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# Quantitative unique continuation for the linear coupled heat equations

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## Abstract

In this paper, we established a quantitative unique continuation results for a coupled heat equations, with the homogeneous Dirichlet boundary condition, on a bounded convex domain  $\Omega$  of  $\mathbb{R}^d$  with smooth boundary  $\partial\Omega$ . Our result shows that the value of the solutions can be determined uniquely by its value on an arbitrary open subset  $\omega$  of  $\Omega$  at any given positive time  $T$ .

**MSC:** 35K05; 93D15

**Keywords:** coupled heat equations; unique continuation; frequency functions

## 1 Introduction

In this paper, we consider an unique continuation of the following linear coupled heat equations:

$$\begin{cases} y_t(x, t) - \Delta y(x, t) - a(x, t)y(x, t) - b(x, t)z(x, t) = 0, & (x, t) \in \Omega \times (0, T), \\ z_t(x, t) - \Delta z(x, t) - c(x, t)y(x, t) - d(x, t)z(x, t) = 0, & (x, t) \in \Omega \times (0, T), \\ y(x, t) = 0, & (x, t) \in \partial\Omega \times (0, T), \\ z(x, t) = 0, & (x, t) \in \partial\Omega \times (0, T), \end{cases} \quad (1.1)$$

where  $\Omega$  is a bounded, convex domain of  $\mathbb{R}^d$  ( $d \geq 1$ ), with a smooth boundary  $\partial\Omega$ . Let  $\omega \subset \Omega$  be a nonempty and open subset of  $\Omega$ , and  $T$  is a positive number. Now we suppose that

$$a(x, t), b(x, t), c(x, t), d(x, t) \in L^\infty(\Omega \times (0, T)) \quad (1.2)$$

and

$$\max\{\|a\|_{L^\infty(\Omega \times (0, T))}, \|b\|_{L^\infty(\Omega \times (0, T))}, \|c\|_{L^\infty(\Omega \times (0, T))}, \|d\|_{L^\infty(\Omega \times (0, T))}\} \leq M, \quad (1.3)$$

where  $M$  is a positive number.

In this paper, we are concerned with the unique continuation for the solution of Eq. (1.1).

The main results obtained in this work can be stated as follows.

**Theorem 1.1** *Let  $\omega$  be a nonempty open subset of  $\Omega$ . Then there are two positive numbers  $p = p(\Omega, \omega)$  and  $C = C(\Omega, \omega)$  such that for each  $T > 0$  and for each potential  $a, b, c, d \in L^\infty(\Omega \times (0, T))$  with condition (1.3), any solution  $(y, z)$  with  $y(\cdot, 0) \equiv y_0(\cdot) \in L^2(\Omega)$ ,  $z(\cdot, 0) \equiv z_0(\cdot) \in L^2(\Omega)$ , to equation (1.1) has the following estimate:*

$$\begin{aligned} & \int_{\Omega} (|y(x, T)|^2 + |z(x, T)|^2) dx \\ & \leq C \exp\left(\frac{C}{T} + C(MT + M^2T^2)\right) \times \left(\int_{\Omega} (|y_0(x)|^2 + |z_0(x)|^2) dx\right)^{1-p} \\ & \quad \times \left(\int_{\omega} (|y(x, T)|^2 + |z(x, T)|^2) dx\right)^p. \end{aligned} \tag{1.4}$$

**Theorem 1.2** *Let  $\omega$  be a nonempty open subset of  $\Omega$ . If the initial data  $(y_0, z_0) \neq (0, 0)$ , and  $y_0, z_0 \in H_0^1(\Omega)$ , then there exists a positive number  $C = C(\Omega, \omega)$  such that for each  $T > 0$  and for each potential  $a, b, c, d \in L^\infty(\Omega \times (0, T))$  with condition (1.3), any solution  $(y, z)$  with  $y(\cdot, 0) \equiv y_0(\cdot)$ ,  $z(\cdot, 0) \equiv z_0(\cdot)$ , to equation (1.1) has the following estimate:*

$$\begin{aligned} & \|y_0\|_{L^2(\Omega)}^2 + \|z_0\|_{L^2(\Omega)}^2 \\ & \leq C e^{\frac{C}{T}} e^{C(MT + M^2T^2)} \exp\left(C(1 + 2M)Te^{CM^2T} \frac{\|y_0\|_{H_0^1(\Omega)}^2 + \|z_0\|_{H_0^1(\Omega)}^2}{\|y_0\|_{L^2(\Omega)}^2 + \|z_0\|_{L^2(\Omega)}^2}\right) \\ & \quad \times \int_{\omega} (|y(x, T)|^2 + |z(x, T)|^2) dx. \end{aligned} \tag{1.5}$$

**Remark 1.1** The constant  $C$  in (1.4) or (1.5) stands for a positive constant only dependent on  $\Omega, \omega$ . This constant varies in different contexts.

The study of unique continuation for the solutions for PDEs began at the early of the last century. Besides its own interesting in PDEs theory, it also plays a key role in both inverse problem and control theory. The first quantitative result of strong unique continuation for parabolic equations was derived in 1974 in the literature [1]. In [1], the authors establish the unique continuation for the parabolic equations with time independent coefficients in terms of the eigenfunctions of the corresponding elliptic operator, and their results did not apply to parabolic equations with time dependent coefficients. After 1988, there were more results of unique continuation for parabolic equations, and we refer the reader to [2–9], and the rich work cited therein. In [10], the author discusses the unique continuation for stochastic counterpart. In our paper, we mainly study this property for the coupled heat equations. To the best of our knowledge, this topic has not been studied in past publications. With the aid of frequency function methods, we can establish these quantitative estimates (see [2, 6]).

The paper is organized as follows. In Section 2, some preliminary results are presented. The proofs of Theorem 1.1 and Theorem 1.2 will be given in Section 3.

