RESEARCH

Open Access



New results on the continuous Weinstein wavelet transform

Hatem Mejjaoli^{1*} and Ahmedou Ould Ahmed Salem²

Dedicated to the Professor Néjib Ben Salem

*Correspondence: hatem.mejjaoli@yahoo.fr *College of Sciences, Department of Mathematics, Taibah University, PO BOX 30002, Al Madinah Al Munawarah, Saudi Arabia Full list of author information is available at the end of the article

Abstract

We consider the continuous wavelet transform S_h^W associated with the Weinstein operator. We introduce the notion of localization operators for S_h^W . In particular, we prove the boundedness and compactness of localization operators associated with the continuous wavelet transform. Next, we analyze the concentration of S_h^W on sets of finite measure. In particular, Benedicks-type and Donoho-Stark's uncertainty principles are given. Finally, we prove many versions of Heisenberg-type uncertainty principles for S_h^W .

MSC: Primary 44A05; secondary 42B10

Keywords: Weinstein operator; Weinstein wavelet transform; localization operators; Schatten-von Neumann class; Heisenberg's type inequalities

1 Introduction

In this paper, we consider the Weinstein operator (also called the Laplace-Bessel differential operator (see [1])) defined on $\mathbb{R}^{d-1} \times (0, \infty)$ by

$$\begin{split} & \bigtriangleup_{\beta} := \sum_{i=1}^{d} \frac{\partial^{2}}{\partial x_{i}^{2}} + \frac{2\beta + 1}{x_{d}} \frac{\partial}{\partial x_{d}} \\ & = \bigtriangleup_{x'} + \mathcal{L}_{\beta, x_{d}}, \quad \beta > -\frac{1}{2}, \end{split}$$

where $\triangle_{x'}$ is the Laplace operator on \mathbb{R}^{d-1} , and \mathcal{L}_{β,x_d} the Bessel operator on $(0,\infty)$ given by

$$\mathcal{L}_{\beta,x_d} := \frac{d^2}{dx_d^2} + \frac{2\beta + 1}{x_d} \frac{d}{dx_d}, \quad \beta > -\frac{1}{2}.$$
 (1.1)

The Weinstein operator \triangle_{β} has several applications in pure and applied mathematics especially in fluid mechanics (see [2]).

The harmonic analysis associated with the Weinstein operator is studied by Ben Nahia and Ben Salem [3, 4]. In particular the authors have introduced and studied the generalized Fourier transform associated with the Weinstein operator. This transform is called the Weinstein transform.



© The Author(s) 2017. This article is distributed under the terms of the Creative Commons Attribution 4.0 International License (http://creativecommons.org/licenses/by/4.0/), which permits unrestricted use, distribution, and reproduction in any medium, provided you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons license, and indicate if changes were made.

In the classical setting, the notion of wavelets was first introduced by Morlet, a French petroleum engineer at ELF-Aquitaine, in connection with his study of seismic traces. The mathematical foundations were given by Grossmann and Morlet [5]. The harmonic analyst Meyer and many other mathematicians became aware of this theory, and they recognized many classical results inside it (see [6–8]). Classical wavelets have wide applications, ranging from signal analysis in geophysics and acoustics to quantum theory and pure mathematics (see [9–11] and the references therein).

Next, the theory of wavelets and continuous wavelet transforms has been extended to hypergroups, in particular, to the Chébli-Trimèche hypergroups (see [12]).

Recently in [13] the authors have introduced and studied the Weinstein wavelet transform S_h^W . In the same paper, for the transform S_h^W , the authors have proved the Plancherel and inversion formulas.

Nowadays, the continuous wavelet transform is one of the useful subjects in harmonic analysis. In this paper we present only two subjects.

The first subject is the new uncertainty principles involving time-frequency representations.

The second subject is the localization operators. These operators were initiated by Daubechies [14–16], and detailed in the book [17] by Wong.

As the harmonic analysis associated with the Weinstein operator has known remarkable development, it is a natural question whether there exist an equivalent of the theory of localization operators and new uncertainty principles for the continuous wavelet transform relating to this harmonic analysis.

The purpose of the present paper is twofold. On one hand, we want to study many versions of quantitative uncertainty principles for the continuous Weinstein wavelet transform. On the other hand, we want to study the localization operators for the continuous Weinstein wavelet transform.

The remainder of this paper is arranged as follows. In Section 2, we recall the main results about the harmonic analysis associated with the Weinstein operator. In Section 3, we introduce and study the two-localization operators associated with the Weinstein continuous wavelet transform. More precisely, the Schatten-von Neumann properties of these two-wavelet localization operators are established, and for trace class localization operators, the traces and the trace class norm inequalities are presented. In Section 4, we study the quantitative analysis of the continuous Weinstein wavelet transform and time-frequency concentration. In particular, we give results on sets of finite measure and Donoho-Stark and Benedicks-type uncertainty principles. In Sections 5 and 6, we prove many versions of the Heisenberg uncertainty inequalities for the Weinstein continuous wavelet transform. Finally, some conclusions are drawn in Section 7.

2 Preliminaries

To confirm the basic and standard notations, we briefly overview the Weinstein operator and related harmonic analysis. The main references are [3, 4].

2.1 Harmonic analysis associated with the Weinstein operator

In this subsection, we collect some notation and results on the Weinstein kernel, the Weinstein transform, and the Weinstein convolution.

We use the following notation:

$$\begin{split} \overline{\mathbb{R}^d_+} &= \mathbb{R}^{d-1} \times [0,\infty); \\ x &= (x_1, \dots, x_{d-1}, x_d) = (x', x_d) \in \overline{\mathbb{R}^d_+}; \\ C_*(\mathbb{R}^d) \text{ is the space of continuous functions on } \mathbb{R}^d \text{ even with respect to the last variable.} \\ C^p_*(\mathbb{R}^d) \text{ is the space of functions of class } C^p \text{ on } \mathbb{R}^d \text{ even with respect to the last variable.} \end{split}$$

 $\mathcal{E}_*(\mathbb{R}^d)$ is the space of C^∞ -functions on \mathbb{R}^d even with respect to the last variable.

 $S_*(\mathbb{R}^d)$ is the Schwartz space of rapidly decreasing functions on \mathbb{R}^d even with respect to the last variable.

 $D_*(\mathbb{R}^d)$ is the space of C^∞ -functions on \mathbb{R}^d with compact support and even with respect to the last variable.

The Weinstein kernel Λ is given by

$$\Lambda(x,z) := e^{i\langle x',z'\rangle} j_{\beta}(x_d z_d) \quad \text{for all } (x,z) \in \mathbb{R}^d \times \mathbb{C}^d,$$
(2.1)

where j_{β} is the normalized Bessel function. The Weinstein kernel satisfies the following properties:

(i) For all $z, t \in \mathbb{C}^d$, we have

$$\Lambda(z,t) = \Lambda(t,z), \qquad \Lambda(z,0) = 1, \text{ and}$$

$$\Lambda(\lambda z,t) = \Lambda(z,\lambda t) \quad \text{for all } \lambda \in \mathbb{C}.$$

$$(2.2)$$

(ii) For all $\nu \in \mathbb{N}^d$, $x \in \mathbb{R}^d$, and $z \in \mathbb{C}^d$, we have

$$|D_{z}^{\nu}\Lambda(x,z)| \le ||x||^{|\nu|} \exp(||x|| || \operatorname{Im} z||),$$
(2.3)

where $D_z^{\nu} = \frac{\partial^{\nu}}{\partial z_1^{\nu_1} \cdots \partial z_d^{\nu_d}}$ and $|\nu| = \nu_1 + \cdots + \nu_d$. In particular,

$$\left|\Lambda(x,y)\right| \le 1 \quad \text{for all } x, y \in \mathbb{R}^d.$$
(2.4)

We denote by $L^p_{\beta}(\overline{\mathbb{R}^d_+})$, $1 \le p \le \infty$, the space of measurable functions on $\overline{\mathbb{R}^d_+}$ such that

$$\begin{split} \|f\|_{L^p_{\beta}(\overline{\mathbb{R}^d_+})} &= \left(\int_{\overline{\mathbb{R}^d_+}} \left|f(y)\right|^p d\lambda_{\beta}(y)\right)^{\frac{1}{p}} < \infty \quad \text{if } 1 \le p < \infty, \\ \|f\|_{L^\infty_{\beta}(\overline{\mathbb{R}^d_+})} &= \mathrm{ess}\sup_{y \in \overline{\mathbb{R}^d_+}} \left|f(y)\right| < \infty \quad \text{if } p = \infty, \end{split}$$

where $d\lambda_{\beta}$ is the measure on $\overline{\mathbb{R}^d_+}$ given by

$$d\lambda_{\beta}(y) := y_d^{2\beta+1} dy' \, dy_d.$$

The Weinstein transform is given for f in $L^1_{\beta}(\overline{\mathbb{R}^d_+})$ by

$$\mathcal{F}_{W}(f)(\xi) = \frac{1}{C(\beta)} \int_{\mathbb{R}^{d+1}_{+}} f(y) \Lambda(-\xi, y) \, d\lambda_{\beta}(y) \quad \text{for all } \xi \in \overline{\mathbb{R}^{d}_{+}},$$
(2.5)

where

$$C(\beta) = \int_{\mathbb{R}^d_+} e^{-\frac{\|y\|^2}{2}} d\lambda_{\beta}(y) = 2^{\beta} \Gamma(\beta+1)(2\pi)^{\frac{d-1}{2}}.$$

Some basic properties of this transform are the following:

(i) For f in $L^1_\beta(\overline{\mathbb{R}^d_+})$,

$$\left\|\mathcal{F}_{W}(f)\right\|_{L^{\infty}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \leq \frac{1}{C(\beta)} \left\|f\right\|_{L^{1}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}.$$
(2.6)

(ii) For all f in $L^1_\beta(\overline{\mathbb{R}^d_+})$, if $\mathcal{F}_W(f)$ belongs to $L^1_\beta(\overline{\mathbb{R}^d_+})$, then

$$f(y) = \frac{1}{C(\beta)} \int_{\mathbb{R}^d_+} \mathcal{F}_W(f)(\xi) \Lambda(\xi, y) \, d\lambda_\beta(\xi) \quad \text{a.e.}$$
(2.7)

(iii) For $f \in \mathcal{S}_*(\mathbb{R}^d)$, if we define

$$\overline{\mathcal{F}_W}(f)(\xi) = \mathcal{F}_W(f)(-\xi),$$

then

$$\mathcal{F}_W \overline{\mathcal{F}_W} = \overline{\mathcal{F}_W} \mathcal{F}_W = Id. \tag{2.8}$$

Proposition 2.1

 (i) The Weinstein transform 𝓕_W is a topological isomorphism from 𝔅_{*}(ℝ^d) onto itself, and for all f in 𝔅_{*}(ℝ^d),

$$\int_{\overline{\mathbb{R}^d_+}} \left| f(y) \right|^2 d\lambda_\beta(y) = \int_{\overline{\mathbb{R}^d_+}} \left| \mathcal{F}_W(f)(\xi) \right|^2 d\lambda_\beta(\xi).$$
(2.9)

(ii) The Weinstein transform $f \to \mathcal{F}_W(f)$ induces an isometric isomorphism from $L^2_\beta(\mathbb{R}^d_+)$ onto itself.

