}\Big|_{\Sigma_{N}} = 0, \quad \theta(0) = \theta_{0}, \quad x \in \Omega,$$
(2.1)

such that $\theta \in L^2(0,T;Y)$, $\theta_t \in L^2(0,T;Y_0^*)$, and

$$0 \le \theta \le le^{\lambda t}$$
 a.e. $(x,t) \in Q_T$, (2.2)

where $l=(\|\overline{\theta}\|_{L^{\infty}(\Omega)}+\|\theta_0\|_{L^{\infty}(\Omega)}+1)$ and $\lambda\geq \|i\|_{L^{\infty}(Q_T)}+\|F\|_{L^{\infty}(Q_T)}.$

Our main difficulty in the proof is that problem (1.1)–(1.6) is degenerate at points where $n, p = +\infty$, and is singular at points where n, p = 0. This difficulty leads us to consider the following auxiliary regularized problem with the initial and boundary conditions (1.4)–(1.6):

$$-\nabla \cdot (\nabla \psi) = p_k - n_k + C(x), \tag{2.3}$$

$$n_t - \nabla \cdot \left(m \left(\frac{1}{n_k + \tau} \right)^{1-m} \nabla n - n \nabla \psi \right) = r(n_k, p_k) (1 - n p_k) + g, \tag{2.4}$$

$$p_t - \nabla \cdot \left(m \left(\frac{1}{p_k + \tau} \right)^{1-m} \nabla p + p \nabla \psi \right) = r(n_k, p_k) (1 - n_k p) + g, \tag{2.5}$$

where $0 < \tau < 1$ and $s_k = \min\{k, \max\{0, s\}\}.$

Now, we use the Schauder fixed point theorem [15] to prove the existence of weak solution to the above regularized problem.

Lemma 2.3. Under hypotheses (H1)–(H4), there exists at least one weak solution of problem (2.3)–(2.5), under conditions (1.4)–(1.6).

Proof. Define the set κ as

$$\kappa = \{ x \in L^2(Q_T) \mid x \ge 0 \text{ a.e. } x \in Q_{T_t} \|x\|_{L^2(Q_T)} \le R \}.$$
 (2.6)

It is obvious that κ is a closed convex set and weakly compact in $L^2(Q_T)$. Given $\tilde{n}, \tilde{p} \in \kappa$, we consider the following problems:

$$-\nabla \cdot (\nabla \psi) = \tilde{p}_k - \tilde{n}_k + C(x), \tag{2.7}$$

$$\psi|_{\Sigma_D} = \psi_D, \qquad \frac{\partial \psi}{\partial \eta}\Big|_{\Sigma_N} = 0,$$
 (2.8)

$$n_t - \nabla \cdot \left(m \left(\frac{1}{\tilde{n}_k + \tau} \right)^{1-m} \nabla n - n \nabla \psi \right) + r \left(\tilde{n}_k, \tilde{p}_k \right) \tilde{p}_k n = r \left(\tilde{n}_k, \tilde{p}_k \right) + g, \tag{2.9}$$

$$n|_{\Sigma_D} = n_D, \quad \frac{\partial n}{\partial \eta}\Big|_{\Sigma_N} = 0, \quad n(0) = n_0, \quad x \in \Omega,$$
 (2.10)

$$p_t - \nabla \cdot \left(m \left(\frac{1}{\widetilde{p}_k + \tau} \right)^{1-m} \nabla p + p \nabla \psi \right) + r(\widetilde{n}_k, \widetilde{p}_k) \widetilde{n}_k p = r(\widetilde{n}_k, \widetilde{p}_k) + g, \tag{2.11}$$

$$p|_{\Sigma_D} = p_D, \quad \frac{\partial p}{\partial \eta}\Big|_{\Sigma_{\Sigma_I}} = 0, \quad p(0) = p_0, \quad x \in \Omega.$$
 (2.12)

We deduce the existence of a unique weak solution of (2.7) and (2.8) with the regularity $\psi \in L^2(0,T;H^1(\Omega)) \cap L^\infty(Q_T)$ from standard theory. Set $\alpha = m(1/(\widetilde{n}_k + \tau))^{1-m}$, $\gamma = \nabla \psi$,

 $v = r(\tilde{n}_k, \tilde{p}_k)\tilde{p}_k$, and $F = r(\tilde{n}_k, \tilde{p}_k) + g$. Lemma 2.2 ensures the existence and uniqueness of weak solution of problem (2.9)-(2.10) with the regularity $n \in L^2(0, T; H^1(\Omega)) \cap L^{\infty}(Q_T)$, where L^{∞} estimate of n is dependent on k, τ and the known data, but independent of R by (2.2). Similar results hold for problem (2.11)-(2.12).

Thus, the map

$$S: \kappa^2 \longrightarrow (L^2(Q_T))^2, \qquad (\widetilde{n}, \widetilde{p}) \longmapsto (n, p)$$
 (2.13)

is well defined and compact. For the above given R, $S(\kappa^2) \hookrightarrow \kappa^2$ holds if we choose R sufficiently large. By using the standard method (for details, see [5] or [9]), we can obtain the continuity of map S. Thus, existence of a fixed point of S follows the weak solution of problem (2.3)–(2.5), under conditions (1.4)–(1.6).