**2 Preliminary results**

Given a positive number  $\lambda$ , we define

$$G_\lambda(x, t) = \frac{1}{(T - t + \lambda)^{d/2}} e^{-\frac{|x-x_0|^2}{4(T-t+\lambda)}}, \quad (x, t) \in \Omega \times (0, T). \tag{2.1}$$

Then, for each  $t \in [0, T]$ , we write

$$H_\lambda(t) = \int_\Omega (|y(x, t)|^2 + |z(x, t)|^2) G_\lambda(x, t) \, dx, \tag{2.2}$$

$$D_\lambda(t) = \int_\Omega (|\nabla y(x, t)|^2 + |\nabla z(x, t)|^2) G_\lambda(x, t) \, dx, \tag{2.3}$$

and

$$N_\lambda(t) = \frac{2D_\lambda(t)}{H_\lambda(t)}, \tag{2.4}$$

where  $(y(x, t), z(x, t))$  are the solutions of equation (1.1). The function  $N_\lambda(t)$  was first discussed in [11], and it was called frequency function (see also [2, 12], and [9]). Throughout this section, we always work under the assumption  $H_\lambda(t) \neq 0$ . Next, we will discuss the properties for the functions  $G_\lambda(x, t)$ ,  $H_\lambda(t)$ ,  $D_\lambda(t)$  and  $N_\lambda(t)$ . Now, we first fix a positive number  $r$  and a point  $x_0$  in the subset  $\omega$  such that  $B_r \subset \omega$ .  $B_r$  denotes the open ball, centered at the point  $x_0$  and of radius  $r$ , in  $\mathbb{R}^d$ . Write  $m = \sup_{x \in \Omega} |x - x_0|^2$ . Lemma 2.1 is taken from [2, 6].

**Lemma 2.1** *For each  $\lambda > 0$ , the function  $G_\lambda$  given in (2.1) holds the following four identities over  $\mathbb{R}^d \times [0, T]$ :*

$$\partial_t G_\lambda(x, t) + \Delta G_\lambda(x, t) = 0, \tag{2.5}$$

$$\nabla G_\lambda(x, t) = \frac{-(x - x_0)}{2(T - t + \lambda)} G_\lambda(x, t), \tag{2.6}$$

$$\partial_t^2 G_\lambda(x, t) = \frac{-1}{2(T - t + \lambda)} G_\lambda(x, t) + \frac{|x_i - x_{0i}|^2}{4(T - t + \lambda)^2} G_\lambda(x, t), \tag{2.7}$$

and, for  $i \neq j$ ,

$$\partial_i \partial_j G_\lambda(x, t) = \frac{(x_i - x_{0i})(x_j - x_{0j})}{4(T - t + \lambda)^2} G_\lambda(x, t). \tag{2.8}$$

**Lemma 2.2** *For each  $\lambda > 0$ , the following identity holds for  $t \in (0, T)$ :*

$$\frac{d}{dt} \ln H_\lambda(t) = -N_\lambda(t) + \frac{2}{H_\lambda(t)} \int_\Omega (y(\partial_t y - \Delta y) + z(\partial_t z - \Delta z)) G_\lambda \, dx. \tag{2.9}$$

*Proof* By a direct computation, we can obtain

$$\begin{aligned} H'_\lambda(t) &= 2 \int_\Omega (y \partial_t y + z \partial_t z) G_\lambda \, dx + \int_\Omega (y^2 + z^2) \partial_t G_\lambda \, dx \\ &= 2 \int_\Omega (y \partial_t y + z \partial_t z) G_\lambda \, dx - \int_\Omega (y^2 + z^2) \Delta G_\lambda \, dx \\ &= 2 \int_\Omega (y \partial_t y + z \partial_t z) G_\lambda \, dx - \int_\Omega \Delta (y^2 + z^2) G_\lambda \, dx \\ &= 2 \int_\Omega (y \partial_t y + z \partial_t z) G_\lambda \, dx - 2 \int_\Omega ((\nabla y)^2 + y \Delta y + (\nabla z)^2 + z \Delta z) G_\lambda \, dx \end{aligned}$$

$$\begin{aligned}
 &= 2 \int_{\Omega} (y(\partial_t y - \Delta y) + z(\partial_t z - \Delta z)) G_{\lambda} \, dx - 2 \int_{\Omega} (|\nabla y|^2 + |\nabla z|^2) G_{\lambda} \, dx \\
 &= 2 \int_{\Omega} (y(\partial_t y - \Delta y) + z(\partial_t z - \Delta z)) G_{\lambda} \, dx - 2D_{\lambda}(t).
 \end{aligned}
 \tag{2.10}$$

Therefore, for each  $t \in (0, T)$ , we have

$$\frac{d}{dt} \ln H_{\lambda}(t) = \frac{H'_{\lambda}(t)}{H_{\lambda}(t)} = -N_{\lambda}(t) + \frac{2}{H_{\lambda}(t)} \int_{\Omega} (y(\partial_t y - \Delta y) + z(\partial_t z - \Delta z)) G_{\lambda} \, dx.
 \tag{2.11}$$

This completes the proof of this lemma. □

**Lemma 2.3** *For each  $\lambda > 0$  and  $t \in (0, T)$ , the functions  $H_{\lambda}$  and  $D_{\lambda}$  defined in (2.2) and (2.3), respectively, satisfy*

$$\begin{aligned}
 2H'_{\lambda}(t)D_{\lambda}(t) &= - \left[ 2 \int_{\Omega} y \left( \partial_t y - \frac{x - x_0}{2(T - t + \lambda)} \nabla y + \frac{1}{2}(\Delta y - \partial_t y) \right) G_{\lambda} \, dx \right. \\
 &\quad \left. + 2 \int_{\Omega} z \left( \partial_t z - \frac{x - x_0}{2(T - t + \lambda)} \nabla z + \frac{1}{2}(\Delta z - \partial_t z) \right) G_{\lambda} \, dx \right]^2 \\
 &\quad + \left( \int_{\Omega} y(\Delta y - \partial_t y) G_{\lambda} \, dx + \int_{\Omega} z(\Delta z - \partial_t z) G_{\lambda} \, dx \right)^2.
 \end{aligned}
 \tag{2.12}$$