For a function $f \in S_*(\mathbb{R}^d)$ and $y \in \overline{\mathbb{R}^d_+}$, the generalized translation $\tau_y f$ is defined by the following relation:

$$\forall \xi \in \mathbb{R}^d, \quad \mathcal{F}_W(\tau_y f)(\xi) = \Lambda(\xi, y) \mathcal{F}_W(f)(\xi).$$
(2.10)

By using the generalized translation, we define the generalized convolution product $f *_W g$ of functions $f, g \in L^1_\beta(\overline{\mathbb{R}^d_+})$ as follows:

$$f *_{W} g(x) = \int_{\overline{\mathbb{R}^{d}_{+}}} \tau_{x} f\left(-y', y_{d}\right) g(y) \, d\lambda_{\beta}(y), \quad x \in \overline{\mathbb{R}^{d}_{+}}.$$
(2.11)

This convolution is commutative and associative. Moreover, we have the following:

Proposition 2.2

- (i) For all $f,g \in L^1_{\beta}(\overline{\mathbb{R}^d_+}), f *_W g$ belongs to $L^1_{\beta}(\overline{\mathbb{R}^d_+}), and$ $\forall \xi \in \overline{\mathbb{R}^d_+}, \quad \mathcal{F}_W(f *_W g)(\xi) = \mathcal{F}_W(f)(\xi)\mathcal{F}_W(g)(\xi).$ (2.12)
- (ii) Let $1 \le p, q, r \le \infty$ be such that $\frac{1}{p} + \frac{1}{q} \frac{1}{r} = 1$. If $f \in L^p_\beta(\overline{\mathbb{R}^d_+})$ and $g \in L^q_\beta(\overline{\mathbb{R}^d_+})$, then $f *_W g \in L^r_\beta(\overline{\mathbb{R}^d_+})$, and

$$\|f *_{W} g\|_{L^{p}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \leq \|f\|_{L^{p}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \|g\|_{L^{q}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}.$$
(2.13)

(iii) Let f, g be in $L^2_{\beta}(\overline{\mathbb{R}^d_+})$. Then $f *_W g \in L^2_{\beta}(\overline{\mathbb{R}^d_+})$ if and only if $\mathcal{F}_W(f)\mathcal{F}_W(g) \in L^2_{\beta}(\overline{\mathbb{R}^d_+})$ and (2.12) holds.

Definition 2.1 Let *E*, *F* be two measurable subsets of $\overline{\mathbb{R}^d_+}$. Then:

- (1) We say that (E, F) is weakly annihilating pair if supp $f \subset E$ and supp $\mathcal{F}_W(f) \subset F$ implies f = 0.
- (2) We say that (E, F) is strongly annihilating pair if there exists a constant $C := C_{\beta}(E, F) > 0$ such that, for every function $f \in L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})$,

$$C(\left\|\mathcal{F}_{W}(f)\right\|_{L^{2}_{\beta}(F^{c})}^{2}+\left\|f\right\|_{L^{2}_{\beta}(E^{c})}^{2}) \geq \left\|f\right\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}^{2}.$$
(2.14)

The constant $C_{\beta}(E, F)$ is called the annihilation constant of (E, F).

Proposition 2.3 [18] Let E, F be two measurable subsets of $\overline{\mathbb{R}^d_+}$ with

$$\int_E d\lambda_\beta(y) < \infty \quad and \quad \int_F d\lambda_\beta(\xi) < \infty.$$

Then the pair (E, F) is a strongly annihilating pair.

2.2 Weinstein wavelets

This subsection gives an introduction to the theory of Weinstein wavelets. The main reference is [13].

Definition 2.2 A Weinstein wavelet on $\overline{\mathbb{R}^d_+}$ is a measurable function *h* on $\overline{\mathbb{R}^d_+}$ satisfying, for almost all $\xi \in \overline{\mathbb{R}^d_+}$, the condition

$$0 < \mathcal{C}_h = \int_0^\infty \left| \mathcal{F}_W(h)(t\xi) \right|^2 \frac{dt}{t} < \infty.$$
(2.15)

Example 2.1 Let $h \in L^2_{\beta}(\mathbb{R}^d_+)$ be given by

$$h := \mathcal{F}_{W}^{-1}(\|y\|^{2} e^{-\frac{1}{2}\|y\|^{2}}).$$
(2.16)

By a simple calculations we see that, for almost all $x \in \mathbb{R}^d_+$,

$$\int_0^\infty \left|\mathcal{F}_W(h)(\lambda x)\right|^2 \frac{d\lambda}{\lambda} < \infty.$$

Thus *h* is an example of a Weinstein wavelet on \mathbb{R}^{d}_{+} .

Let b > 0, and let h be in $L^2_{\beta}(\overline{\mathbb{R}^d_+})$. The dilation of h by b is defined by

$$\forall y \in \overline{\mathbb{R}^d_+}, \quad h_b(y) = \frac{1}{b^{2\beta+1+d}} h\left(\frac{y}{b}\right). \tag{2.17}$$

It is easy to see that $h_b \in L^2_\beta(\overline{\mathbb{R}^d_+})$ and

$$\forall \xi \in \mathbb{R}^d_+, \quad \mathcal{F}_W(h_b)(\xi) = \mathcal{F}_W(h)(b\xi). \tag{2.18}$$

We introduce the family $h_{b,y}$, b > 0 and $y \in \overline{\mathbb{R}^d_+}$, of Weinstein wavelets on $\overline{\mathbb{R}^d_+}$ in $L^2_\beta(\overline{\mathbb{R}^d_+})$ defined by

$$\forall x \in \overline{\mathbb{R}^d_+}, \quad h_{b,y}(x) = b^{\beta + \frac{d+1}{2}} \tau_x h_b(-y', y_d).$$
(2.19)

Note that

$$\forall b > 0, \forall y \in \overline{\mathbb{R}^d_+}, \quad \|h_{b,y}\|_{L^2_\beta(\overline{\mathbb{R}^d_+})} \le \|h\|_{L^2_\beta(\overline{\mathbb{R}^d_+})}.$$

$$(2.20)$$

Notation We denote:

$$\Omega_{d+1} = \left\{ (b, y) = (b, y_1, \dots, y_d) \in \mathbb{R}^{d+1}, b > 0 \text{ and } y \in \overline{\mathbb{R}^d_+} \right\}$$

 $L^p_{\mu_\beta}(\Omega_{d+1})$, $1 \le p \le \infty$, is the space of measurable functions f on Ω_{d+1} such that

$$\begin{split} \|f\|_{L^p_{\mu_{\beta}}(\Omega_{d+1})} &\coloneqq \left(\int_{\Omega_{d+1}} \left|f(b,y)\right|^p d\mu_{\beta}(b,y)\right)^{\frac{1}{p}} < \infty, \quad 1 \le p < \infty, \\ \|f\|_{L^{\infty}_{\mu_{\beta}}(\Omega_{d+1})} &\coloneqq \operatorname{ess} \sup_{(b,y) \in \Omega_{d+1}} \left|f(b,y)\right| < \infty, \quad p = \infty, \end{split}$$

where the measure $d\mu_{\beta}$ is defined by

$$\forall (b,y) \in \Omega_{d+1}, \quad d\mu_{\beta}(b,y) = \frac{y_d^{2\beta+1} dy db}{b^{2\beta+d+2}}.$$

Definition 2.3 Let *h* be a Weinstein wavelet on $\overline{\mathbb{R}^d_+}$ in $L^2_\beta(\overline{\mathbb{R}^d_+})$. The Weinstein continuous wavelet transform \mathcal{S}^W_h on $\overline{\mathbb{R}^d_+}$ is defined for regular functions *f* on $\overline{\mathbb{R}^d_+}$ by

$$\mathcal{S}_{h}^{W}(f)(b,y) = \int_{\mathbb{R}^{d}_{+}} f(y)\overline{h_{b,y}(x)} \, d\lambda_{\beta}(x), \quad (b,y) \in \Omega_{d+1}.$$
(2.21)

It is easy to see that

$$\forall (b, y) \in \Omega_{d+1}, \quad \mathcal{S}_h^W(f)(b, y) = b^{\beta + \frac{d+1}{2}} f *_W \overline{h_b}(y). \tag{2.22}$$

Proposition 2.4 Let f be in $L^2_{\beta}(\overline{\mathbb{R}^d_+})$, and let h be a Weinstein wavelet. Then we have

$$\mathcal{S}_{h}^{W}(f)(b,y) = b^{\beta + \frac{d+1}{2}} \mathcal{F}_{W}^{-1} \big(\mathcal{F}_{W}(f)(\xi) \mathcal{F}_{W}(\overline{h})(b\xi) \big)(y), \quad (b,y) \in \Omega_{d+1},$$
(2.23)

and

$$\|\mathcal{S}_{h}^{W}f\|_{L^{\infty}_{\mu_{\beta}}(\Omega_{d+1})} \le \|f\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \|h\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}.$$
(2.24)

Proposition 2.5 (Covariance properties) Let h be a Weinstein wavelet. The transformation S_h^W is a bounded linear operator from $L_{\beta}^2(\overline{\mathbb{R}_+^d})$ into the space of continuous bounded functions on Ω_{d+1} . Moreover, we have the following covariance property: For any $f \in L_{\beta}^2(\overline{\mathbb{R}_+^d})$, $(b, y) \in \Omega_{d+1}$ and $c \in \mathbb{R}^*$,

$$\left[\mathcal{S}_{h}^{W}\left(\frac{1}{c^{\beta+\frac{d}{2}}}f\left(\frac{\cdot}{c}\right)\right)\right](b,y) = \left[\mathcal{S}_{h}^{W}(f)\right]\left(\frac{b}{c},\frac{y}{c}\right).$$
(2.25)

Theorem 2.1 (Plancherel's formula for \mathcal{S}_h^W) Let h be a Weinstein wavelet. For all f in $L^2_{\beta}(\overline{\mathbb{R}^d_+})$, we have

$$\int_{\mathbb{R}^{d}_{+}} |f(x)|^{2} d\lambda_{\beta}(x) = \frac{1}{C_{h}} \int_{0}^{\infty} \int_{\mathbb{R}^{d}_{+}} |\mathcal{S}_{h}^{W}(f)(b,y)|^{2} d\mu_{\beta}(b,y).$$
(2.26)

Corollary 2.1 (Parseval's formula) Let *h* be a Weinstein wavelet. For all f_1, f_2 in $L^2_\beta(\mathbb{R}^d_+)$, we have

$$\int_{\mathbb{R}^d_+} f_1(x)\overline{f_2(x)} \, d\lambda_\beta(x) = \frac{1}{\mathcal{C}_h} \int_0^\infty \int_{\mathbb{R}^d_+} \mathcal{S}_h^W(f_1)(b, y) \overline{\mathcal{S}_h^W(f_2)(b, y)} \, d\mu_\beta(b, y).$$
(2.27)

3 Localization operators for the Weinstein continuous wavelet transform3.1 Preliminaries

We denote by $B(L^2_{\beta}(\overline{\mathbb{R}^d_+}))$ the space of bounded operators from $L^2_{\beta}(\overline{\mathbb{R}^d_+})$ into $L^2_{\beta}(\overline{\mathbb{R}^d_+})$. The singular values $(s_n(M))_{n\in\mathbb{N}}$ of a compact operator $M \in B(L^2_{\beta}(\overline{\mathbb{R}^d_+}))$ are the eigenvalues of the positive self-adjoint operator $|M| = \sqrt{M^*M}$. For $1 \le p < \infty$, the Schatten class S_p is the space of compact operators whose singular values lie in l^p . Hence S_p is equipped with the norm

$$\|M\|_{S_p} := \left(\sum_{j=1}^{\infty} (s_j(M))^p\right)^{\frac{1}{p}}.$$
(3.1)

In particular, S_2 is the space of Hilbert-Schmidt operators, and S_1 is the space of trace-class operators. It is well known that the trace of an operator $M \in S_1$ is defined by

$$\operatorname{tr}(M) = \sum_{j=1}^{\infty} \langle M \nu_j, \nu_j \rangle, \qquad (3.2)$$

where $(v_i)_i$ is any orthonormal basis for $L^2_{\beta}(\overline{\mathbb{R}^d_+})$. Moreover, if *M* is positive, then

$$tr(M) = \|M\|_{S_1}.$$
(3.3)

Moreover, a compact operator M on the Hilbert space $L^2_\beta(\overline{\mathbb{R}^d_+})$ is Hilbert-Schmidt if the positive operator M^*M is in the space of trace class S_1 . Then

$$\|M\|_{HS}^{2} := \|M\|_{S_{2}}^{2} = \|M^{*}M\|_{S_{1}} = \operatorname{tr}(M^{*}M) = \sum_{j=1}^{\infty} \|M\nu_{j}\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}^{2}$$
(3.4)

for any orthonormal basis $(v_j)_j$ basis for $L^2_\beta(\mathbb{R}^d_+)$.