Lemma 2.4. The weak solutions of problem (2.3)–(2.5), under conditions (1.4)–(1.6), satisfy the estimates

$$0 \le n(x,t), \quad p(x,t) \le k, \quad \text{a.e. } (x,t) \in Q_T,$$
 (2.14)

if one chooses k such that

$$k \ge M := \max \left\{ \sup_{\Omega} \{n_0, p_0\}, \sup_{\Sigma_D} \{n_D, p_D\} \right\} e^{(\overline{r} + \|g\|_{L^{\infty}(\Omega)} + \|C(x)\|_{L^{\infty}(\Omega)})T}.$$
 (2.15)

Proof. Set $N = ne^{-\beta t}$ and $P = pe^{-\beta t}$; here $\beta > 0$ is to be determined. Then, (N, P) satisfies

$$N_{t} + \beta N - \nabla \cdot \left(m \left(\frac{1}{(e^{\beta t} N)_{k} + \tau} \right)^{1 - m} \nabla N - N \nabla \psi \right) = e^{-\beta t} \left[r(n_{k}, p_{k}) (1 - n p_{k}) + g \right], \tag{2.16}$$

$$P_{t} + \beta P - \nabla \cdot \left(m \left(\frac{1}{\left(e^{\beta t}P\right)_{k} + \tau} \right)^{1-m} \nabla P + P \nabla \psi \right) = e^{-\beta t} \left[r(n_{k}, p_{k}) \left(1 - n_{k}p \right) + g \right], \quad (x, t) \in Q_{T},$$

$$(2.17)$$

$$(e^{\beta t}N, e^{\beta t}P) = (n_D, p_D), \quad (x, t) \in \Sigma_D,$$
 (2.18)

$$\left(\frac{\partial N}{\partial \eta}, \frac{\partial P}{\partial \eta}\right) = (0, 0), \quad (x, t) \in \Sigma_N, \tag{2.19}$$

$$(N, P) = (n_0, p_0), \quad x \in \Omega, \ t = 0.$$
 (2.20)

To obtain the upper bound, we compare N and P with $z(t) := k_0 e^{(\mu - \beta)t}$, where

$$k_0 \ge \max\left\{\sup_{\Omega} \{n_0, p_0\}, \sup_{\Sigma_D} \{n_D, p_D\}\right\},$$
 (2.21)

and μ such that $\mu > \beta$ is a large constant which can be determined by the following calculations. Then, we could take $(N-z)^+$ and $(P-z)^+$ as test functions in (2.16) and (2.17), respectively, and obtain

$$\frac{1}{2} \int_{\Omega} \left[(N-z)^{+2} (t) + (P-z)^{+2} (t) \right] + \beta \int_{0}^{t} \int_{\Omega} \left[(N-z)^{+2} + (P-z)^{+2} \right] \\
+ \int_{0}^{t} \int_{\Omega} m \left(\frac{1}{(e^{\beta t} N)_{k} + \tau} \right)^{1-m} \left| \nabla (N-z)^{+} \right|^{2} + \int_{0}^{t} \int_{\Omega} m \left(\frac{1}{(e^{\beta t} P)_{k} + \tau} \right)^{1-m} \left| \nabla (P-z)^{+} \right|^{2} \\
\leq -\mu \int_{0}^{t} \int_{\Omega} z \left[(N-z)^{+} + (P-z)^{+} \right] + \int_{0}^{t} \int_{\Omega} \left[N \cdot \nabla (N-z)^{+} - P \cdot \nabla (P-z)^{+} \right] \cdot \nabla \psi \\
+ \int_{0}^{t} \int_{\Omega} e^{-\beta t} \left[r(n_{k}, p_{k}) + g \right] \left[(N-z)^{+} + (P-z)^{+} \right] \\
\leq -\mu \int_{0}^{t} \int_{\Omega} z \left[(N-z)^{+} + (P-N)^{+} \right] + \frac{1}{2} \int_{0}^{t} \int_{\Omega} \left(p_{k} - n_{k} + C(x) \right) \left[(N-z)^{+2} - (P-z)^{+2} \right] \\
+ \int_{0}^{t} \int_{\Omega} \left(\overline{r} + g \right) z \left[(N-z)^{+} (P-N)^{+} \right] + \int_{0}^{t} \int_{\Omega} z \left(p_{k} - n_{k} + C(x) \right) \left[(N-z)^{+} + (P-z)^{+} \right] \\
\leq \left(-\mu + \overline{r} + \|g\|_{L^{\infty}(\Omega)} + \|C(x)\|_{L^{\infty}(\Omega)} \right) \int_{0}^{t} \int_{\Omega} z \left[(N-z)^{+} + (P-z)^{+} \right] \\
+ \frac{1}{2} \|C(x)\|_{L^{\infty}(\Omega)} \int_{0}^{t} \int_{\Omega} \left[(N-z)^{+2} + (P-z)^{+2} \right], \tag{2.22}$$

where in the last inequality we have used

$$(p_k - n_k) \left[(N - z)^+ - (P - z)^+ \right] = e^{-\beta t} (p_k - n_k) \left[(n - k_0 e^{\mu t})^+ - (p - k_0 e^{\mu t})^+ \right] \le 0.$$
 (2.23)