*Proof* By (2.5), (2.6), we can rewrite  $H'_{\lambda}(t)$  as follows:

$$\begin{aligned}
 H'_{\lambda}(t) &= 2 \int_{\Omega} (y \partial_t y + z \partial_t z) G_{\lambda} \, dx - \int_{\Omega} (y^2 + z^2) \Delta G_{\lambda} \, dx \\
 &= 2 \int_{\Omega} (y \partial_t y + z \partial_t z) G_{\lambda} \, dx + \int_{\Omega} \nabla (y^2 + z^2) \nabla G_{\lambda} \, dx \\
 &= 2 \int_{\Omega} (y \partial_t y + z \partial_t z) G_{\lambda} \, dx - \int_{\Omega} (2y \nabla y + 2z \nabla z) \frac{x - x_0}{2(T - t + \lambda)} G_{\lambda} \, dx \\
 &= 2 \int_{\Omega} y \left( \partial_t y - \frac{x - x_0}{2(T - t + \lambda)} \nabla y \right) G_{\lambda} \, dx + 2 \int_{\Omega} z \left( \partial_t z - \frac{x - x_0}{2(T - t + \lambda)} \nabla z \right) G_{\lambda} \, dx \\
 &= 2 \int_{\Omega} y \left( \partial_t y - \frac{x - x_0}{2(T - t + \lambda)} \nabla y + \frac{1}{2}(\Delta y - \partial_t y) \right) G_{\lambda} \, dx \\
 &\quad + 2 \int_{\Omega} z \left( \partial_t z - \frac{x - x_0}{2(T - t + \lambda)} \nabla z + \frac{1}{2}(\Delta z - \partial_t z) \right) G_{\lambda} \, dx \\
 &\quad - \int_{\Omega} y(\Delta y - \partial_t y) G_{\lambda} \, dx - \int_{\Omega} z(\Delta z - \partial_t z) G_{\lambda} \, dx.
 \end{aligned}
 \tag{2.13}$$

It follows from (2.6) that for each  $t \in (0, T)$

$$\begin{aligned}
 D_{\lambda}(t) &= \int_{\Omega} (|\nabla y|^2 + |\nabla z|^2) G_{\lambda} \, dx \\
 &= \int_{\Omega} \nabla y \nabla y G_{\lambda} \, dx + \int_{\Omega} \nabla z \nabla z G_{\lambda} \, dx \\
 &= \int_{\Omega} \operatorname{div}(y \nabla y G_{\lambda}) \, dx - \int_{\Omega} y \operatorname{div}(\nabla y G_{\lambda}) \, dx + \int_{\Omega} \operatorname{div}(z \nabla z G_{\lambda}) \, dx \\
 &\quad - \int_{\Omega} z \operatorname{div}(\nabla z G_{\lambda}) \, dx
 \end{aligned}$$

$$\begin{aligned}
 &= - \int_{\Omega} y \Delta y G_{\lambda} \, dx + \int_{\Omega} y \nabla y \frac{x - x_0}{2(T - t + \lambda)} G_{\lambda} \, dx - \int_{\Omega} z \Delta z G_{\lambda} \, dx \\
 &\quad + \int_{\Omega} z \nabla z \frac{x - x_0}{2(T - t + \lambda)} G_{\lambda} \, dx \\
 &= - \int_{\Omega} y \left( \partial_t y - \frac{x - x_0}{2(T - t + \lambda)} \nabla y + \frac{1}{2} (\Delta y - \partial_t y) \right) G_{\lambda} \, dx - \frac{1}{2} \int_{\Omega} y (\Delta y - \partial_t y) G_{\lambda} \, dx \\
 &\quad - \int_{\Omega} z \left( \partial_t z - \frac{x - x_0}{2(T - t + \lambda)} \nabla z + \frac{1}{2} (\Delta z - \partial_t z) \right) G_{\lambda} \, dx \\
 &\quad - \frac{1}{2} \int_{\Omega} z (\Delta z - \partial_t z) G_{\lambda} \, dx. \tag{2.14}
 \end{aligned}$$

This, combining with (2.13), shows that

$$\begin{aligned}
 2H'_{\lambda}(t)D_{\lambda}(t) &= - \left[ 2 \int_{\Omega} y \left( \partial_t y - \frac{x - x_0}{2(T - t + \lambda)} \nabla y + \frac{1}{2} (\Delta y - \partial_t y) \right) G_{\lambda} \, dx \right. \\
 &\quad \left. + 2 \int_{\Omega} z \left( \partial_t z - \frac{x - x_0}{2(T - t + \lambda)} \nabla z + \frac{1}{2} (\Delta z - \partial_t z) \right) G_{\lambda} \, dx \right]^2 \\
 &\quad + \left( \int_{\Omega} y (\Delta y - \partial_t y) G_{\lambda} \, dx + \int_{\Omega} z (\Delta z - \partial_t z) G_{\lambda} \, dx \right)^2. \tag{2.15}
 \end{aligned}$$

This completes the proof of this lemma. □

**Lemma 2.4** *For each  $\lambda > 0$  and  $t \in (0, T)$ , then we have*

$$\begin{aligned}
 D'_{\lambda}(t) &= - \int_{\partial\Omega} |\nabla y|^2 \partial_{\nu} G_{\lambda} \, d\sigma + 2 \int_{\partial\Omega} \partial_{\nu} y (\nabla y \nabla G_{\lambda}) \, d\sigma \\
 &\quad - \int_{\partial\Omega} |\nabla z|^2 \partial_{\nu} G_{\lambda} \, d\sigma + 2 \int_{\partial\Omega} \partial_{\nu} z (\nabla z \nabla G_{\lambda}) \, d\sigma \\
 &\quad - 2 \int_{\Omega} \left( \partial_t y - \frac{x - x_0}{2(T - t + \lambda)} \nabla y + \frac{1}{2} (\Delta y - \partial_t y) \right)^2 G_{\lambda} \, dx \\
 &\quad - 2 \int_{\Omega} \left( \partial_t z - \frac{x - x_0}{2(T - t + \lambda)} \nabla z + \frac{1}{2} (\Delta z - \partial_t z) \right)^2 G_{\lambda} \, dx \\
 &\quad + \frac{1}{2} \int_{\Omega} (\Delta y - \partial_t y)^2 G_{\lambda} \, dx + \frac{1}{2} \int_{\Omega} (\Delta z - \partial_t z)^2 G_{\lambda} \, dx + \frac{1}{(T - t + \lambda)} D_{\lambda}(t). \tag{2.16}
 \end{aligned}$$