For consistency, we define $S_{\infty} := B(L^2_{\beta}(\overline{\mathbb{R}^d_+}))$ equipped with norm

$$\|M\|_{S_{\infty}} := \sup_{\nu \in L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}}): \|\nu\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} = 1} \|M\nu\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}.$$
(3.5)

3.2 Boundedness

In this section, h and k will be two Weinstein wavelets such that

$$\|h\|_{L^2_\beta(\overline{\mathbb{R}^d_+})} = \|k\|_{L^2_\beta(\overline{\mathbb{R}^d_+})} = 1.$$

Definition 3.1 The localization operator associated with the symbol *a* and two Weinstein wavelets is denoted by $\mathcal{L}_{h,k}(a)$ and defined on $L^2_{\beta}(\overline{\mathbb{R}^d_+})$ by

$$\mathcal{L}_{h,k}(a)(f)(x) = \frac{1}{\sqrt{\mathcal{C}_h \mathcal{C}_k}} \int_{\Omega_{d+1}} a(b, y) \mathcal{S}_h^W(f)(b, y) k_{b,y}(x) \, d\mu_\beta(b, y), \quad x \in \overline{\mathbb{R}_+^d}.$$

Often, it is more convenient to interpret the definition of $\mathcal{L}_{h,k}(a)$ in a weak sense, that is, for f, g in $L^2_\beta(\overline{\mathbb{R}^d_+})$,

$$\left\langle \mathcal{L}_{h,k}(a)(f),g\right\rangle_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} = \frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}} \int_{\Omega_{d+1}} a(b,y)\mathcal{S}_{h}^{W}(f)(b,y)\overline{\mathcal{S}_{k}^{W}(g)(b,y)} \,d\mu_{\beta}(b,y), \quad g \in L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}}).$$
(3.6)

Proposition 3.1 *The adjoint of* $\mathcal{L}_{h,k}(a)$ *is* $\mathcal{L}_{k,h}(\overline{a})$ *.*

Proof For all *f*, *g* in $L^2_{\beta}(\overline{\mathbb{R}^d_+})$, from (3.6) it immediately follows that

$$\begin{split} \left\langle \mathcal{L}_{h,k}(a)(f),g\right\rangle_{L^2_{\beta}(\overline{\mathbb{R}^d_+})} &= \frac{1}{\sqrt{\mathcal{C}_h \mathcal{C}_k}} \int_{\Omega_{d+1}} a(b,y) \mathcal{S}_h^W(f)(b,y) \overline{\mathcal{S}_k^W(g)(b,y)} \, d\mu_{\beta}(b,y) \\ &= \overline{\frac{1}{\sqrt{\mathcal{C}_h \mathcal{C}_k}} \int_{\Omega_{d+1}} \overline{a(b,y)} \mathcal{S}_k^W(g)(b,y) \overline{\mathcal{S}_h^W(f)(b,y)} \, d\mu_{\beta}(b,y)} \\ &= \overline{\left\langle \mathcal{L}_{k,h}(\overline{a})(g),f \right\rangle_{L^2_{\beta}(\overline{\mathbb{R}^d_+})}} = \left\langle f, \mathcal{L}_{k,h}(\overline{a})(g) \right\rangle_{L^2_{\beta}(\overline{\mathbb{R}^d_+})}. \end{split}$$

Thus we get

$$\mathcal{L}_{h,k}^*(a) = \mathcal{L}_{k,h}(\overline{a}). \tag{3.7}$$

The objective of this subsection is to prove that the operators

$$\mathcal{L}_{h,k}(a): L^2_{\beta}\left(\overline{\mathbb{R}^d_+}\right) \to L^2_{\beta}\left(\overline{\mathbb{R}^d_+}\right)$$

are bounded for all symbols $a \in L^p_{\mu_\beta}(\Omega_{d+1})$. We first consider this problem for a in $L^1_{\mu_\beta}(\Omega_{d+1})$ and next in $L^\infty_{\mu_\beta}(\Omega_{d+1})$, and then by interpolation theory we deduce the result.

Proposition 3.2 Let a be in $L^1_{\mu_{\beta}}(\Omega_{d+1})$. Then the localization operator $\mathcal{L}_{h,k}(a)$ is in S_{∞} , and

$$\left\|\mathcal{L}_{h,k}(a)\right\|_{S_{\infty}} \leq \frac{1}{\sqrt{\mathcal{C}_h \mathcal{C}_k}} \|a\|_{L^1_{\mu_\beta}(\Omega_{d+1})}.$$

Proof Let *f* and *g* be in $L^2_\beta(\overline{\mathbb{R}^d_+})$. From relations (3.6) and (2.24) we have

$$\begin{split} \left| \left\langle \mathcal{L}_{h,k}(a)(f), g \right\rangle_{L^2_{\beta}(\overline{\mathbb{R}^d_+})} \right| &\leq \frac{1}{\sqrt{\mathcal{C}_h \mathcal{C}_k}} \int_{\Omega_{d+1}} \left| a(b,y) \right| \left| \mathcal{S}_h^W(f)(b,y) \right| \left| \overline{\mathcal{S}_k^W(g)(b,y)} \right| d\mu_{\beta}(b,y) \\ &\leq \frac{1}{\sqrt{\mathcal{C}_h \mathcal{C}_k}} \left\| \mathcal{S}_h^W(f) \right\|_{L^\infty_{\mu\beta}(\Omega_{d+1})} \left\| \mathcal{S}_k^W(g) \right\|_{L^\infty_{\mu\beta}(\Omega_{d+1})} \| a \|_{L^1_{\mu\beta}(\Omega_{d+1})} \\ &\leq \frac{1}{\sqrt{\mathcal{C}_h \mathcal{C}_k}} \left\| f \right\|_{L^2_{\beta}(\overline{\mathbb{R}^d_+})} \| g \|_{L^2_{\beta}(\overline{\mathbb{R}^d_+})} \| a \|_{L^1_{\mu\beta}(\Omega_{d+1})}. \end{split}$$

Thus,

$$\left\|\mathcal{L}_{h,k}(a)\right\|_{S_{\infty}} \leq \frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}} \|a\|_{L^{1}_{\mu_{\beta}}(\Omega_{d+1})}.$$

Proposition 3.3 Let a be in $L^{\infty}_{\mu_{\beta}}(\Omega_{d+1})$. Then the localization operator $\mathcal{L}_{h,k}(a)$ is in S_{∞} , and

$$\left\|\mathcal{L}_{h,k}(a)\right\|_{S_{\infty}}\leq \|a\|_{L^{\infty}_{\mu_{\beta}}(\Omega_{d+1})}.$$

Proof For all functions f and g in $L^2_{\beta}(\overline{\mathbb{R}^d_+})$, by Hölder's inequality we have

$$\begin{split} \left| \left\langle \mathcal{L}_{h,k}(a)(f), g \right\rangle_{L^2_{\beta}(\overline{\mathbb{R}^d_+})} \right| &\leq \frac{1}{\sqrt{\mathcal{C}_h \mathcal{C}_k}} \int_{\Omega_{d+1}} \left| a(b, y) \right| \left| \mathcal{S}_h^W(f)(b, y) \right| \left| \overline{\mathcal{S}_k^W(g)(b, y)} \right| d\mu_{\beta}(b, y) \\ &\leq \frac{1}{\sqrt{\mathcal{C}_h \mathcal{C}_k}} \left\| a \right\|_{L^\infty_{\mu_{\beta}}(\Omega_{d+1})} \left\| \mathcal{S}_h^W(f) \right\|_{L^2_{\mu_{\beta}}(\Omega_{d+1})} \left\| \mathcal{S}_k^W(g) \right\|_{L^2_{\mu_{\beta}}(\Omega_{d+1})}. \end{split}$$

Applying the Plancherel formula (2.26) for \mathcal{S}_h^W and \mathcal{S}_k^W , we obtain

$$\left|\left(\mathcal{L}_{h,k}(a)(f),g\right)_{L^2_\beta(\overline{\mathbb{R}^d_+})}\right| \leq \|a\|_{L^\infty_{\mu\beta}(\Omega_{d+1})} \|f\|_{L^2_\beta(\overline{\mathbb{R}^d_+})} \|g\|_{L^2_\beta(\overline{\mathbb{R}^d_+})}.$$

Thus,

$$\left\|\mathcal{L}_{h,k}(a)\right\|_{S_{\infty}} \leq \|a\|_{L^{\infty}_{\mu_{\beta}}(\Omega_{d+1})}.$$

Remark 3.1 For $a = a_1 + a_\infty$ in $L^1_{\mu_\beta}(\Omega_{d+1}) + L^\infty_{\mu_\beta}(\Omega_{d+1})$, from Proposition 3.2 and Proposition 3.3 we deduce that $\mathcal{L}_{h,k}(a)$ is in S_∞ and

$$\left\|\mathcal{L}_{h,k}(a)\right\|_{S_{\infty}} \leq \frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}} \|a_{1}\|_{L^{1}_{\mu_{\beta}}(\Omega_{d+1})} + \|a_{\infty}\|_{L^{\infty}_{\mu_{\beta}}(\Omega_{d+1})} \leq \max\left\{\frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}}, 1\right\} \|a\|,$$

where

$$\|a\| = \inf \{ \|a_1\|_{L^1_{\mu_\beta}(\Omega_{d+1})} + \|a_\infty\|_{L^\infty_{\mu_\beta}(\Omega_{d+1})} : a = a_1 + a_\infty, a_1 \in L^1_{\mu_\beta}(\Omega_{d+1}), a_\infty \in L^\infty_{\mu_\beta}(\Omega_{d+1}) \}.$$

We now can prove that $\mathcal{L}_{h,k}(a)$ is in S_{∞} for every symbol a in $L^p_{\mu_{\beta}}(\Omega_{d+1})$, $1 \le p \le \infty$.

Theorem 3.1 Let a be in $L^p_{\mu_\beta}(\Omega_{d+1}), p \in [1, \infty]$. Then there exists a bounded linear operator $\mathcal{L}_{h,k}(a): L^2_{\beta}(\overline{\mathbb{R}^d_+}) \to L^2_{\beta}(\overline{\mathbb{R}^d_+})$ such that

$$\left\|\mathcal{L}_{h,k}(a)\right\|_{S_{\infty}} \leq \left(\frac{1}{\mathcal{C}_{h}\mathcal{C}_{k}}\right)^{\frac{1}{2p}} \left\|a\right\|_{L^{p}_{\mu_{\beta}}(\Omega_{d+1})}.$$

Proof Let *f* be in $L^2_{\beta}(\overline{\mathbb{R}^d_+})$. We consider the operator

$$\mathcal{T}: L^1_{\mu_{\beta}}(\Omega_{d+1}) \cap L^{\infty}_{\mu_{\beta}}(\Omega_{d+1}) \to L^2_{\beta}(\mathbb{R}^d)$$

given by

$$\mathcal{T}(a) := \mathcal{L}_{h,k}(a)(f).$$

Then, by Proposition 3.2 and Proposition 3.3,

$$\|\mathcal{T}(a)\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \leq \frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}} \|f\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \|a\|_{L^{1}_{\mu_{\beta}}(\Omega_{d+1})}$$
(3.8)

and

$$\|\mathcal{T}(a)\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \leq \|f\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \|a\|_{L^{\infty}_{\mu_{\beta}}(\Omega_{d+1})}.$$
(3.9)

Therefore, by (3.8), (3.9), and the Riesz-Thorin interpolation argument ([19, Theorem 2], see also [17, Theorem 2.11]) we get that \mathcal{T} may be uniquely extended to a linear operator on $L^p_{\mu_{\beta}}(\Omega_{d+1}), 1 \le p \le \infty$, and we have

$$\|\mathcal{L}_{h,k}(a)(f)\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} = \|\mathcal{T}(a)\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \le \left(\frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}}\right)^{\frac{1}{p}} \|f\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \|a\|_{L^{p}_{\mu_{\beta}}(\Omega_{d+1})}.$$
(3.10)

Since (3.10) is true for all $f \in L^2_\beta(\overline{\mathbb{R}^d_+})$, the result is proved.