Choosing $\beta = (1/2) \| C(x) \|_{L^{\infty}(\Omega)}$ and $\mu \ge \overline{r} + \| g \|_{L^{\infty}(\Omega)} + \| C(x) \|_{L^{\infty}(\Omega)}$, we get

$$\frac{1}{2} \int_{\Omega} \left[(N-z)^{+^2} (t) + (P-z)^{+^2} (t) \right] + m \left(\frac{1}{k+\tau} \right)^{1-m} \int_{0}^{t} \int_{\Omega} \left[\left| \nabla (N-z)^{+} \right|^{2} + \left| \nabla (P-z)^{+} \right|^{2} \right] \le 0.$$
(2.24)

Hence, $(N-z)^+ + (P-z)^+ = 0$ a.e. $(x,t) \in Q_T$; this implies $n \le k_0 e^{\mu t}$ and $p \le k_0 e^{\mu t}$.

Thus, we obtain a weak solution of the following problem:

$$-\nabla \cdot (\nabla \psi) = p - n + C(x), \tag{2.25}$$

$$n_t - \nabla \cdot \left(m \left(\frac{1}{n+\tau} \right)^{1-m} \nabla n - n \nabla \psi \right) = r(n,p)(1-np) + g, \tag{2.26}$$

$$p_t - \nabla \cdot \left(m \left(\frac{1}{p+\tau} \right)^{1-m} \nabla p + p \nabla \psi \right) = r(n,p)(1-np) + g, \tag{2.27}$$

with the initial and boundary conditions (1.4)–(1.6).

In what follows, we give some estimates of the weak solutions to the above problem uniformly in τ , which are necessary in the proof of Theorem 2.1.

Lemma 2.5. The solutions of problem (2.25)–(2.27), under conditions (1.4)–(1.6), satisfy the estimates

$$m^{2} \int_{0}^{T} \int_{\Omega} \left[\left(\frac{1}{n+\tau} \right)^{2-2m} \left| \nabla n \right|^{2} + \left(\frac{1}{p+\tau} \right)^{2-2m} \left| \nabla p \right|^{2} \right] \le C, \tag{2.28}$$

where C is independent of τ .

Proof. First of all, a standard elliptic estimate gives

$$\|\nabla \psi\|_{L^2(Q_T)} \le C\left(1 + \|n\|_{L^2(Q_T)} + \|p\|_{L^2(Q_T)}\right) \le C. \tag{2.29}$$

Let

$$\varphi(n, n_D) = m \int_{n_D}^n \left(\frac{1}{s+\tau}\right)^{1-m} \mathrm{d}s. \tag{2.30}$$

Since $\varphi(n, n_D)|_{\Gamma_D} = 0$ and

$$\int_{0}^{T} \int_{\Omega} |\nabla \varphi(n, n_{D})|^{2} \leq 2m^{2} \int_{0}^{T} \int_{\Omega} \left[\left(\frac{1}{n+\tau} \right)^{2-2m} |\nabla n|^{2} + \left(\frac{1}{n_{D}+\tau} \right)^{2-2m} |\nabla n_{D}|^{2} \right] \\
\leq 2m^{2} \left(\frac{1}{\tau} \right)^{2-2m} \left(\|n\|_{L^{2}(0,T;H^{1}(\Omega))}^{2} + \|n_{D}\|_{L^{2}(0,T;H^{1}(\Omega))}^{2} \right) \leq C(\tau), \tag{2.31}$$

we have $\varphi(n, n_D) \in L^2(0, T; Y)$. Then,

$$m^{2} \int_{0}^{T} \int_{\Omega} \left(\frac{1}{n+\tau}\right)^{2-2m} |\nabla n|^{2} = m^{2} \int_{0}^{T} \int_{\Omega} \left(\frac{1}{n+\tau}\right)^{1-m} \nabla n \cdot \left(\frac{1}{n_{D}+\tau}\right)^{1-m} \nabla n_{D}$$

$$- \int_{0}^{T} \int_{\Omega} n_{t} \varphi(n, n_{D}) + \int_{0}^{T} \int_{\Omega} n \nabla \psi \cdot \nabla \varphi(n, n_{D})$$

$$+ \int_{0}^{T} \int_{\Omega} \left[r(n, p)(1-np) + g\right] \varphi(n, n_{D})$$

$$= I_{1} + I_{2} + I_{3} + I_{4}.$$

$$(2.32)$$

By Young's inequality and (2.29), we have

$$I_{1} \leq \frac{m^{2}}{4} \int_{0}^{\infty} \int_{\Omega} \left(\frac{1}{n+\tau} \right)^{2-2m} |\nabla n|^{2} + \int_{0}^{T} \int_{\Omega} |\nabla (n_{D}^{m})|^{2}, \tag{2.33}$$