*Proof* By a direct computation, we can derive

$$\begin{aligned}
 D'_{\lambda}(t) &= 2 \int_{\Omega} \nabla y \partial_t \nabla y G_{\lambda} \, dx + \int_{\Omega} |\nabla y|^2 \partial_t G_{\lambda} \, dx \\
 &\quad + 2 \int_{\Omega} \nabla z \partial_t \nabla z G_{\lambda} \, dx + \int_{\Omega} |\nabla z|^2 \partial_t G_{\lambda} \, dx \\
 &= 2 \int_{\Omega} \nabla y \partial_t \nabla y G_{\lambda} \, dx - \int_{\Omega} |\nabla y|^2 \Delta G_{\lambda} \, dx \\
 &\quad + 2 \int_{\Omega} \nabla z \partial_t \nabla z G_{\lambda} \, dx - \int_{\Omega} |\nabla z|^2 \Delta G_{\lambda} \, dx \\
 &= 2 \int_{\Omega} \operatorname{div}(\nabla y \partial_t y G_{\lambda}) \, dx - 2 \int_{\Omega} \partial_t y \operatorname{div}(\nabla y G_{\lambda}) \, dx - \int_{\Omega} |\nabla y|^2 \Delta G_{\lambda} \, dx
 \end{aligned}$$

$$\begin{aligned}
 &+ 2 \int_{\Omega} \operatorname{div}(\nabla z \partial_t z G_{\lambda}) \, dx - 2 \int_{\Omega} \partial_t z \operatorname{div}(\nabla z G_{\lambda}) \, dx - \int_{\Omega} |\nabla z|^2 \Delta G_{\lambda} \, dx \\
 &= -2 \int_{\Omega} \partial_t y \Delta y G_{\lambda} \, dx - 2 \int_{\Omega} \partial_t y \nabla y \nabla G_{\lambda} \, dx - \int_{\Omega} |\nabla y|^2 \Delta G_{\lambda} \, dx \\
 &\quad - 2 \int_{\Omega} \partial_t z \Delta z G_{\lambda} \, dx - 2 \int_{\Omega} \partial_t z \nabla z \nabla G_{\lambda} \, dx - \int_{\Omega} |\nabla z|^2 \Delta G_{\lambda} \, dx \\
 &= -2 \int_{\Omega} \partial_t y \Delta y G_{\lambda} \, dx + 2 \int_{\Omega} \partial_t y \nabla y \frac{x - x_0}{2(T - t + \lambda)} G_{\lambda} \, dx - \int_{\Omega} |\nabla y|^2 \Delta G_{\lambda} \, dx \\
 &\quad - 2 \int_{\Omega} \partial_t z \Delta z G_{\lambda} \, dx + 2 \int_{\Omega} \partial_t z \nabla z \frac{x - x_0}{2(T - t + \lambda)} G_{\lambda} \, dx - \int_{\Omega} |\nabla z|^2 \Delta G_{\lambda} \, dx. \tag{2.17}
 \end{aligned}$$

However,

$$\begin{aligned}
 |\nabla y|^2 \Delta G_{\lambda} &= \operatorname{div}(|\nabla y|^2 \nabla G_{\lambda}) - 2 \operatorname{div}(\nabla y (\nabla y \nabla G_{\lambda})) \\
 &\quad + 2 \Delta y \nabla y \nabla G_{\lambda} + 2 \sum_{i=1}^d \nabla y \partial_i y \partial_i \nabla G_{\lambda} \tag{2.18}
 \end{aligned}$$

and

$$\begin{aligned}
 |\nabla z|^2 \Delta G_{\lambda} &= \operatorname{div}(|\nabla z|^2 \nabla G_{\lambda}) - 2 \operatorname{div}(\nabla z (\nabla z \nabla G_{\lambda})) \\
 &\quad + 2 \Delta z \nabla z \nabla G_{\lambda} + 2 \sum_{i=1}^d \nabla z \partial_i z \partial_i \nabla G_{\lambda}. \tag{2.19}
 \end{aligned}$$

Now, we write

$$A = \int_{\partial\Omega} |\nabla y|^2 \partial_\nu G_{\lambda} \, d\sigma - 2 \int_{\partial\Omega} \partial_\nu y (\nabla y \nabla G_{\lambda}) \, d\sigma \tag{2.20}$$

and

$$B = \int_{\partial\Omega} |\nabla z|^2 \partial_\nu G_{\lambda} \, d\sigma - 2 \int_{\partial\Omega} \partial_\nu z (\nabla z \nabla G_{\lambda}) \, d\sigma. \tag{2.21}$$

Then, by (2.18), (2.19), (2.20), (2.21), we obtain

$$\begin{aligned}
 \int_{\Omega} |\nabla y|^2 \Delta G_{\lambda} \, dx &= A + 2 \int_{\Omega} \Delta y \nabla y \nabla G_{\lambda} \, dx + 2 \sum_{i=1}^d \int_{\Omega} \nabla y \partial_i y \partial_i \nabla G_{\lambda} \, dx \\
 &= A - 2 \int_{\Omega} \Delta y \nabla y \frac{x - x_0}{2(T - t + \lambda)} G_{\lambda} \, dx \\
 &\quad + 2 \sum_{i=1}^d \int_{\Omega} \partial_i y \partial_i y \left( \frac{-1}{2(T - t + \lambda)} G_{\lambda} + \frac{|x - x_0|^2}{4(T - t + \lambda)^2} G_{\lambda} \right) \, dx \\
 &\quad + 2 \sum_{i \neq j} \int_{\Omega} \partial_i y \partial_j y \frac{(x_i - x_{0i})(x_j - x_{0j})}{4(T - t + \lambda)^2} G_{\lambda} \, dx \\
 &= A - 2 \int_{\Omega} \Delta y \nabla y \frac{x - x_0}{2(T - t + \lambda)} G_{\lambda} \, dx + \int_{\Omega} |\nabla y|^2 \left( \frac{-1}{(T - t + \lambda)} \right) G_{\lambda} \, dx \\
 &\quad + 2 \int_{\Omega} \left( \frac{x - x_0}{2(T - t + \lambda)} \nabla y \right)^2 G_{\lambda} \, dx. \tag{2.22}
 \end{aligned}$$