3.3 Schatten-von Neumann properties for $\mathcal{L}_{h,k}(a)$

Let us begin with the following statement.

Proposition 3.4 Let a be in $L^1_{\mu_{\beta}}(\Omega_{d+1})$. Then the localization operator

$$\mathcal{L}_{h,k}(a): L^2_{\beta}\left(\overline{\mathbb{R}^d_+}\right) \to L^2_{\beta}\left(\overline{\mathbb{R}^d_+}\right)$$

is in S_2 , and

$$\left\|\mathcal{L}_{h,k}(a)\right\|_{S_2} \leq \frac{1}{\sqrt{\mathcal{C}_h \mathcal{C}_k}} \|a\|_{L^1_{\mu_\beta}(\Omega_{d+1})}$$

Proof Let $\{\phi_j, j = 1, 2, ...\}$ be an orthonormal basis for $L^2_\beta(\overline{\mathbb{R}^d_+})$. Then by (3.6), Fubini's theorem, Parseval's identity, and relations (2.21) and (3.7) we have

$$\begin{split} &\sum_{j=1}^{\infty} \left\| \mathcal{L}_{h,k}(a)(\phi_{j}) \right\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}}))}^{2} \\ &= \sum_{j=1}^{\infty} \langle \mathcal{L}_{h,k}(a)(\phi_{j}), \mathcal{L}_{h,k}(a)(\phi_{j}) \rangle_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \\ &= \sum_{j=1}^{\infty} \frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}} \int_{\Omega_{d+1}} a(b,y) \langle \phi_{j}, h_{b,y} \rangle_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \overline{\langle \mathcal{L}_{h,k}(a)(\phi_{j}), k_{b,y} \rangle_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}} \, d\mu_{\beta}(b,y) \\ &= \frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}} \int_{\Omega_{d+1}} a(b,y) \sum_{j=1}^{\infty} \langle \phi_{j}, h_{b,y} \rangle_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \langle \mathcal{L}^{*}_{h,k}(a)(k_{b,y}), \phi_{j} \rangle_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \, d\mu_{\beta}(b,y) \\ &= \frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}} \int_{\Omega_{d+1}} a(b,y) \langle \mathcal{L}^{*}_{h,k}(a)k_{b,y}, h_{b,y} \rangle_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \, d\mu_{\beta}(b,y). \end{split}$$

Thus, we get

$$\sum_{j=1}^{\infty} \left\| \mathcal{L}_{h,k}(a)(\phi_{j}) \right\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}^{2} \leq \frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}} \int_{\Omega_{d+1}} \left| a(b,y) \right| \left\| \mathcal{L}^{*}_{h,k}(a) \right\|_{S_{\infty}} d\mu_{\beta}(b,y)$$
$$= \frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}} \left\| \mathcal{L}^{*}_{h,k}(a) \right\|_{S_{\infty}} \|a\|_{L^{1}_{\mu_{\beta}}(\Omega_{d+1})} < \infty.$$
(3.11)

So by (3.11) and Proposition 2.8 in Wong [17]

$$\mathcal{L}_{h,k}(a): L^2_{\beta}\left(\overline{\mathbb{R}^d_+}\right) \to L^2_{\beta}\left(\overline{\mathbb{R}^d_+}\right)$$

is in the Hilbert-Schmidt class S_2 and hence compact.

Proposition 3.5 Let a be a symbol in $L^p_{\mu_{\beta}}(\Omega_{d+1})$, $p \in [1, \infty)$. Then the operator $\mathcal{L}_{h,k}(a)$ is compact.

Proof Let *a* be in $L^p_{\mu_{\beta}}(\Omega_{d+1})$, and let a sequence $(a_n)_{n \in \mathbb{N}} \in L^1_{\mu_{\beta}}(\Omega_{d+1}) \cap L^{\infty}_{\mu_{\beta}}(\Omega_{d+1})$ be such that $a_n \to a$ in $L^p_{\mu_{\beta}}(\Omega_{d+1})$ as $n \to \infty$. Then by Theorem 3.1

$$\left\|\mathcal{L}_{h,k}(a_n)-\mathcal{L}_{h,k}(a)\right\|_{S_{\infty}}\leq \left(rac{1}{\sqrt{\mathcal{C}_h\mathcal{C}_k}}
ight)^{rac{1}{p}}\left\|a_n-a\|_{L^p_{\mu_{eta}}(\Omega_{d+1})}.$$

Hence $\mathcal{L}_{h,k}(a_n) \to \mathcal{L}_{h,k}(a)$ in S_{∞} as $n \to \infty$. On the other hand, since by Proposition 3.4 $\mathcal{L}_{h,k}(a_n)$ is in S_2 and hence compact, it follows that $\mathcal{L}_{h,k}(a)$ is compact. \Box

Theorem 3.2 Let a be in $L^1_{\mu_{\beta}}(\Omega_{d+1})$. Then

$$\frac{2}{\mathcal{C}_{h} + \mathcal{C}_{k}} \|\widetilde{a}\|_{L^{1}_{\mu_{\beta}}(\Omega_{d+1})} \leq \left\| \mathcal{L}_{h,k}(a) \right\|_{S_{1}} \leq \frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}} \|a\|_{L^{1}_{\mu_{\beta}}(\Omega_{d+1})},$$
(3.12)

where \tilde{a} is given by

$$\widetilde{a}(b,y) = \left\langle \mathcal{L}_{h,k}(a)(h_{b,y}), k_{b,y} \right\rangle_{L^2_\beta(\overline{\mathbb{R}^d_+})}, \quad (b,y) \in \Omega_{d+1}.$$

Proof Since *a* is in $L^1_{\mu_\beta}(\Omega_{d+1})$, $\mathcal{L}_{h,k}(a)$ is in S_2 by Proposition 3.4. By [17, Theorem 2.2] there exist an orthonormal basis { $\phi_{j,j} = 1, 2, ...$ } for the orthogonal complement of the kernel of $\mathcal{L}_{h,k}(a)$, consisting of eigenvectors of $|\mathcal{L}_{h,k}(a)|$, and an orthonormal set { $\varphi_{j,j} = 1, 2, ...$ } in $L^2_{\beta}(\mathbb{R}^d_+)$ such that

$$\mathcal{L}_{h,k}(a)(f) = \sum_{j=1}^{\infty} s_j \langle f, \phi_j \rangle_{L^2_{\beta}(\overline{\mathbb{R}^d_+})} \varphi_j,$$
(3.13)

where $s_j, j \in \mathbb{N}$, are the positive singular values of $\mathcal{L}_{h,k}(a)$ corresponding to ϕ_j . Thus, we get

$$\left\|\mathcal{L}_{h,k}(a)\right\|_{S_1}=\sum_{j=1}^{\infty}s_j=\sum_{j=1}^{\infty}\langle\mathcal{L}_{h,k}(a)(\phi_j),\varphi_j\rangle_{L^2_{\beta}(\overline{\mathbb{R}^d_+})}.$$

Hence, by Fubini's theorem, Cauchy-Schwarz's inequality, Bessel inequality, and relations (2.21) and (2.20) we obtain

$$\begin{split} \left\| \mathcal{L}_{h,k}(a) \right\|_{S_{1}} &= \sum_{j=1}^{\infty} \left\langle \mathcal{L}_{h,k}(a)(\phi_{j}), \varphi_{j} \right\rangle_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \\ &= \sum_{j=1}^{\infty} \frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}} \int_{\Omega_{d+1}} a(b, y) \mathcal{S}_{h}^{W}(\phi_{j})(b, y) \overline{\mathcal{S}_{k}^{W}(\varphi_{j})(b, y)} \, d\mu_{\beta}(b, y) \\ &\leq \frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}} \int_{\Omega_{d+1}} \left| a(b, y) \right| \\ &\quad \times \left(\sum_{j=1}^{\infty} \left| \mathcal{S}_{h}^{W}(\phi_{j})(b, y) \right|^{2} \right)^{\frac{1}{2}} \left(\sum_{j=1}^{\infty} \left| \mathcal{S}_{k}^{W}(\varphi_{j})(b, y) \right|^{2} \right)^{\frac{1}{2}} d\mu_{\beta}(b, y) \\ &\leq \frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}} \int_{\Omega_{d+1}} \left| a(b, y) \right| \|h_{b, y}\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \|k_{b, y}\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} d\mu_{\beta}(b, y) \\ &\leq \frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}} \|a\|_{L^{1}_{\mu_{\beta}}(\Omega_{d+1})}. \end{split}$$

Thus

$$\left\|\mathcal{L}_{h,k}(a)\right\|_{S_1} \leq \frac{1}{\sqrt{\mathcal{C}_h \mathcal{C}_k}} \|a\|_{L^1_{\mu_\beta}(\Omega_{d+1})}.$$

We now prove that $\mathcal{L}_{h,k}(a)$ satisfies the first member of (3.12). Indeed, from (3.13) we have

$$\begin{aligned} \widetilde{a}(b,y) &| = \left| \left\langle \mathcal{L}_{h,k}(a)(h_{b,y}), k_{b,y} \right\rangle_{L^2_{\beta}(\overline{\mathbb{R}^d_+})} \right| \\ &= \left| \sum_{j=1}^{\infty} s_j \langle h_{b,y}, \phi_j \rangle_{L^2_{\beta}(\overline{\mathbb{R}^d_+})} \langle \varphi_j, k_{b,y} \rangle_{L^2_{\beta}(\overline{\mathbb{R}^d_+})} \right| \\ &\leq \frac{1}{2} \sum_{j=1}^{\infty} s_j \left(\left| \langle h_{b,y}, \phi_j \rangle_{L^2_{\beta}(\overline{\mathbb{R}^d_+})} \right|^2 + \left| \langle k_{b,y}, \varphi_j \rangle_{L^2_{\beta}(\overline{\mathbb{R}^d_+})} \right|^2 \right). \end{aligned}$$

Then, using Fubini's theorem, we obtain

$$\begin{split} \int_{\Omega_{d+1}} \left| \widetilde{a}(b,y) \right| d\mu_{\beta}(b,y) &= \frac{1}{2} \sum_{j=1}^{\infty} s_j \left(\int_{\Omega_{d+1}} \left| \langle h_{b,y}, \phi_j \rangle_{L^2_{\beta}(\overline{\mathbb{R}^d_+})} \right|^2 d\mu_{\beta}(b,y) \right. \\ &+ \int_{\Omega_{d+1}} \left| \langle k_{b,y}, \varphi_j \rangle_{L^2_{\beta}(\overline{\mathbb{R}^d_+})} \right|^2 d\mu_{\beta}(b,y) \right) \! . \end{split}$$

Thus, applying Plancherel's identity for \mathcal{S}_h^W and \mathcal{S}_k^W , we get

$$\int_{\Omega_{d+1}} \left| \widetilde{a}(b,y) \right| d\mu_{\beta}(b,y) \leq \frac{\mathcal{C}_h + \mathcal{C}_k}{2} \sum_{j=1}^{\infty} s_j = \frac{\mathcal{C}_h + \mathcal{C}_k}{2} \left\| \mathcal{L}_{h,k}(a) \right\|_{S_1}.$$

The proof is complete.