$$I_{3} \leq \frac{1}{8} \int_{0}^{T} \int_{\Omega} |\nabla \varphi(n, n_{D})|^{2} + 4 \int_{0}^{T} \int_{\Omega} |n \nabla \psi|^{2}$$

$$\leq \frac{m^{2}}{4} \int_{0}^{T} \int_{\Omega} \left(\frac{1}{n+\tau}\right)^{2-2m} |\nabla n|^{2} + \frac{m^{2}}{4} \int_{0}^{T} \int_{\Omega} |\nabla (n_{D}^{m})|^{2} + 4M^{2} ||\nabla \psi||_{L^{2}(Q_{T})}^{2}$$

$$\leq \frac{m^{2}}{4} \int_{0}^{T} \int_{\Omega} \left(\frac{1}{n+\tau}\right)^{2-2m} |\nabla n|^{2} + C.$$
(2.34)

Furthermore,

$$I_{4} \leq m \left[\overline{r} (1 + M^{2}) + \|g\|_{L^{\infty}(\Omega)} \right] \int_{0}^{T} \int_{\Omega} \left[\int_{0}^{n} \left(\frac{1}{s + \tau} \right)^{1 - m} ds + \int_{0}^{n_{D}} \left(\frac{1}{s + \tau} \right)^{1 - m} ds \right]$$

$$\leq 2m \left[\overline{r} (1 + M^{2}) + \|g\|_{L^{\infty}(\Omega)} \right] |Q_{T}| \int_{0}^{M} \left(\frac{1}{s + \tau} \right)^{1 - m} ds \leq C.$$

$$(2.35)$$

For the term I_2 , we see that

$$|I_{2}| \leq \left| \int_{0}^{T} \int_{\Omega} n \varphi_{t}(n, n_{D}) \right| + \left| \int_{\Omega} n(T) \varphi(n, n_{D})(T) \right| + \left| \int_{\Omega} n(T) \varphi(n, n_{D})(0) \right|$$

$$\leq m \left| \int_{0}^{T} \int_{\Omega} n n_{t} \left(\frac{1}{n+\tau} \right)^{1-m} \right| + 4M |\Omega| \int_{0}^{M} \left(\frac{1}{s+\tau} \right)^{1-m} ds$$

$$\leq m \left| \int_{0}^{T} \int_{\Omega} \frac{\partial}{\partial t} \left(\int_{0}^{n} s \left(\frac{1}{s+\tau} \right)^{1-m} ds \right) \right| + C$$

$$\leq m \left| \int_{\Omega} \int_{n_{0}(x)}^{n(x,T)} s \left(\frac{1}{s+\tau} \right)^{1-m} ds \right| + C \leq C.$$

$$(2.36)$$

Inserting (2.33)–(2.36) into (2.32), we finally get

$$m^{2} \int_{0}^{T} \int_{\Omega} \left(\frac{1}{n+\tau} \right)^{2-2m} |\nabla n|^{2} \le C. \tag{2.37}$$

A similar estimate holds for *p*, and then we complete the proof.

Lemma 2.6. The weak solutions n and p of problem (2.25)–(2.27), under conditions (1.4)–(1.6), satisfy the following estimate:

$$\|\nabla n\|_{L^2(O_T)}, \|\nabla p\|_{L^2(O_T)} \le C.$$
 (2.38)

Proof. We will only prove that (2.38) holds for n. The rest can be derived similarly. Without loss of generality, we assume $0 < \tau < 1$. Then, we conclude that

$$\int_{0}^{T} \int_{\Omega} |\nabla n|^{2} = \int_{0}^{T} \int_{\Omega} (n+\tau)^{2-2m} \left[\left(\frac{1}{n+\tau} \right)^{1-m} |\nabla n| \right]^{2} \\
\leq (M+1)^{2-2m} \int_{0}^{T} \int_{\Omega} \left[\left(\frac{1}{n+\tau} \right)^{1-m} |\nabla n| \right]^{2} \leq C.$$
(2.39)

Following Lemmas 2.5 and 2.6, we can easily get the following lemma.

Lemma 2.7. The weak solutions n and p of problem (2.25)–(2.27), under conditions (1.4)–(1.6), satisfy the following estimate:

$$||n_t||_{L^2(0,T;Y^*)}, ||p_t||_{L^2(0,T;Y^*)} \le C.$$
 (2.40)