In the same way, we get

$$\begin{aligned} \int_{\Omega} |\nabla z|^2 \Delta G_{\lambda} dx &= B - 2 \int_{\Omega} \Delta z \nabla z \frac{x - x_0}{2(T - t + \lambda)} G_{\lambda} dx + \int_{\Omega} |\nabla z|^2 \left( \frac{-1}{(T - t + \lambda)} \right) G_{\lambda} dx \\ &\quad + 2 \int_{\Omega} \left( \frac{x - x_0}{2(T - t + \lambda)} \nabla z \right)^2 G_{\lambda} dx. \end{aligned} \tag{2.23}$$

Thus, we can rewrite (2.17) as

$$\begin{aligned} D'_{\lambda}(t) &= -A - B - 2 \int_{\Omega} \left( \frac{x - x_0}{2(T - t + \lambda)} \nabla y \right)^2 G_{\lambda} dx - 2 \int_{\Omega} \left( \frac{x - x_0}{2(T - t + \lambda)} \nabla z \right)^2 G_{\lambda} dx \\ &\quad + 2 \int_{\Omega} (\Delta y + \partial_t y) \frac{x - x_0}{2(T - t + \lambda)} \nabla y G_{\lambda} dx + 2 \int_{\Omega} (\Delta z + \partial_t z) \frac{x - x_0}{2(T - t + \lambda)} \nabla z G_{\lambda} dx \\ &\quad - 2 \int_{\Omega} \partial_t y \Delta y G_{\lambda} dx - 2 \int_{\Omega} \partial_t z \Delta z G_{\lambda} dx + \frac{1}{(T - t + \lambda)} D_{\lambda}(t) \\ &= -A - B - 2 \int_{\Omega} \left( \partial_t y - \frac{x - x_0}{2(T - t + \lambda)} \nabla y + \frac{1}{2} (\Delta y - \partial_t y) \right)^2 G_{\lambda} dx \\ &\quad - 2 \int_{\Omega} \left( \partial_t z - \frac{x - x_0}{2(T - t + \lambda)} \nabla z + \frac{1}{2} (\Delta z - \partial_t z) \right)^2 G_{\lambda} dx \\ &\quad + 2 \int_{\Omega} \frac{1}{4} (\Delta y + \partial_t y)^2 G_{\lambda} dx - 2 \int_{\Omega} \partial_t y \Delta y G_{\lambda} dx \\ &\quad + 2 \int_{\Omega} \frac{1}{4} (\Delta z + \partial_t z)^2 G_{\lambda} dx - 2 \int_{\Omega} \partial_t z \Delta z G_{\lambda} dx + \frac{1}{(T - t + \lambda)} D_{\lambda}(t) \\ &= -A - B - 2 \int_{\Omega} \left( \partial_t y - \frac{x - x_0}{2(T - t + \lambda)} \nabla y + \frac{1}{2} (\Delta y - \partial_t y) \right)^2 G_{\lambda} dx \\ &\quad - 2 \int_{\Omega} \left( \partial_t z - \frac{x - x_0}{2(T - t + \lambda)} \nabla z + \frac{1}{2} (\Delta z - \partial_t z) \right)^2 G_{\lambda} dx \\ &\quad + \frac{1}{2} \int_{\Omega} (\Delta y - \partial_t y)^2 G_{\lambda} dx + \frac{1}{2} \int_{\Omega} (\Delta z - \partial_t z)^2 G_{\lambda} dx + \frac{1}{(T - t + \lambda)} D_{\lambda}. \end{aligned} \tag{2.24}$$

This completes the proof of this lemma. □

**Lemma 2.5** For each  $\lambda > 0$  and  $t \in (0, T)$ , it follows that

$$\frac{d}{dt} [(T - t + \lambda) N_{\lambda}(t)] \leq 4M^2(T + \lambda). \tag{2.25}$$

*Proof* Firstly, we compute  $N'_{\lambda}(t)$  as  $t \in (0, T)$ . By (2.24) and (2.15), we derive that

$$\begin{aligned} N'_{\lambda}(t) &= 2 \left( \frac{1}{H_{\lambda}(t)} \right)^2 [D'_{\lambda}(t) H_{\lambda}(t) - D_{\lambda}(t) H'_{\lambda}(t)] \\ &= \frac{2}{H_{\lambda}(t)} \left\{ -A - B + \frac{1}{(T - t + \lambda)} D_{\lambda}(t) \right. \\ &\quad - 2 \int_{\Omega} \left( \partial_t y - \frac{x - x_0}{2(T - t + \lambda)} \nabla y + \frac{1}{2} (\Delta y - \partial_t y) \right)^2 G_{\lambda} dx + \frac{1}{2} \int_{\Omega} (\Delta y - \partial_t y)^2 G_{\lambda} dx \\ &\quad \left. - 2 \int_{\Omega} \left( \partial_t z - \frac{x - x_0}{2(T - t + \lambda)} \nabla z + \frac{1}{2} (\Delta z - \partial_t z) \right)^2 G_{\lambda} dx + \frac{1}{2} \int_{\Omega} (\Delta z - \partial_t z)^2 G_{\lambda} dx \right\} \end{aligned}$$