Corollary 3.1 For a in $L^1_{\mu_{\beta}}(\Omega_{d+1})$, we have the trace formula

$$\operatorname{tr}(\mathcal{L}_{h,k}(a)) = \frac{1}{\sqrt{\mathcal{C}_h \mathcal{C}_k}} \int_{\Omega_{d+1}} a(b, y) \langle k_{b,y}, h_{b,y} \rangle_{L^2_\beta(\overline{\mathbb{R}^d_+})} d\mu_\beta(b, y).$$
(3.14)

Proof By Proposition 3.2, $\mathcal{L}_{h,k}(a) \in S_1$. Then, using (3.2), we get

$$\begin{aligned} \operatorname{tr}(\mathcal{L}_{h,k}(a)) &= \sum_{j=1}^{\infty} \langle \mathcal{L}_{h,k}(a)(\phi_{j}), \phi_{j} \rangle_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \\ &= \sum_{j=1}^{\infty} \frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}} \int_{\Omega_{d+1}} a(b, y) \langle \phi_{j}, h_{b,y} \rangle_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \overline{\langle \phi_{j}, k_{b,y} \rangle_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}} \, d\mu_{\beta}(b, y) \\ &= \frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}} \int_{\Omega_{d+1}} \sum_{j=1}^{\infty} \langle \phi_{j}, h_{b,y} \rangle_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \overline{\langle \phi_{j}, k_{b,y} \rangle_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}} \, d\mu_{\beta}(b, y) \\ &= \frac{1}{\sqrt{\mathcal{C}_{h}\mathcal{C}_{k}}} \int_{\Omega_{d+1}} \langle k_{b,y}, h_{b,y} \rangle_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \, d\mu_{\beta}(b, y), \end{aligned}$$

and the proof is complete.

Corollary 3.2 Let a be in $L^p_{\mu_{\beta}}(\Omega_{d+1})$, $1 \le p \le \infty$. Then, the localization operator

$$\mathcal{L}_{h,k}(a): L^2_{\beta}(\overline{\mathbb{R}^d_+}) \longrightarrow L^2_{\beta}(\overline{\mathbb{R}^d_+})$$

is in S_p , and

$$\left\|\mathcal{L}_{h,k}(a)\right\|_{S_p} \leq \left(\frac{1}{\sqrt{\mathcal{C}_h \mathcal{C}_k}}\right)^{\frac{1}{p}} \|a\|_{L^p_{\mu_\beta}(\Omega_{d+1})}.$$

Proof The result follows from Proposition 3.3 and Theorem 3.2 and by interpolation [17, Theorems 2.10 and 2.11]. $\hfill \Box$

4 Continuous Weinstein wavelet transform and time-frequency concentration

In this section, we suppose that the Weinstein wavelet *h* belongs to $L^1_{\beta}(\overline{\mathbb{R}^d_+}) \cap L^2_{\beta}(\overline{\mathbb{R}^d_+})$.

Proposition 4.1 Let h be a Weinstein wavelet on $\overline{\mathbb{R}^d_+}$ in $L^2_\beta(\overline{\mathbb{R}^d_+})$. Then, $\mathcal{S}^W_h(L^2_\beta(\overline{\mathbb{R}^d_+}))$ is a reproducing kernel Hilbert space with kernel

$$\mathcal{K}_{h}(b',y';b,y) := \frac{1}{\mathcal{C}_{h}} \int_{\mathbb{R}^{4}_{+}} h_{b',y'}(x) \overline{h_{b,y}(x)} \, d\lambda_{\beta}(x).$$

$$(4.1)$$

The kernel is pointwise bounded:

$$\forall (b', y'), (b, y) \in \Omega_{d+1}, \quad \left| \mathcal{K}_h(b', x'; b, y) \right| \le \frac{\|h\|_{L^2_\beta(\overline{\mathbb{R}^d_+})}^2}{C_h}.$$

$$(4.2)$$

Proof Let $f \in L^2_{\beta}(\overline{\mathbb{R}^d_+})$. We have

$$\forall (b,y) \in \Omega_{d+1}, \quad \mathcal{S}_h^W(f)(b,y) = \int_{\mathbb{R}^d_+} f(x) \overline{h_{b,y}(x)} \, d\lambda_\beta(x).$$

Using relation (2.27), we obtain

$$\mathcal{S}_{h}^{W}(f)(b,y) = \frac{1}{\mathcal{C}_{h}} \int_{0}^{\infty} \int_{\mathbb{R}^{d}_{+}} \mathcal{S}_{h}^{W}(f)(b',y') \overline{\mathcal{S}_{h}^{W}(h_{b,y})(b',y')} d\mu_{\beta}(b',y').$$

Applying Proposition 2.2(iii), we find that, for all b, b' > 0 and $y, y' \in \overline{\mathbb{R}^d_+}$, the function

$$y' \mapsto \mathcal{S}_{h}^{W}(h_{b,y})(b',y') = \frac{1}{\mathcal{C}_{h}} \int_{\mathbb{R}^{d}_{+}} h_{b',y'}(x) \overline{h_{b,y}(x)} \, d\lambda_{\beta}(x)$$

belongs to $L^2_\beta(\overline{\mathbb{R}^d_+})$. Therefore, the result is obtained.

In the following we denote:

$$\begin{split} P_h &: L^2_{\mu_{\beta}}(\Omega_{d+1}) \to L^2_{\mu_{\beta}}(\Omega_{d+1}) \text{ is the orthogonal projection from } L^2_{\mu_{\beta}}(\Omega_{d+1}) \text{ onto} \\ \mathcal{S}_h^W(L^2_{\beta}(\overline{\mathbb{R}^d_+})). \\ P_U &: L^2_{\mu_{\beta}}(\Omega_{d+1}) \to L^2_{\mu_{\beta}}(\Omega_{d+1}) \text{ is the orthogonal projection on } L^2_{\mu_{\beta}}(\Omega_{d+1}) \text{ defined by} \end{split}$$

 $\forall v \in L^2_{\mu_{\mathcal{B}}}(\Omega_{d+1}), \quad P_U v = \chi_U v,$

where χ_U denotes the characteristic function of $U \subset \Omega_{d+1}$ with

$$0 < \mu_{\beta}(U) := \int_{U} d\mu_{\beta}(b, y) < \infty.$$

We put

$$\|P_{U}P_{h}\| := \sup \{ \|P_{U}P_{h}\nu\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})} : \nu \in L^{2}_{\mu_{\beta}}(\Omega_{d+1}), \|\nu\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})} = 1 \}.$$

We further prove the concentration of $\mathcal{S}_h^W(f)$ in small sets.

Proposition 4.2 Let h be a Weinstein wavelet, and let $U \subset \Omega_{d+1}$ with

$$\mu_{\beta}(U) < \frac{C_h}{\|h\|_{L^2_{\beta}(\overline{\mathbb{R}^4_+})}^2}.$$

Then, for all $f \in L^2_{\beta}(\overline{\mathbb{R}^d_+})$ *, we have*

$$\left\|\mathcal{S}_{h}^{W}(f) - \chi_{U}\mathcal{S}_{h}^{W}(f)\right\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})} \geq \sqrt{\mathcal{C}_{h}} \sqrt{1 - \frac{\|h\|^{2}_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}{\mathcal{C}_{h}}\mu_{\beta}(U)\|f\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})},\tag{4.3}$$

where χ_U denotes the characteristic function of U.

Proof From Plancherel's formula (Theorem 2.1) we have

$$\mathcal{C}_{h} \|f\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}^{2} = \|\mathcal{S}_{h}^{W}(f)\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})}^{2} = \|\mathcal{S}_{h}^{W}(f)\|_{L^{2}_{\mu_{\beta}}(U)}^{2} + \|\mathcal{S}_{h}^{W}(f)\|_{L^{2}_{\mu_{\beta}}(U^{c})}^{2}.$$
(4.4)

On the other hand, from relation (2.24) we have

$$\int_{U} \left| \mathcal{S}_{h}^{W}(f)(b,y) \right|^{2} d\mu_{\beta}(b,y) \leq \left\| \mathcal{S}_{h}^{W}(f) \right\|_{L^{\infty}_{\mu_{\beta}}(\Omega_{d+1})}^{2} \mu_{\beta}(U) \\ \leq \mu_{\beta}(U) \left\| f \right\|_{L^{2}_{\beta}(\mathbb{R}^{d}_{+})}^{2} \left\| h \right\|_{L^{2}_{\beta}(\mathbb{R}^{d}_{+})}^{2}.$$
(4.5)

Thus, using (4.4) and (4.5), we deduce the result.

Remark 4.1 If $\mathcal{S}_{h}^{W}(f)$ is supported in U and $\mu_{\beta}(U) < \frac{C_{h}}{\|h\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{4}_{+}})}^{2}}$, then f = 0.

We further prove the concentration of $S_h^W(f)$ in arbitrary sets of finite measures.

Theorem 4.1 Let h be a Weinstein wavelet, and let $U \subset \Omega_{d+1}$ with $0 < \mu_{\beta}(U) < \infty$. If $P_h(L^2_{\mu_{\beta}}(\Omega_{d+1})) \cap P_U(L^2_{\mu_{\beta}}(\Omega_{d+1})) = \{0\}$, then there exists a constant C := C(h, U) > 0 such that, for all $f \in L^2_{\beta}(\mathbb{R}^d_+)$,

$$\left\|\mathcal{S}_{h}^{W}(f) - \chi_{U}\mathcal{S}_{h}^{W}(f)\right\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})} \ge C\|f\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}.$$
(4.6)

For the proof of this theorem, we need the following lemma.

Lemma 4.1 ([20]) Let \mathcal{H}_1 and \mathcal{H}_2 be two closed subspaces of a Hilbert space \mathcal{H} such that $\mathcal{H}_1 \cap \mathcal{H}_2 = \{0\}$. Let $P_{\mathcal{H}_1}$ and $P_{\mathcal{H}_2}$ denote the corresponding orthogonal projections, and assume that the product $P_{\mathcal{H}_1}P_{\mathcal{H}_2}$ is a compact operator. Then, there exists a constant C > 0 such that, for $f \in \mathcal{H}$,

$$\|P_{\mathcal{H}_{1}^{\perp}}f\|_{\mathcal{H}} + \|P_{\mathcal{H}_{2}^{\perp}}f\|_{\mathcal{H}} \ge C\|f\|_{\mathcal{H}}.$$
(4.7)

Proof of Theorem 4.1 Define \mathcal{H}_1 and \mathcal{H}_2 by

$$\mathcal{H}_1 := P_U \left(L^2_{\mu_\beta}(\Omega_{d+1}) \right), \qquad \mathcal{H}_2 := P_h \left(L^2_{\mu_\beta}(\Omega_{d+1}) \right).$$

Proceeding as in [21], we prove that

$$\|P_{U}P_{h}\|_{HS} := \left(\int_{\Omega_{d+1} \times \Omega_{d+1}} |\chi_{U}(b,y)|^{2} |\mathcal{K}_{h}(b',y';b,y)|^{2} d\mu_{\beta}(b',y') d\mu_{\beta}(b,y) \right)^{\frac{1}{2}} \\ \leq \frac{\|h\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}}{\sqrt{\mathcal{C}_{h}}} \sqrt{\mu_{\beta}(U)} < \infty.$$
(4.8)

Hence, $P_{U}P_{h}$ is a Hilbert-Schmidt operator and therefore compact. Now, Lemma 4.1 implies the existence of a constant C > 0 such that (4.7) holds for $P_{\mathcal{H}_1} := P_U$ and $P_{\mathcal{H}_2} := P_h$. Since

$$P_{\mathcal{H}_{2}^{\perp}}(\mathcal{S}_{h}^{W}(f)) = (Id - P_{h})\mathcal{S}_{h}^{W}(f) = 0,$$

this leads to (4.6).

Definition 4.1 Let *h* be a Weinstein wavelet, and let $U \subset \Omega_{d+1}$ such that $0 < \mu_{\beta}(U) < \infty$. Then

- (1) We say that *U* is weakly annihilating if any function $f \in L^2_\beta(\overline{\mathbb{R}^d_+})$ vanishes when its Weinstein wavelet transform $S_h^W(f)$ with respect to the Weinstein wavelet *h* is supported in *U*.
- (2) We say that *U* is strongly annihilating if there exists a constant $C_{\beta}(U) > 0$ such that, for every function $f \in L^2_{\beta}(\mathbb{R}^d_+)$,

$$C_{\beta}(U) \left\| \mathcal{S}_{h}^{W}(f) - \chi_{U} \mathcal{S}_{h}^{W}(f) \right\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})} \ge \left\| f \right\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \left\| h \right\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}.$$
(4.9)

The constant $C_{\beta}(U)$ is called the annihilation constant of U.