Proof of Theorem 2.1. Let $\{\psi_{\tau}, n_{\tau}, p_{\tau}\}$ be the sequence of solutions of problem (2.25)–(2.27), under conditions (1.4)–(1.6). By passing to a subsequence, if necessary, from Lemmas 2.4–2.7, we infer that

$$((n_{\varepsilon})_{t}, (p_{\varepsilon})_{t}) \longrightarrow (n_{t}, p_{t}) \text{ weakly in } L^{2}(0, T; Y^{*}),$$

$$(\psi_{\varepsilon}, n_{\varepsilon}, p_{\varepsilon}) \longrightarrow (\psi, n, p) \text{ strongly in } L^{2}(Q_{T}), \text{ a.e. in } Q_{T},$$

$$(\psi_{\varepsilon}, n_{\varepsilon}, p_{\varepsilon}) \longrightarrow (\psi, n, p) \text{ weakly in } L^{2}(0, T; H^{1}(\Omega)),$$

$$(\nabla (n+\tau)^{m}, \nabla (p+\tau)^{m}) \longrightarrow (\nabla (n^{m}), \nabla (p^{m})) \text{ weakly in } L^{2}(Q_{T}).$$

$$(2.41)$$

Now, we can conclude that (ψ, n, p) is a weak solution of problem (1.1)–(1.6) by standard method, and then complete the proof of Theorem 2.1.

3. Extinction

Besides the assumptions (H1)–(H5), we need for our extinction result the following structural condition on r(n,p).

(H6) $r(n,p) \le \rho_0(n+p)$ for some positive constant ρ_0 .

Because r(n,p) denotes the generation term in the recombination-generation rate, this condition means that the generation rate grows very slowly in (n,p) when ρ_0 is very small.

Additionally, we replace boundary conditions (1.4)-(1.5) with homogeneous Dirichlet conditions

$$n = p = \psi = 0, \qquad x \in \partial\Omega \times (0, T).$$
 (3.1)

It is obvious that the existence of weak solutions of problem (1.1)–(1.3), under conditions (3.1) and (1.6), can be obtained in easier way than the case of mixed boundary value problem.

Before stating our main result in this section, we list the following important conditions in Theorem 3.1:

$$||C(x)||_{L^{\infty}(\Omega)} + 2\rho_0 \le \frac{1}{2} C_0^{-2} |\Omega|^{-2m/(m+1) + (N-2)/N} \times \left(||n_0||_{L^{m+1}(\Omega)}^{m+1} + ||p_0||_{L^{m+1}(\Omega)}^{m+1} \right)^{(m-1)/(m+1)},$$
(3.2)

$$||n_0||_{L^{m+1}(\Omega)}^{m+1} + ||p_0||_{L^{m+1}(\Omega)}^{m+1} \le \left(\frac{C_0^{-2}|\Omega|^{-2m/(m+1)+(N-2)/N}}{||C(x)||_{L^{\infty}(\Omega)} + 2\rho_0}\right)^{(1+m)/(1-m)},\tag{3.3}$$

$$||C(x)||_{L^{\infty}(\Omega)} + 2\rho_0 \le \frac{2ms}{(m+s)^2} C_0^{-2} (||n_0||_{L^{s+1}(\Omega)}^{s+1} + ||p_0||_{L^{s+1}(\Omega)}^{s+1})^{(m-1)/(s+1)}, \tag{3.4}$$

$$||n_0||_{L^{s+1}(\Omega)}^{s+1} + ||p_0||_{L^{s+1}(\Omega)}^{s+1} \le \left[\frac{4msC_0^{-2}}{(m+s)^2(||C(x)||_{L^{\infty}(\Omega)} + 2\rho_0)} \right]^{(1+s)/(1-m)}, \tag{3.5}$$

where s = (N(1-m)/2) - 1 and C_0 is the embedding constant of the following inequality:

$$||u||_{L^{2N/(N-2)}(\Omega)} \le C_0 ||\nabla u||_{L^2(\Omega)}, \quad \forall u \in H^1_0(\Omega).$$
 (3.6)

Theorem 3.1. Let (H1)–(H6) be fulfilled, let $g(x) \equiv 0$, and let (ψ, n, p) be a weak solution of problem (1.1)–(1.3), under conditions (3.1) and (1.6). Then, one has the following.

(i) If (3.2) or (3.3) below is satisfied, then u(x,t) and p(x,t) vanish in the finite time for $(N-2)/(N+2) \le m < 1$, and

$$||n(\cdot,t)||_{L^{m+1}(\Omega)}^{m+1} + ||p(\cdot,t)||_{L^{m+1}(\Omega)}^{m+1} \le B_1 e^{-\alpha_1 t}, \quad t \in [0,T_{01}),$$

$$||n(\cdot,t)||_{L^{m+1}(\Omega)}^{m+1} + ||p(\cdot,t)||_{L^{m+1}(\Omega)}^{m+1}$$

$$\leq \left[\left(\left\| n(\cdot, T_{01}) \right\|_{L^{m+1}(\Omega)}^{m+1} + \left\| p(\cdot, T_{01}) \right\|_{L^{m+1}(\Omega)}^{m+1} \right)^{(1-m)/(1+m)} - \frac{C_1(1-m)}{1+m} t \right]^{(1+m)/(1-m)}, \qquad (3.7)$$

$$t \in [T_{01}, T_1),$$

$$||n(\cdot,t)||_{L^{m+1}(\Omega)}^{m+1} + ||p(\cdot,t)||_{L^{m+1}(\Omega)}^{m+1} \equiv 0, \quad t \in [T_1,\infty).$$