$$\begin{aligned}
 & - \left( \frac{1}{H_\lambda(t)} \right)^2 \left\{ - \left[ 2 \int_\Omega y \left( \partial_t y - \frac{x-x_0}{2(T-t+\lambda)} \nabla y + \frac{1}{2} (\Delta y - \partial_t y) \right) G_\lambda dx \right. \right. \\
 & + 2 \int_\Omega z \left( \partial_t z - \frac{x-x_0}{2(T-t+\lambda)} \nabla z + \frac{1}{2} (\Delta z - \partial_t z) \right) G_\lambda dx \left. \right]^2 \\
 & + \left[ \int_\Omega y (\Delta y - \partial_t y) G_\lambda dx + \int_\Omega z (\Delta z - \partial_t z) G_\lambda dx \right]^2 \left. \right\} \\
 = & \frac{1}{(T-t+\lambda)} N_\lambda - 2 \frac{A+B}{H_\lambda(t)} \\
 & + \left( \frac{2}{H_\lambda(t)} \right)^2 \left\{ \left[ \int_\Omega y \left( \partial_t y - \frac{x-x_0}{2(T-t+\lambda)} \nabla y + \frac{1}{2} (\Delta y - \partial_t y) \right) G_\lambda dx \right. \right. \\
 & + \int_\Omega z \left( \partial_t z - \frac{x-x_0}{2(T-t+\lambda)} \nabla z + \frac{1}{2} (\Delta z - \partial_t z) \right) G_\lambda dx \left. \right]^2 \\
 & - \int_\Omega \left( \partial_t y - \frac{x-x_0}{2(T-t+\lambda)} \nabla y + \frac{1}{2} (\Delta y - \partial_t y) \right)^2 G_\lambda dx \int_\Omega (y^2 + z^2) G_\lambda dx \\
 & - \int_\Omega \left( \partial_t z - \frac{x-x_0}{2(T-t+\lambda)} \nabla z + \frac{1}{2} (\Delta z - \partial_t z) \right)^2 G_\lambda dx \int_\Omega (y^2 + z^2) G_\lambda dx \left. \right\} \\
 & - \left( \frac{1}{H_\lambda(t)} \right)^2 \left[ \int_\Omega y (\Delta y - \partial_t y) G_\lambda dx + \int_\Omega z (\Delta z - \partial_t z) G_\lambda dx \right]^2 \\
 & + \frac{1}{H_\lambda(t)} \left[ \int_\Omega (\Delta y - \partial_t y)^2 G_\lambda dx + \int_\Omega (\Delta z - \partial_t z)^2 G_\lambda dx \right]. \tag{2.26}
 \end{aligned}$$

Let

$$\begin{aligned}
 \alpha &= \partial_t y - \frac{x-x_0}{2(T-t+\lambda)} \nabla y + \frac{1}{2} (\Delta y - \partial_t y), \\
 \beta &= \partial_t z - \frac{x-x_0}{2(T-t+\lambda)} \nabla z + \frac{1}{2} (\Delta z - \partial_t z).
 \end{aligned}$$

It follows from the Cauchy-Schwarz inequality that

$$\begin{aligned}
 & \int_\Omega \alpha^2 G_\lambda dx \int_\Omega (y^2 + z^2) G_\lambda dx + \int_\Omega \beta^2 G_\lambda dx \int_\Omega (y^2 + z^2) G_\lambda dx \\
 &= \int_\Omega (\alpha^2 + \beta^2) G_\lambda dx \int_\Omega (y^2 + z^2) G_\lambda dx \\
 &\geq \left[ \int_\Omega y \alpha G_\lambda dx + \int_\Omega z \beta G_\lambda dx \right]^2.
 \end{aligned}$$

It, together with (2.26), shows that

$$\begin{aligned}
 & N'_\lambda(t) - \frac{1}{(T-t+\lambda)} N_\lambda(t) + 2 \frac{A+B}{H_\lambda(t)} \\
 & - \frac{1}{H_\lambda(t)} \left[ \int_\Omega (\Delta y - \partial_t y)^2 G_\lambda dx + \int_\Omega (\Delta z - \partial_t z)^2 G_\lambda dx \right] \leq 0. \tag{2.27}
 \end{aligned}$$

In what follows, we will discuss the properties of  $A$  and  $B$ . Since  $y = z = 0$  on  $\partial\Omega$ , we have  $\nabla y = \partial_\nu y \nu$ ,  $\nabla z = \partial_\nu z \nu$  on  $\partial\Omega$ . If the domain  $\Omega$  is convex, we can derive that  $((x-x_0) \cdot \nu) \geq 0$ .



According to (2.20) and (2.6), then

$$\begin{aligned}
 A &= \int_{\partial\Omega} |\nabla y|^2 \partial_\nu G_\lambda \, d\sigma - 2 \int_{\partial\Omega} \partial_\nu y (\nabla y \nabla G_\lambda) \, d\sigma \\
 &= -\frac{1}{2(T-t+\lambda)} \int_{\partial\Omega} |\nabla y|^2 ((x-x_0) \cdot \nu) G_\lambda \, d\sigma + \frac{1}{(T-t+\lambda)} \\
 &\quad \times \int_{\partial\Omega} \partial_\nu y ((x-x_0) \nabla y) G_\lambda \, d\sigma \\
 &= -\frac{1}{2(T-t+\lambda)} \int_{\partial\Omega} |\nabla y|^2 ((x-x_0) \cdot \nu) G_\lambda \, d\sigma + \frac{1}{(T-t+\lambda)} \\
 &\quad \times \int_{\partial\Omega} |\partial_\nu y|^2 ((x-x_0) \cdot \nu) G_\lambda \, d\sigma \\
 &= \frac{1}{2(T-t+\lambda)} \times \int_{\partial\Omega} |\partial_\nu y|^2 ((x-x_0) \cdot \nu) G_\lambda \, d\sigma \geq 0.
 \end{aligned}$$

In the same way, we get

$$B = \frac{1}{2(T-t+\lambda)} \times \int_{\partial\Omega} |\partial_\nu z|^2 ((x-x_0) \cdot \nu) G_\lambda \, d\sigma \geq 0.$$

Combining with (2.27), we get

$$\begin{aligned}
 &(T-t+\lambda)N'_\lambda(t) - N_\lambda(t) \\
 &\quad - \frac{(T-t+\lambda)}{H_\lambda(t)} \left[ \int_\Omega (\Delta y - \partial_t y)^2 G_\lambda \, dx + \int_\Omega (\Delta z - \partial_t z)^2 G_\lambda \, dx \right] \leq 0. \tag{2.28}
 \end{aligned}$$