The analogue of this definition was introduced by Ghobber and Omri [22] in the cadre of windowed Hankel transform.

Remark 4.2

- (1) It is clear that every strongly annihilating set is also a weakly.
- (1) It is clear that every strongly annumating set to use L(2) From Proposition 4.2 we see that any set $U \subset \Omega_{d+1}$ with $\mu_{\beta}(U) < \frac{C_h}{\|h\|_{L^2_{\beta}(\overline{\mathbb{R}^d_+})}^2}$ is strongly

annihilating.

- (3) As the operator $P_U P_h$ is Hilbert-Schmidt and hence compact, from [23] we have that if U is weakly annihilating, then it is also strongly annihilating.
- (4) If $||P_UP_h|| < 1$, then, for all $f \in L^2_\beta(\overline{\mathbb{R}^d_+})$,

$$\frac{1}{\sqrt{1 - \|P_{U}P_{h}\|^{2}}} \|\chi_{U^{c}}\mathcal{S}_{h}^{W}(f)\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})} \ge \sqrt{\mathcal{C}_{h}} \|f\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})} \|h\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}.$$
(4.10)

(5) Following [23, p.88], we have that if *U* is strongly annihilating, then $||P_UP_h|| < 1$.

In the following, we prove the Benedicks uncertainty principle for the Weinstein wavelet transform under some condition on the Weinstein wavelet. We note that the Benedickstype uncertainty principle for the windowed Hankel transform was studied by Ghobber and Omri [22].

Theorem 4.2 Let h be a Weinstein wavelet such that

$$\int_{\{\xi \in \mathbb{R}^{\vec{d}}_{+}: \mathcal{F}_{W}(\bar{h})(\xi) \neq 0\}} d\lambda_{\beta}(\xi) < \infty.$$
(4.11)

For any subset $U \subset \Omega_{d+1}$ *such that*

$$\int_{\mathbb{R}^d_+} \chi_U(b,y) \, d\lambda_\beta(y) < \infty$$

for almost every b > 0, we have

$$P_h(L^2_{\mu_\beta}(\Omega_{d+1})) \cap P_U(L^2_{\mu_\beta}(\Omega_{d+1})) = \{0\}.$$
(4.12)

Proof Let $F \in P_h(L^2_{\mu_\beta}(\Omega_{d+1})) \cap P_U(L^2_{\mu_\beta}(\Omega_{d+1})) \setminus \{0\}$. Then by the definition of the operators P_U and P_h there exists a function $f \in L^2_\beta(\overline{\mathbb{R}^d_+})$ such that $F = \mathcal{S}_h^W(f)$ and $\operatorname{supp} F \subset U$.

Let b > 0 be such that $\int_{\mathbb{R}^d_+} \chi_U(b, y) d\lambda_\beta(y) < \infty$. Consider the function F_b defined by

$$F_b(y) = \mathcal{S}_h^W(f)(b, y), \quad y \in \overline{\mathbb{R}_+^d}.$$

Then we have

$$\operatorname{supp} F_b \subset \left\{ y \in \overline{\mathbb{R}^d_+} : (b, y) \in U \right\}$$

and

$$\int_{\operatorname{supp} F_b} d\lambda_\beta(\xi) < \infty.$$

On the other hand, using (2.23) and hypothesis (4.11), we get

$$\int_{\{\xi\in\overline{\mathbb{R}^d_+}:\mathcal{F}_W(F_b)(\xi)\neq 0\}} d\lambda_\beta(\xi) < \infty.$$

Using Proposition 2.3, we obtain that $F_b = 0$ for every b > 0, and hence F = 0.

Consequently, we obtain the following improvement.

Corollary 4.1 Let h be a Weinstein wavelet on $\overline{\mathbb{R}^d_+}$ in $L^2_\beta(\overline{\mathbb{R}^d_+})$ such that

$$\int_{\{\xi\in\overline{\mathbb{R}^d_+}:\mathcal{F}_W(\bar{h})(\xi)\neq 0\}}d\lambda_\beta(\xi)<\infty.$$

Then, for any subset $U \subset \Omega_{d+1}$ such that $0 < \mu_{\beta}(U) < \infty$, there exists a constant C := C(h, U) > 0 such that, for all $f \in L^2_{\beta}(\overline{\mathbb{R}^d_+})$, we have

$$\|\mathcal{S}_{h}^{W}(f) - \chi_{U}\mathcal{S}_{h}^{W}(f)\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})} \ge C\|f\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}.$$
(4.13)

Now we will derive a sufficient condition by means of which we can recover a signal F belonging to $L^2_{\mu\beta}(\Omega_{d+1})$ from the knowledge of its truncated version following the Donoho-Stark criterion [24].

Let *h* be a Weinstein wavelet function. A signal $F \in L^2_{\mu_{\beta}}(\Omega_{d+1})$ is transmitted to a receiver who knows that $F \in S_h^W(L^2_{\beta}(\overline{\mathbb{R}^d_+}))$. Suppose that the observation of F is corrupted by a noise $n \in L^2_{\mu_{\beta}}(\Omega_{d+1})$ (which is nonetheless assumed to be small) and unregistered values on $U \in \Omega_{d+1}$. Thus, the observable function *r* satisfies

$$r(b,y) = \begin{cases} F(b,y) + n(b,y) & \text{if } (b,y) \in U^c, \\ 0 & \text{if } (b,y) \in U. \end{cases}$$
(4.14)

Here we have assumed without loss of generality that n = 0 on U. Equivalently,

$$r = (Id - P_U)F + n. (4.15)$$

We say that F can be stably reconstructed from r if there exist a linear operator

 $L_{U,h}: L^2_{\mu_{\beta}}(\Omega_{d+1}) \to L^2_{\mu_{\beta}}(\Omega_{d+1})$

and a constant C(U, h) such that

$$\|F - L_{U,h}(r)\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})} \le C(U,h) \|n\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})}.$$
(4.16)

Proceeding as in [24], it is easy to prove the following:

Proposition 4.3 Let h be a Weinstein wavelet function such that

$$\int_{\{\xi\in\overline{\mathbb{R}^d_+}:\mathcal{F}_W(\bar{h})(\xi)\neq 0\}} d\lambda_\beta(\xi) < \infty.$$

Let $U \subset \Omega_{d+1}$ with $0 < \mu_{\beta}(U) < \infty$. Then F can be stably reconstructed from r. The constant C(U,h) in (4.16) is not larger than $(1 - \|P_UP_h\|)^{-1}$.

The identity

$$Q = (Id - P_{U}P_{h})^{-1} = \sum_{j=0}^{\infty} (P_{U}P_{h})^{j}$$

suggests an algorithm for computing Q(r). Finally, using a method similar to that in [24], we give an algorithm for computing $L_{U,h}$. Indeed, put

$$F_k = \sum_{j=0}^k (P_U P_h)^j r.$$

Then
$$F_k \to Q(r)$$
 as $k \to \infty$. Now

$$F_0 = r,$$

$$F_1 = r + P_U P_h F_0,$$

$$F_2 = r + P_U P_h F_1,$$

$$\dots$$
(4.17)

and so on. The iteration converges at a geometric rate to the fixed point

$$F = r + P_U P_h F.$$

Algorithms of type (4.17), have been applied to a host of problems in signal recovery; see [24] and others.

5 Heisenberg-type uncertainty inequalities for the Weinstein wavelet transform

In this section, we establish many Heisenberg-type uncertainty inequalities for the Weinstein wavelet transform.

Proposition 5.1 We assume that the Weinstein wavelet h satisfies $\|h\|_{L^2_{\theta}(\mathbb{R}^d_+)} = 1$.

- *Let s* > 0*. Then we have the following uncertainty inequalities.*
- (1) There exists a constant $C_1(\beta, s) > 0$ such that, for all $f \in L^2_{\beta}(\overline{\mathbb{R}^d_+})$,

$$\left\| \left\| \left(\frac{1}{b}, y\right) \right\|^{s} \mathcal{S}_{h}^{W}(f) \right\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})} \geq C_{1}(\beta, s) \left\| f \right\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}.$$
(5.1)

(2) There exists a constant $C_2(\beta, s) > 0$ such that, for all $f \in L^2_{\beta}(\overline{\mathbb{R}^d_+})$,

$$\|\|y\|^{s} \mathcal{S}_{h}^{W}(f)\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})} \||b|^{-s} \mathcal{S}_{h}^{W}(f)\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})} \ge C_{2}(\beta, s) \|f\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}^{2}.$$
(5.2)

(3) There exists a constant $C_3(\beta, s) > 0$ such that, for all $f \in L^2_\beta(\overline{\mathbb{R}^d_+})$ and every measurable subset U of finite measure $0 < \mu_\beta(U) < \infty$,

$$\left\|\mathcal{S}_{h}^{W}(f)\right\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})} \leq C_{3}(\beta,s)\sqrt{\mu_{\beta}(U)}\left\|\left\|\left(\frac{1}{b},y\right)\right\|^{s}\mathcal{S}_{h}^{W}(f)\right\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})}.$$
(5.3)

Proof (1) Let $r \in (0,1]$ be such that $\mu_{\beta}(U_r) < C_h$, where

$$U_r = \left\{ (b, y) \in \Omega_{d+1} : \left\| \left(\frac{1}{b}, y \right) \right\| < r \right\}.$$

From Proposition 4.2 we have

$$\begin{split} \|f\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}^{2} &\leq \frac{1}{\mathcal{C}_{h} - \mu_{\beta}(U_{r})} \int_{U_{r}^{c}} \left|\mathcal{S}_{h}^{W}(f)(b, y)\right|^{2} d\mu_{\beta}(b, y) \\ &\leq \frac{1}{r^{2s}(\mathcal{C}_{h} - \mu_{\beta}(U_{r}))} \int_{\|(\frac{1}{b}, y)\| \geq r} \left\|\left(\frac{1}{b}, y\right)\right\|^{2s} |\mathcal{S}_{h}^{W}(f)(b, y)|^{2} d\mu_{\beta}(b, y) \\ &\leq \frac{1}{r^{2s}(\mathcal{C}_{h} - \mu_{\beta}(U_{r}))} \left\|\left\|\left(\frac{1}{b}, y\right)\right\|^{s} \mathcal{S}_{h}^{W}(f)\right\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})}^{2}. \end{split}$$

Thus we have obtained the result with $C_1(\beta, s) := r^s \sqrt{C_h - \mu_\beta(U_r)}$.

Now we will prove (2). Indeed, using the fact that $\|(\frac{1}{b}, y)\|^s \le 2^s (|b|^{-s} + \|y\|^s)$ in (5.1), we get

$$\left\| \|y\|^{s} \mathcal{S}_{h}^{W}(f) \right\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})}^{2} + \left\| |b|^{-s} \mathcal{S}_{h}^{W}(f) \right\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})}^{2} \geq \frac{C_{1}^{2}(\beta, s)}{2^{2s}} \|f\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}^{2}$$

Replacing *f* by $\delta_t f := \frac{1}{t^{\beta+\frac{d+1}{2}}} f(\frac{\cdot}{t})$ in the previous inequality, by (2.25) and by a suitable change of variables we obtain:

$$t^{2s} \| \|y\|^{s} \mathcal{S}_{h}^{W}(f) \|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})}^{2} + t^{-2s} \| |b|^{-s} \mathcal{S}_{h}^{W}(f) \|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})}^{2} \geq \frac{C_{1}^{2}(\beta, s)}{2^{2s}} \|f\|_{L^{2}_{\beta}(\overline{\mathbb{R}_{+}^{d}})}^{2}$$

Then (5.2) follows by minimizing the left-hand side of that inequality over t > 0. Now we will prove (3). Indeed, using the estimates

$$\left\|\mathcal{S}_{h}^{W}(f)\right\|_{L^{2}_{\mu_{\beta}}(U_{r})} \leq \sqrt{\mu_{\beta}(U)} \left\|\mathcal{S}_{h}^{W}(f)\right\|_{L^{\infty}_{\mu_{\beta}}(\Omega_{d+1})}$$

and

$$\left\|\mathcal{S}_{h}^{W}(f)\right\|_{L^{\infty}_{\mu_{\beta}}(\Omega_{d+1})} \leq \left\|f\right\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})},$$

we get

$$\left\|\mathcal{S}_{h}^{W}(f)\right\|_{L^{2}_{\mu_{\beta}}(U_{r})} \leq \sqrt{\mu_{\beta}(U)} \|f\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}.$$

Finally, the result immediately follows from (5.1).