(ii) If (3.4) or (3.5) below is satisfied, then u(x,t) and p(x,t) vanish in the finite time for 0 < m < (N-2)/(N+2), and

$$||n(\cdot,t)||_{L^{s+1}(\Omega)}^{s+1} + ||p(\cdot,t)||_{L^{s+1}(\Omega)}^{s+1} \le B_2 e^{-\alpha_2 t}, \quad t \in [0,T_{02}),$$

$$||n(\cdot,t)||_{L^{s+1}(\Omega)}^{s+1} + ||p(\cdot,t)||_{L^{s+1}(\Omega)}^{s+1} ||n(\cdot,t)||_{L^{s+1}(\Omega)}^{s+1} + ||p(\cdot,t)||_{L^{s+1}(\Omega)}^{s+1}$$

$$\leq \left[\left(\left\| n(\cdot, T_{02}) \right\|_{L^{s+1}(\Omega)}^{s+1} + \left\| p(\cdot, T_{02}) \right\|_{L^{s+1}(\Omega)}^{s+1} \right)^{(1-m)/(1+s)} - \frac{C_2(1-m)}{1+s} t \right]^{(1+s)/(1-m)}, \qquad (3.8)$$

$$t \in [T_{02}, T_2),$$

$$||n(\cdot,t)||_{L^{s+1}(\Omega)}^{s+1} + ||p(\cdot,t)||_{L^{s+1}(\Omega)}^{s+1} \equiv 0, \quad t \in [T_2,\infty).$$

Here, C_1 , C_2 , T_1 , and T_2 will be given by (3.22), (3.32), (3.25), and (3.33), respectively.

To obtain the above result, we will use the following lemma which is of crucial importance to the proof.

Lemma 3.2. Let $0 < k < q \le 1$, and let $y(t) \ge 0$ be a solution of the differential inequality

$$\frac{dy}{dt} + \beta y^k \le \sigma y^q \quad (t \ge 0), \ y(0) = y_0 \ge 0, \tag{3.9}$$

where β and σ are positive constants. If

(i)
$$y_0 < \left(\frac{\beta}{\sigma}\right)^{1/(q-k)}$$
 or (ii) $\sigma \le \frac{\beta y_0^{k-q}}{2}$ (3.10)

holds, then there exist $\alpha > 0$ and B > 0 such that

$$0 \le y(t) \le Be^{-\alpha t}, \quad t \ge 0. \tag{3.11}$$

Proof. We will use upper and lower solutions' methods to complete the proof. The proof of case (i) is basically the same as method of case (ii). So, we are only devoted to the proof of case (i). In fact, we only need to choose α , B properly such that $Be^{-\alpha t}$ is an upper solution of (3.9). That is, α and B must satisfy

$$-\alpha B e^{-\alpha t} + \beta B^k e^{-\alpha kt} \ge \sigma B^q e^{-\alpha qt} \quad (t \ge 0), \ B \ge y_0. \tag{3.12}$$

Choose *B* such that $\ln B < \ln(\beta/\sigma)/(q-k)$ and α sufficiently small. Then,

$$\frac{\sigma}{\beta}B^{q-k} + \frac{\alpha}{\beta}B^{1-k} \le 1,\tag{3.13}$$

which implies that

$$\alpha(q-k)t \ge \ln\left(\frac{\sigma}{\beta}B^{q-k} + \frac{\alpha}{\beta}B^{1-k}\right), \quad \forall t \ge 0.$$
 (3.14)

Noting that $0 < q \le 1$, therefore we have

$$e^{-\alpha(k-q)t} \ge \frac{\sigma}{\beta} B^{q-k} + \frac{\alpha}{\beta} B^{1-k} e^{-(1-q)\alpha t}, \quad \forall t \ge 0, \tag{3.15}$$

which is equivalent to (3.12) if $y_0 \le B < (\beta/\sigma)^{1/(q-k)}$. The lemma is proved.

Proof of Theorem 3.1. We consider first the case $(N-2)/(N+2) \le m < 1$. Multiplying (1.2) and (1.3) by n^m and p^m , respectively, and employing the Poisson equation (1.1), we get

$$\frac{1}{m+1} \frac{\mathrm{d}}{\mathrm{d}t} \int_{\Omega} (n^{m+1} + p^{m+1}) \mathrm{d}x + \int_{\Omega} [|\nabla(n^m)|^2 + |\nabla(p^m)|^2] \mathrm{d}x$$

$$= \frac{m}{m+1} \int_{\Omega} (p-n+C(x)) (n^{m+1} + p^{m+1}) \mathrm{d}x + \int_{\Omega} r(n,p) (1-np) (n^m + p^m) \mathrm{d}x$$

$$= I_1 + I_2. \tag{3.16}$$

The embedding theorem gives

$$||n||_{L^{m+1}(\Omega)}^{m} \leq |\Omega|^{m/(m+1)-(N-2)/2N} ||n^{m}||_{L^{2N/(N-2)}(\Omega)}$$

$$\leq C_{0} |\Omega|^{m/(m+1)-(N-2)/2N} ||\nabla(n^{m})||_{L^{2}(\Omega)},$$
(3.17)

and similarly for p.