Using equations (1.1) and (2.28) can be written as

$$\begin{aligned}
 (T-t+\lambda)N'_\lambda(t) - N_\lambda(t) &\leq \frac{(T-t+\lambda)}{H_\lambda(t)} \left[ \int_\Omega (ay + bz)^2 G_\lambda \, dx + \int_\Omega (cy + dz)^2 G_\lambda \, dx \right] \\
 &\leq \frac{4M^2(T-t+\lambda)}{H_\lambda(t)} \int_\Omega (|y|^2 + |z|^2) G_\lambda \, dx \\
 &\leq 4M^2(T+\lambda).
 \end{aligned}$$

Thus,

$$\frac{d}{dt} [(T-t+\lambda)N_\lambda(t)] \leq 4M^2(T+\lambda), \quad \forall t \in (0, T).$$

This completes the proof of this lemma. □

Let

$$\mathcal{K}_{M,y,z,T} = 2 \ln \left( \frac{\int_\Omega (|y(x,0)|^2 + |z(x,0)|^2) \, dx}{\int_\Omega (|y(x,T)|^2 + |z(x,T)|^2) \, dx} \right) + \frac{m}{T} + 8MT + 4M^2 T^2 + \frac{d}{2}.$$

Then we have the following lemma.

**Lemma 2.6** For each  $\lambda > 0$ , we have

$$\lambda N_\lambda(T) + \frac{d}{2} \leq \left(\frac{\lambda}{T} + 1\right) \mathcal{K}_{M,y,z,T}. \tag{2.29}$$

*Proof* Integrating (2.25) over  $(t, T)$ , we have

$$\begin{aligned} \lambda N_\lambda(T) &\leq (T - t + \lambda)N_\lambda(t) + 4M^2(T + \lambda)(T - t) \\ &\leq (T + \lambda)N_\lambda(t) + 4M^2T(T + \lambda). \end{aligned}$$

Integrating the above over  $(0, \frac{T}{2})$ , we obtain

$$\frac{T}{2} \lambda N_\lambda(T) \leq (T + \lambda) \int_0^{\frac{T}{2}} N_\lambda(t) dt + 2M^2T^2(T + \lambda). \tag{2.30}$$

By integrating (2.11) over  $(0, \frac{T}{2})$ , we get

$$\begin{aligned} \int_0^{\frac{T}{2}} N_\lambda(t) dt &= - \int_0^{\frac{T}{2}} \frac{H'_\lambda(t)}{H_\lambda(t)} dt + \int_0^{\frac{T}{2}} \frac{2}{H_\lambda(t)} \int_\Omega [y(\partial_t y - \Delta y) + z(\partial_t z - \Delta z)] G_\lambda dx dt \\ &= - \ln \frac{H_\lambda(\frac{T}{2})}{H_\lambda(0)} + \int_0^{\frac{T}{2}} \frac{2}{H_\lambda(t)} \int_\Omega [-y(ay + bz) - z(cy + dz)] G_\lambda dx dt \\ &\leq \ln \frac{H_\lambda(0)}{H_\lambda(\frac{T}{2})} + \int_0^{\frac{T}{2}} \frac{2}{H_\lambda(t)} \int_\Omega 2M(y^2 + z^2) G_\lambda dx dt \\ &\leq \ln \frac{H_\lambda(0)}{H_\lambda(\frac{T}{2})} + 2MT. \end{aligned}$$

This, alone with (2.30), shows that

$$\frac{T}{2} \lambda N_\lambda(T) \leq (T + \lambda) \left[ \ln \frac{H_\lambda(0)}{H_\lambda(\frac{T}{2})} + 2MT + 2M^2T^2 \right].$$

We have

$$\begin{aligned} \frac{H_\lambda(0)}{H_\lambda(\frac{T}{2})} &= \frac{\int_\Omega (|y(x, 0)|^2 + |z(x, 0)|^2) (T + \lambda)^{-\frac{d}{2}} \cdot e^{-\frac{|x-x_0|^2}{4(T+\lambda)}} dx}{\int_\Omega (|y(x, \frac{T}{2})|^2 + |z(x, \frac{T}{2})|^2) (\frac{T}{2} + \lambda)^{-\frac{d}{2}} \cdot e^{-\frac{|x-x_0|^2}{4(\frac{T}{2}+\lambda)}} dx} \\ &\leq \frac{\int_\Omega (|y(x, 0)|^2 + |z(x, 0)|^2) dx \cdot (\frac{T}{2} + \lambda)^{\frac{d}{2}}}{\int_\Omega (|y(x, \frac{T}{2})|^2 + |z(x, \frac{T}{2})|^2) \cdot e^{-\frac{|x-x_0|^2}{4(\frac{T}{2}+\lambda)}} dx \cdot (T + \lambda)^{\frac{d}{2}}} \\ &\leq \frac{\int_\Omega (|y(x, 0)|^2 + |z(x, 0)|^2) dx \cdot (\frac{T}{2} + \lambda)^{\frac{d}{2}}}{\int_\Omega (|y(x, \frac{T}{2})|^2 + |z(x, \frac{T}{2})|^2) dx \cdot (T + \lambda)^{\frac{d}{2}}} \cdot e^{\frac{m}{4(\frac{T}{2}+\lambda)}} \\ &\leq e^{\frac{m}{2T}} \frac{\int_\Omega (|y(x, 0)|^2 + |z(x, 0)|^2) dx}{\int_\Omega (|y(x, \frac{T}{2})|^2 + |z(x, \frac{T}{2})|^2) dx}. \end{aligned}$$

Therefore,

$$\frac{T}{2} \lambda N_\lambda(T) \leq (T + \lambda) \left[ \frac{m}{2T} + \ln \left( \frac{\int_\Omega (|y(x, 0)|^2 + |z(x, 0)|^2) dx}{\int_\Omega (|y(x, \frac{T}{2})|^2 + |z(x, \frac{T}{2})|^2) dx} \right) + 2MT + 2M^2 T^2 \right].$$