Now we recall the uncertainty principle of Heisenberg type for the Weinstein transform.

Proposition 5.2 ([25]) For s, t > 0, there exists a positive constant $C(\beta, s, t)$ such that, for every $f \in L^2_\beta(\overline{\mathbb{R}^d_+})$, we have the inequality

$$\left\| \|\xi\|^{s} \mathcal{F}_{W}(f) \right\|_{L^{2}_{\beta}(\mathbb{R}^{d}_{+})}^{\frac{t}{s+t}} \left\| \|y\|^{t} f \right\|_{L^{2}_{\beta}(\mathbb{R}^{d}_{+})}^{\frac{s}{s+t}} \ge C(\beta, s, t) \|f\|_{L^{2}_{\beta}(\mathbb{R}^{d}_{+})}.$$
(5.4)

Theorem 5.1 Let s, t > 0, and let h be a Weinstein wavelet on $\overline{\mathbb{R}^d_+}$ in $L^2_\beta(\overline{\mathbb{R}^d_+})$. Then, for all f in $L^2_\beta(\overline{\mathbb{R}^d_+})$, we have the inequality

$$\|\|y\|^{s} \mathcal{S}_{h}^{W}(f)\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})}^{\frac{t}{t+s}} \|\|\xi\|^{t} \mathcal{F}_{W}(f)\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}^{\frac{s}{s+t}} \geq C(\beta, s, t) \mathcal{C}_{h}^{\frac{t}{2(s+t)}} \|f\|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}.$$
(5.5)

Proof Let us assume the nontrivial case that both integrals on the left-hand side of (5.5) are finite. We get from the admissibility condition (2.15) for *h* that

$$\int_0^\infty \int_{\mathbb{R}^d_+} \|\xi\|^{2t} \big| \mathcal{F}_W(h)(b\xi) \big|^2 \big| \mathcal{F}_W(f)(\xi) \big|^2 \frac{d\lambda_\beta(\xi) \, da}{a} = \mathcal{C}_h \int_{\mathbb{R}^d_+} \|\xi\|^{2t} \big| \mathcal{F}_W(f)(\xi) \big|^2 \, d\lambda_\beta(\xi).$$

Using relation (2.23), we obtain

$$\int_{0}^{\infty} \int_{\mathbb{R}^{d}_{+}} \|\xi\|^{2t} \left| \mathcal{F}_{W} \left(\mathcal{S}_{h}^{W}(f)(b, \cdot) \right)(\xi) \right|^{2} d\mu_{\beta}(b, \xi) = \mathcal{C}_{h} \int_{\mathbb{R}^{d}_{+}} \|\xi\|^{2t} \left| \mathcal{F}_{W}(f)(\xi) \right|^{2} d\lambda_{\beta}(\xi).$$
(5.6)

Moreover, using (5.4), we get

$$\begin{aligned} \forall b > 0, \quad \left(\int_{\mathbb{R}^d_+} \|\xi\|^{2t} \left| \mathcal{F}_W \left(\mathcal{S}_h^W(f)(b, \cdot) \right)(\xi) \right|^2 d\lambda_\beta(\xi) \right)^{\frac{s}{s+t}} \\ & \times \left(\int_{\mathbb{R}^d_+} \|y\|^{2s} \left| \mathcal{S}_h^W(f)(b,y) \right|^2 d\lambda_\beta(y) \right)^{\frac{t}{s+t}} \\ & \ge C^2(\beta, s, t) \int_{\mathbb{R}^d_+} \left| \mathcal{S}_h^W(f)(b,y) \right|^2 d\lambda_\beta(y). \end{aligned}$$

Integrating with respect to $\frac{db}{b^{2\beta+d+2}}$, we obtain

$$\begin{split} &\int_0^\infty \biggl[\left(\int_{\mathbb{R}^d_+} \|\xi\|^{2t} \big| \mathcal{F}_W \big(\mathcal{S}_h^W(f)(b, \cdot) \big)(\xi) \big|^2 \, d\lambda_\beta(\xi) \right)^{\frac{s}{s+t}} \\ & \times \left(\int_{\mathbb{R}^d_+} \|y\|^{2s} \big| \mathcal{S}_h^W(f)(b, y) \big|^2 \, d\lambda_\beta(y) \right)^{\frac{t}{s+t}} \biggr] \frac{db}{b^{2\beta+d+2}} \\ & \ge C^2(\beta, s, t) \int_0^\infty \int_{\mathbb{R}^d_+} \big| \mathcal{S}_h^W(f)(b, y) \big|^2 \, d\lambda_\beta(y) \frac{db}{b^{2\beta+d+2}}. \end{split}$$

The left-hand side of this inequality may be estimated from above using Hölder's inequality. The right-hand side can be rewritten by Plancherel's formula for S_h^W . Therefore, from (5.6) we get

$$\begin{split} \left(\int_0^\infty \int_{\overline{\mathbb{R}^d_+}} \|\xi\|^{2t} \big| \mathcal{F}_W \big(\mathcal{S}_h^W(f)(b, \cdot) \big)(\xi) \big|^2 \, d\mu_\beta(b, \xi) \Big)^{\frac{s}{s+t}} \\ & \times \left(\int_0^\infty \int_{\overline{\mathbb{R}^d_+}} \|y\|^{2s} \big| \mathcal{S}_h^W(f)(b, y) \big|^2 \, d\mu_\beta(b, y) \Big)^{\frac{t}{s+t}} \\ &= \mathcal{C}_h^{\frac{s}{s+t}} \left(\int_{\overline{\mathbb{R}^d_+}} \|\xi\|^{2t} \big| \mathcal{F}_W \big(f(\xi)\big) \big|^2 \, d\lambda_\beta(\xi) \Big)^{\frac{s}{s+t}} \end{split}$$

$$\times \left(\int_0^\infty \int_{\overline{\mathbb{R}^d_+}} \|y\|^{2s} \left| \mathcal{S}_h^W(f)(b,y) \right|^2 d\mu_\beta(b,y) \right)^{\frac{1}{s+t}}$$

$$\ge C^2(\beta,s,t) \mathcal{C}_h \|f\|_{L^2_\beta(\overline{\mathbb{R}^d_+})}^2.$$

This proves the result.

6 Heisenberg-type uncertainty inequalities for the modified Weinstein wavelet transform

There is no uncertainty principle of Heisenberg type for $S_h^W(f)(b, y)$ with respect to *b*.

For this, we consider the Weinstein wavelet h defined by relation (2.16) and introduce the modified Weinstein continuous wavelet transform $\widetilde{\mathcal{S}_h^W}$ given by

$$\widetilde{\mathcal{S}_{h}^{W}}f(b,y) = b^{2\beta+d+1}\mathcal{F}_{W}^{-1}\left(\mathcal{F}_{W}f(b\xi)\overline{\mathcal{F}_{W}(h)(\xi)}\right)(y), \quad (b,y) \in \Omega_{d+1}.$$
(6.1)

Using this transform, we obtain the following theorem.

Theorem 6.1 For s, t > 0 and every f in $L^2_{\beta}(\overline{\mathbb{R}^d_+})$, we have

$$\|b^{t}\widetilde{\mathcal{S}_{h}^{W}}f\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})}^{\frac{s}{s+t}}\|\|y\|^{s}f\|_{L^{2}_{\beta}(\overline{\mathbb{R}_{+}^{d}})}^{\frac{t}{s+t}} \ge C(\beta,s,t)\big(\mathcal{M}\big(e^{-y^{2}}\big)(2t-4)\big)^{\frac{s}{2(s+t)}}\|f\|_{L^{2}_{\beta}(\overline{\mathbb{R}_{+}^{d}})},\tag{6.2}$$

where

$$\mathcal{M}: H \mapsto \mathcal{M}(H)(y) = \int_0^\infty H(r) \frac{dr}{r^{y+1}}$$

is the classical Mellin transform.

Proof In the following, we assume that

$$\int_0^\infty \int_{\mathbb{R}^d_+} b^{2t} \left| \widetilde{\mathcal{S}_h^W} f(b,y) \right|^2 d\mu_\beta(b,y) < \infty \quad \text{and} \quad \int_{\mathbb{R}^d_+} \|y\|^{2s} \left| f(y) \right|^2 d\lambda_\beta(y) < \infty.$$

Otherwise, (6.2) is trivially satisfied. Using Fubini's theorem and (2.26), we have

$$\begin{split} &\int_{\Omega_{d+1}} b^{2t} \left| \left(\widetilde{\mathcal{S}_h^{W}} f \right)(b, y) \right|^2 d\mu_{\beta}(b, y) \\ &= \int_0^\infty b^{2t} \int_{\mathbb{R}^d_+} b^{2\beta+d+1} \left| \mathcal{F}_W f(b\xi) \right|^2 \left| \mathcal{F}_W(h)(\xi) \right|^2 d\lambda_{\beta}(\xi) \frac{db}{b} \\ &= \int_{\mathbb{R}^d_+} \int_0^\infty b^{2t} \left| \mathcal{F}_W f(\xi) \right|^2 \left| \mathcal{F}_W(h) \left(\frac{\lambda}{b} \right) \right|^2 \frac{db}{b} d\lambda_{\beta}(\xi) \\ &= \int_{\mathbb{R}^d_+} \Lambda(\xi) \left| \mathcal{F}_W(f)(\xi) \right|^2 d\lambda_{\beta}(\xi) \end{split}$$

with

$$\Lambda(\xi) = \int_0^\infty b^{2t} \left| \mathcal{F}_W(h) \left(\frac{\lambda}{b} \right) \right|^2 \frac{db}{b}.$$

Using (2.16), by simple calculations we get

$$\begin{split} \Lambda(\xi) &= \int_0^\infty b^{2t} \left| \mathcal{F}_W(h) \left(\frac{\lambda}{b} \right) \right|^2 \frac{db}{b} \\ &= \left(\int_0^\infty e^{-y^2} \frac{dy}{y^{2t-3}} \right) \|\xi\|^{2t} \\ &= \left(\mathcal{M}(e^{-y^2})(2t-4) \right) \|\xi\|^{2t}. \end{split}$$

Thus

$$\begin{split} &\left(\int_{0}^{\infty}\int_{\mathbb{R}^{d}_{+}}b^{2t}\big|\widetilde{\mathcal{S}_{h}^{W}}f(b,y)\big|^{2}\,d\mu_{\beta}(b,y)\right)^{\frac{s}{s+t}}\left(\int_{\mathbb{R}^{d}_{+}}\|y\|^{2s}\big|f(y)\big|^{2}\,d\lambda_{\beta}(y)\right)^{\frac{t}{s+t}}\\ &=\left(\mathcal{M}\big(e^{-y^{2}}\big)(2t-4)\big)^{\frac{s}{(s+t)}}\\ &\times\left(\int_{\mathbb{R}^{d}_{+}}\|\xi\|^{2t}\big|\mathcal{F}_{W}(f)(\xi)\big|^{2}\,d\lambda_{\beta}(\xi)\right)^{\frac{s}{s+t}}\left(\int_{\mathbb{R}^{d}_{+}}\|y\|^{2s}\big|f(y)\big|^{2}\,d\lambda_{\beta}(y)\right)^{\frac{t}{s+t}}.\end{split}$$

Now, the result is obtained from Proposition 5.2.