By Young's inequality, we have

$$I_2 \le \rho_0 \int_{\Omega} (n+p) (n^m + p^m) dx \le 2\rho_0 \int_{\Omega} (n^{m+1} + p^{m+1}) dx.$$
 (3.18)

By (3.16)–(3.18), taking into account $(p - n)(n^{m+1} - p^{m+1}) \le 0$, we obtain

$$\frac{\mathrm{d}}{\mathrm{d}t}f_1(t) + (m+1)C_0^{-2}|\Omega|^{-2m/(m+1)+(N-2)/N}f_1(t)^{2m/(m+1)} \le (m+1)(\|C(x)\|_{L^{\infty}(\Omega)} + 2\rho_0)f_1(t), \tag{3.19}$$

where

$$f_1(t) = \|n\|_{L^{m+1}(\Omega)}^{m+1} + \|p\|_{L^{m+1}(\Omega)}^{m+1}.$$
(3.20)

By Lemma 3.2, there exist α_1 and B_1 such that

$$0 \le f_1(t) \le B_1 e^{-\alpha_1 t}, \quad t \ge 0, \tag{3.21}$$

if (3.2) or (3.3) holds.

Furthermore, there exist T_{01} such that

$$(1+m)\left[C_{0}^{-2}|\Omega|^{-2m/(m+1)+(N-2)/N} - \left(\|C(x)\|_{L^{\infty}(\Omega)} + 2\rho_{0}\right)f_{1}(t)^{(1-m)/(1+m)}\right]$$

$$\geq (1+m)\left[C_{0}^{-2}|\Omega|^{-2m/(m+1)+(N-2)/N} - \left(\|C(x)\|_{L^{\infty}(\Omega)} + 2\rho_{0}\right)\left(B_{1}e^{-\alpha_{1}T_{01}}\right)^{(1-m)/(1+m)}\right]$$

$$:= C_{1} > 0$$

$$(3.22)$$

holds for $t \in [T_{01}, \infty)$. Thus, when $t \in [T_{01}, \infty)$, (3.19) turns to

$$\frac{\mathrm{d}}{\mathrm{d}t}f_1(t) + C_1 f_1(t)^{2m/(m+1)} \le 0, \quad t \ge T_{01}. \tag{3.23}$$

By a direct calculation, we deduce that

$$f_1(t) \le \left[f_1(T_{01})^{(1-m)/(1+m)} - \frac{C_1(1-m)}{1+m} t \right]^{(1+m)/(1-m)}, \quad t \in [T_{01}, T_1),$$

$$f_1(t) = 0, \quad t \in [T_1, \infty), \tag{3.24}$$

where

$$T_{1} = \frac{1+m}{C_{1}(1-m)} \left(\left\| n(\cdot, T_{01}) \right\|_{L^{m+1}(\Omega)}^{m+1} + \left\| p(\cdot, T_{01}) \right\|_{L^{m+1}(\Omega)}^{m+1} \right)^{(1-m)/(1+m)}. \tag{3.25}$$

For the case 0 < m < (N-2)/(N+2), using n^s and p^s as test functions in (1.2) and (1.3), respectively, we obtain

$$\frac{1}{s+1} \frac{\mathrm{d}}{\mathrm{d}t} \int_{\Omega} (n^{s+1} + p^{s+1}) \mathrm{d}x + \frac{4ms}{(m+s)^2} \int_{\Omega} \left[\left| \nabla (n^{(m+s)/2}) \right|^2 + \left| \nabla (p^{(m+s)/2}) \right|^2 \right] \mathrm{d}x$$

$$= \frac{s}{s+1} \int_{\Omega} (p-n+C(x)) (n^{s+1} + p^{s+1}) \mathrm{d}x + \int_{\Omega} r(n,p) (1-np) (n^s + p^s) \mathrm{d}x. \tag{3.26}$$

To simplify, we denote

$$f_2(t) = \|n\|_{L^{s+1}(\Omega)}^{s+1} + \|p\|_{L^{s+1}(\Omega)}^{s+1}. \tag{3.27}$$

Hence, we have

$$\frac{\mathrm{d}}{\mathrm{d}t}f_2(t) + \frac{4(s+1)ms}{(m+s)^2}C_0^{-2}f_2(t)^{(m+s)/(1+s)} \le (s+1)(\|C(x)\|_{L^{\infty}(\Omega)} + 2\rho_0)f_2(t),\tag{3.28}$$

where we have used

$$\|u\|_{L^{s+1}(\Omega)}^{(m+s)/2} = \left(\int_{\Omega} u^{(m+s)/2 \cdot 2N/(N-2)} dx\right)^{(N-2)/2N} \le C_0 \|\nabla(u^{(m+s)/2})\|_{L^2(\Omega)}, \quad \forall u \in H_0^1(\Omega).$$
(3.29)

By Lemma 3.2, there exist α_2 and B_2 such that

$$0 \le f_2(t) \le B_2 e^{-\alpha_2 t}, \quad t \ge 0, \tag{3.30}$$

if (3.4) or (3.5) holds.