Remark 6.1

- (i) It is easy to see that $\mathcal{M}(e^{-y^2})(2t-4) < \infty$ if and only if t < 2.
- (ii) In the classical setting, the notion of the modified continuous wavelet transform was first introduced by Wilczok [20].

Corollary 6.1 For s, t > 0 and f in $L^2_{\beta}(\overline{\mathbb{R}^d_+})$, we have

$$\begin{split} \| \|y\|^{s} \mathcal{S}_{h}^{W} f \|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})}^{\frac{1}{s+t}} \| b^{t} \widetilde{\mathcal{S}_{h}^{W}} f \|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})}^{\frac{1}{s+t}} \\ &\geq C(\beta, s, t) \mathcal{C}_{h}^{\frac{t}{2(s+t)}} \left(\mathcal{M}(e^{-y^{2}})(2t-4) \right)^{\frac{s}{2(s+t)}} \| f \|_{L^{2}_{\beta}(\mathbb{R}^{d}_{+})}. \end{split}$$

Proof By the preceding we have

$$\begin{split} &\left(\int_{0}^{\infty}\int_{\mathbb{R}^{d}_{+}}b^{2t}\big|\widetilde{\mathcal{S}_{h}^{W}}f(b,y)\big|^{2}\,d\mu_{\beta}(b,y)\right)^{\frac{s}{s+t}}\left(\int_{0}^{\infty}\int_{\mathbb{R}^{d}_{+}}\|y\|^{2s}\big|\mathcal{S}_{h}^{W}f(b,y)\big|^{2}\,d\mu_{\beta}(b,y)\right)^{\frac{t}{s+t}} \\ &=\left(\mathcal{M}\big(e^{-y^{2}}\big)(2t-4)\big)^{\frac{s}{(s+t)}}\left(\int_{\mathbb{R}^{d}_{+}}\|\xi\|^{2t}\big|\mathcal{F}_{W}(f)(\xi)\big|^{2}\,d\lambda_{\beta}(\xi)\right)^{\frac{s}{s+t}} \\ &\times\left(\int_{0}^{\infty}\int_{\mathbb{R}^{d}_{+}}\|y\|^{2s}\big|\mathcal{S}_{h}^{W}f(b,y)\big|^{2}\,d\mu_{\beta}(b,y)\right)^{\frac{t}{s+t}}. \end{split}$$

The result follows from relation (5.5).

Corollary 6.2 For s, t > 0 and for any f in $L^2_{\beta}(\overline{\mathbb{R}^d_+})$, we have

$$\begin{split} \| \| (b,y) \|^{s} \mathcal{S}_{h}^{W} f \|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})}^{\frac{t}{s+t}} \| \| (b,y) \|^{t} \widetilde{\mathcal{S}_{h}^{W}} f \|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})}^{\frac{s}{s+t}} \\ &\geq C(\beta,s,t) \mathcal{C}_{h}^{\frac{s}{2(s+t)}} \left(\mathcal{M} \left(e^{-y^{2}} \right) (2t-4) \right)^{\frac{s}{2(s+t)}} \| f \|_{L^{2}_{\beta}(\overline{\mathbb{R}^{d}_{+}})}. \end{split}$$

Proof The result follows from Corollary 6.1 and the fact that

$$||(b,y)||^{2s} \ge ||y||^{2s}$$
 and $||(b,y)||^{2t} \ge b^{2t}$.

In the following, we give the local-type uncertainty principle.

Corollary 6.3 For s > 0, subsets $U \in \Omega_{d+1}$ such that $0 < \mu_{\beta}(U) < \infty$, and f in $L^2_{\beta}(\mathbb{R}^d_+)$, we have

$$\begin{split} &\int_{\Omega_{d+1}} \chi_{\mathcal{U}}(b,y) \left| \mathcal{S}_{h}^{W} f(b,y) \right|^{2} d\mu_{\beta}(b,y) \\ &\leq C_{\beta}(s) \left\| \left\| (b,y) \right\|^{s} \mathcal{S}_{h}^{W} f \right\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})}^{\frac{1}{2}} \left\| \left\| (b,y) \right\|^{s} \widetilde{\mathcal{S}_{h}^{W}} f \right\|_{L^{2}_{\mu_{\beta}}(\Omega_{d+1})}^{\frac{1}{2}}, \end{split}$$

where

$$C_{\beta}(s) \coloneqq \frac{\mu_{\beta}(U) \|h\|_{L^2_{\beta}(\overline{\mathbb{R}^d_+})}^2}{C(\beta, s, s) \sqrt{\mathcal{C}_h(\mathcal{M}(e^{-y^2})(2t-4))}}$$

Proof From relation (4.5), we have

$$\int_0^\infty \int_{\mathbb{R}^d_+} \chi_U(b,y) \left| \mathcal{S}_h^W f(b,y) \right|^2 d\mu_\beta(b,y) \le \mu_\beta(U) \|h\|_{L^2_\beta(\overline{\mathbb{R}^d_+})}^2 \|f\|_{L^2_\beta(\overline{\mathbb{R}^d_+})}^2.$$

On the other hand, from Corollary 6.2 we have

$$\|f\|_{L^2_{\beta}(\overline{\mathbb{R}^d_+})}^2 \leq \frac{(\int_{\Omega_{d+1}} \|(b,y)\|^{2s} |\widetilde{\mathcal{S}^W_h}f(b,y)|^2 \, d\mu_{\beta}(b,y))^{\frac{1}{2}} (\int_{\Omega_{d+1}} \|(b,y)\|^{2s} |\mathcal{S}^W_hf(b,y)|^2 \, d\mu_{\beta}(b,y))^{\frac{1}{2}}}{C^2(s,s)\sqrt{\mathcal{C}_h(\mathcal{M}(e^{-y^2})(2t-4))}}.$$

The result immediately follows.

7 Conclusions

This paper is devoted to developing the localization operator theory for the continuous Weinstein wavelet transform. Given a symbol *a* and two Weinstein wavelets *h*, *k*, we investigate the multilinear mapping from $(a, h, k) \in L^p_{\mu_\beta}(\Omega_{d+1}) \times L^2_\beta(\overline{\mathbb{R}^d_+}) \times L^2_\beta(\overline{\mathbb{R}^d_+}), 1 \le p \le \infty$, to the localization operator $\mathcal{L}_{h,k}(a)$, and we give sufficient conditions for $\mathcal{L}_{h,k}(a)$ to be bounded or to belong to a Schatten class. Our results are formulated in terms of time-frequency analysis.

The uncertainty principles in the context of the continuous Weinstein wavelet transform are also established.

Acknowledgements

The authors are deeply indebted to the referees for providing constructive comments and helps in improving the contents of this article. The first author thanks the emeritus Professor Khalifa Trimèche for his help.

Competing interests

The authors declare that they have no competing interests regarding the present manuscript.

Authors' contributions

Both authors worked in coordination. Both authors carried out the proof and read and approved the final version of the manuscript.

Author details

¹College of Sciences, Department of Mathematics, Taibah University, PO BOX 30002, Al Madinah Al Munawarah, Saudi Arabia. ²College of Sciences, Department of Mathematics, King Khalid University, Mohayil, Saudi Arabia.

Publisher's Note

Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

Received: 25 April 2017 Accepted: 6 October 2017 Published online: 28 October 2017

References

- 1. Guliyev, V, Hassanov, J: Sobolev-Morrey type inequality for Riesz potentials, associated with the Laplace-Bessel differential operator. Fract. Calc. Appl. Anal. 9(1), 17-32 (2006)
- 2. Brelot, M: Equation de Weinstein et potentiels de Marcel Riesz. In: Séminaire de Théorie de Potentiel Paris, No. 3. Lecture Notes in Math., vol. 681, pp. 18-38. Springer, Berlin (1978)
- 3. Ben Nahia, Z, Ben Salem, N: Spherical harmonics and applications associated with the Weinstein operator. In: Potential Theory - ICPT 94. De Gruyter Proceedings in Mathematics, pp. 235-241. de Gruyter, Berlin (1996)
- 4. Ben Nahia, Z, Ben Salem, N: On a mean value property associated with the Weinstein operator. In: Potential Theory ICPT 94. De Gruyter Proceedings in Mathematics, pp. 243-253. de Gruyter, Berlin (1996)
- Grossmann, A, Morlet, J: Decomposition of Hardy functions into square integrable wavelets of constant shape. SIAM J. Math. Anal. 15(4), 723-736 (1984). doi:10.1137/0515056
- 6. Chui, CK: An Introduction to Wavelets. Academic Press, Waltham (1992)
- 7. Koornwinder, TH: The continuous wavelet transform. In: Wavelets: An Elementary Treatment of Theory and Applications, pp. 27-48. World Scientific, Singapore (1993)
- 8. Meyer, Y: Wavelets and Operators. Cambridge University Press, Cambridge (1995)
- 9. Daubechies, I: Ten Lectures on Wavelets. CBMS-NSF Regional Conference Series in Applied Mathematics, vol. 61. SIAM, Philadelphia (1992)
- Goupilland, P, Grossmann, A, Morlet, J: Cycle octave and related transforms in seismic signal analysis. Geoexploration 23, 85-102 (1984)
- 11. Holschneider, M: Wavelets: An Analysis Tool. Clarendon Press, Oxford (1995)
- 12. Trimèche, K: Generalized Wavelets and Hypergroups. Gordon & Breach, New York (1997)
- Gasmi, A, Ben Mohamed, H, Bettaieb, N: Inversion of the Weinstein intertwining operator and its dual using Weinstein wavelets. An. Ştiinţ. Univ. 'Ovidius' Constanţa, Ser. Mat. 24(1), 289-307 (2016)
- 14. Daubechies, I: Time-frequency localization operators: a geometric phase space approach. IEEE Trans. Inf. Theory 34(4), 605-612 (1988)
- Daubechies, I, Paul, T: Time-frequency localisation operators a geometric phase space approach: II. The use of dilations. Inverse Probl. 4(3), 661-680 (1988)
- Daubechies, I: The wavelet transform, time-frequency localization and signal analysis. IEEE Trans. Inf. Theory 36(5), 961-1005 (1990)
- Wong, MW: Wavelet Transforms and Localization Operators. Operator Theory: Advances and Applications, vol. 136. Springer, Berlin (2002)
- 18. Ghobber, S, Jaming, P: Uncertainty principles for integral operators. Stud. Math. 220, 197-220 (2014)
- 19. Stein, EM: Interpolation of linear operators. Trans. Am. Math. Soc. 83, 482-492 (1956)
- Wilczok, E: New uncertainty principles for the continuous Gabor transform and the continuous wavelet transform. Doc. Math. 5, 201-226 (2000)
- 21. Mejjaoli, H, Sraieb, N: Uncertainty principles for the continuous Dunkl Gabor transform and the Dunkl continuous wavelet transform. Mediterr. J. Math. 5, 443-466 (2008)
- 22. Ghobber, S, Omri, S: Time-frequency concentration of the windowed Hankel transform. Integral Transforms Spec. Funct. 25(6), 481-496 (2014)
- Havin, V, Jöricke, B: The Uncertainty Principle in Harmonic Analysis. Ergebnisse der Mathematik und ihrer Grenzgebiete 3. Folge, vol. 28. Springer, Berlin (1994)
- 24. Donoho, DL, Stark, PB: Uncertainty principles and signal recovery. SIAM J. Appl. Math. 49, 906-931 (1989)
- Ben Salem, N, Ziad, N: A general form of Heisenberg-Pauli-Weyl uncertainty inequality for the Weinstein transform. Integral Transforms Spec. Funct. 24, 401-409 (2015)

Submit your manuscript to a SpringerOpen[®] journal and benefit from:

- ► Convenient online submission
- ► Rigorous peer review
- ► Open access: articles freely available online
- ► High visibility within the field
- Retaining the copyright to your article

Submit your next manuscript at > springeropen.com