By a simple analysis similar to that of the case $(N-2)/(N+2) \le m < 1$, we obtain

$$f_{2}(t) \leq \left[f_{2}(T_{02})^{(1-m)/(1+s)} - \frac{C_{2}(1-m)}{1+s} t \right]^{(1+s)/(1-m)}, \quad t \in [T_{02}, T_{2}),$$

$$f_{2}(t) \equiv 0, \quad t \in [T_{2}, \infty),$$

$$(3.31)$$

where

$$C_2 = (s+1) \left[\frac{4ms}{(m+s)^2} C_0^{-2} - (\|C(x)\|_{L^{\infty}(\Omega)} + 2\rho_0) (B_2 e^{-\alpha_2 T_{02}})^{(1-m)/(1+s)} \right] > 0, \tag{3.32}$$

$$T_{2} = \frac{1+s}{C_{2}(1-m)} \left(\|n(\cdot, T_{02})\|_{L^{s+1}(\Omega)}^{s+1} + \|p(\cdot, T_{02})\|_{L^{s+1}(\Omega)}^{s+1} \right)^{(1-m)/(1+s)}.$$
 (3.33)

References

- [1] P. A. Markowich, C. A. Ringhofer, and C. Schmeiser, Semiconductor Equations, Springer, Vienna, Austria, 1990.
- [2] H. Gajewski, "On existence, uniqueness and asymptotic behavior of solutions of the basic equations for carrier transport in semiconductors," *Zeitschrift für Angewandte Mathematik und Mechanik*, vol. 65, no. 2, pp. 101–108, 1985.

- [3] W. F. Fang and K. Ito, "On the time-dependent drift-diffusion model for semiconductors," *Journal of Differential Equations*, vol. 117, no. 2, pp. 245–280, 1995.
- [4] W. F. Fang and K. Ito, "Global solutions of the time-dependent drift-diffusion semiconductor equations," *Journal of Differential Equations*, vol. 123, no. 2, pp. 523–566, 1995.
- [5] A. Jüngel, "On the existence and uniqueness of transient solutions of a degenerate nonlinear drift-diffusion model for semiconductors," *Mathematical Models & Methods in Applied Sciences*, vol. 4, no. 5, pp. 677–703, 1994.
- [6] A. Jüngel, "Qualitative behavior of solutions of a degenerate nonlinear drift-diffusion model for semiconductors," *Mathematical Models & Methods in Applied Sciences*, vol. 5, no. 4, pp. 497–518, 1995.
- [7] J. I. Díaz, G. Galiano, and A. Jüngel, "On a quasilinear degenerate system arising in semiconductors theory—part I: existence and uniqueness of solutions," *Nonlinear Analysis: Real World Applications*, vol. 2, no. 3, pp. 305–336, 2001.
- [8] J. I. Díaz, G. Galiano, and A. Jüngel, "On a quasilinear degenerate system arising in semiconductor theory—part II: localization of vacuum solutions," *Nonlinear Analysis: Theory, Methods & Applications*, vol. 36, no. 5, pp. 569–594, 1999.
- [9] P. Guan and B. Wu, "Existence of weak solutions to a degenerate time-dependent semiconductor equations with temperature effect," *Journal of Mathematical Analysis and Applications*, vol. 332, no. 1, pp. 367–380, 2007.
- [10] S. N. Antontsev and S. I. Shmarev, "A model porous medium equation with variable exponent of nonlinearity: existence, uniqueness and localization properties of solutions," *Nonlinear Analysis: Theory, Methods & Applications*, vol. 60, no. 3, pp. 515–545, 2005.
- [11] R. Ferreira and J. L. Vazquez, "Extinction behaviour for fast diffusion equations with absorption," *Nonlinear Analysis: Theory, Methods & Applications*, vol. 43, no. 8, pp. 943–985, 2001.
- [12] H. Yuan, S. Lian, W. Gao, X. Xu, and C. Cao, "Extinction and positivity for the evolution *p*-Laplacian equation in Rⁿ," *Nonlinear Analysis: Theory, Methods & Applications*, vol. 60, no. 6, pp. 1085–1091, 2005.
- [13] Y. Li and J. Wu, "Extinction for fast diffusion equations with nonlinear sources," *Electronic Journal of Differential Equations*, vol. 2005, no. 23, pp. 1–7, 2005.
- [14] W. Liu, M. Wang, and B. Wu, "Extinction and decay estimates of solutions for a class of porous medium equations," *Journal of Inequalities and Applications*, vol. 2007, Article ID 87650, 8 pages, 2007.
- [15] O. A. Ladyžhenskaya and N. N. Vral'tseva, Linear and Quasilinear Elliptic Equations, Academic Press, New York, NY, USA, 1968.