Using the energy estimate

$$\frac{\int_\Omega (|y(x, T)|^2 + |z(x, T)|^2) dx}{\int_\Omega (|y(x, \frac{T}{2})|^2 + |z(x, \frac{T}{2})|^2) dx} \leq e^{2MT},$$

we obtain

$$\begin{aligned} \lambda N_\lambda(T) &\leq \left( \frac{\lambda}{T} + 1 \right) \left[ 2 \ln \left( \frac{\int_\Omega (|y(x, 0)|^2 + |z(x, 0)|^2) dx}{\int_\Omega (|y(x, T)|^2 + |z(x, T)|^2) dx} \right) + \frac{m}{T} + 8MT + 4M^2 T^2 \right] \\ &= \left( \frac{\lambda}{T} + 1 \right) \left( \mathcal{K}_{M,y,z,T} - \frac{d}{2} \right). \end{aligned}$$

Thus,

$$\begin{aligned} \lambda N_\lambda(T) + \frac{d}{2} &\leq \left( \frac{\lambda}{T} + 1 \right) \mathcal{K}_{M,y,z,T} - \frac{\lambda}{T} \cdot \frac{d}{2} \\ &\leq \left( \frac{\lambda}{T} + 1 \right) \mathcal{K}_{M,y,z,T}. \end{aligned}$$

This is (2.29). □

While most of the proof of the following lemma is similar to Lemma 3 of [2], we would rather give the proof in detail for the sake of completeness.

**Lemma 2.7** *For each  $T > 0$  and  $y_0 \in L^2(\Omega)$ ,  $z_0 \in L^2(\Omega)$ , the solution  $y$  and  $z$ , with  $y(\cdot, 0) = y_0(\cdot)$ ,  $z(\cdot, 0) = z_0(\cdot)$ , to (1.1) holds the estimate:*

$$\begin{aligned} &\left[ 1 - \frac{8\lambda}{r^2} \left( \frac{\lambda}{T} + 1 \right) \mathcal{K}_{M,y,z,T} \right] \int_\Omega |x - x_0|^2 (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \\ &\leq 8\lambda \left( \frac{\lambda}{T} + 1 \right) \mathcal{K}_{M,y,z,T} \int_{B_r} (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx. \end{aligned}$$

*Proof* We first point out the following fact: for any  $f \in H_0^1(\Omega)$  and for each  $\lambda > 0$ , we have  $0 \leq \int_\Omega |\nabla(f(x) \exp(-\frac{|x-x_0|^2}{8\lambda}))|^2 dx$ . By computing the right-hand term, we get

$$\begin{aligned} \int_\Omega \frac{|x - x_0|^2}{8\lambda} |f(x)|^2 e^{-\frac{|x-x_0|^2}{4\lambda}} dx &\leq 2\lambda \int_\Omega |\nabla f(x)|^2 e^{-\frac{|x-x_0|^2}{4\lambda}} dx \\ &\quad + \frac{d}{2} \int_\Omega |f(x)|^2 e^{-\frac{|x-x_0|^2}{4\lambda}} dx. \end{aligned} \tag{2.31}$$

It follows from (2.31) and (2.4) that

$$\begin{aligned} &\int_\Omega |x - x_0|^2 (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \\ &\leq 8\lambda \left( 2\lambda \int_\Omega (|\nabla y(x, T)|^2 + |\nabla z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \right) \end{aligned}$$

$$\begin{aligned}
 & + \frac{d}{2} \int_{\Omega} (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \\
 & \leq 8\lambda \left( \lambda N_{\lambda}(T) + \frac{d}{2} \right) \int_{\Omega} (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \\
 & \leq 8\lambda \left( \lambda N_{\lambda}(T) + \frac{d}{2} \right) \left[ \int_{B_r} (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \right. \\
 & \quad \left. + \frac{1}{r^2} \int_{\Omega \setminus B_r} |x - x_0|^2 (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \right]. \tag{2.32}
 \end{aligned}$$

It, combining with (2.29), shows that

$$\begin{aligned}
 & \int_{\Omega} |x - x_0|^2 (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \\
 & \leq 8\lambda \left( \frac{\lambda}{T} + 1 \right) \mathcal{K}_{M,y,z,T} \left[ \int_{B_r} (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \right. \\
 & \quad \left. + \frac{1}{r^2} \int_{\Omega} |x - x_0|^2 (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \right]. \tag{2.33}
 \end{aligned}$$

We get

$$\begin{aligned}
 & \left[ 1 - \frac{8\lambda}{r^2} \left( \frac{\lambda}{T} + 1 \right) \mathcal{K}_{M,y,z,T} \right] \int_{\Omega} |x - x_0|^2 (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \\
 & \leq 8\lambda \left( \frac{\lambda}{T} + 1 \right) \mathcal{K}_{M,y,z,T} \int_{B_r} (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx. \tag{2.34}
 \end{aligned}$$

This completes the proof of the lemma. □

### 3 Results and discussion

#### 3.1 The proof of Theorem 1.1

*Proof* We start with proving (1.4). We take  $\lambda > 0$  in the estimate of Lemma 2.7 to be such that

$$\frac{8\lambda}{r^2} \left( \frac{\lambda}{T} + 1 \right) \mathcal{K}_{M,y,z,T} = \frac{1}{2}. \tag{3.1}$$

By direct computation, we have

$$\lambda = \frac{1}{2} \left( -T + \sqrt{T^2 + \frac{Tr^2}{4\mathcal{K}_{M,y,z,T}}} \right).$$

Since  $\frac{m}{T} \leq \mathcal{K}_{M,y,z,T}$ , it follows that

$$\begin{aligned}
 \frac{1}{\lambda} & = 2 \frac{T + \sqrt{T^2 + \frac{Tr^2}{4\mathcal{K}_{M,y,z,T}}}}{\frac{Tr^2}{4\mathcal{K}_{M,y,z,T}}} \\
 & = 8 \left( T + \sqrt{T^2 + \frac{Tr^2}{4\mathcal{K}_{M,y,z,T}}} \right) \frac{1}{Tr^2} \mathcal{K}_{M,y,z,T}
 \end{aligned}$$

$$\begin{aligned} &\leq 8 \left( 2T + \sqrt{\frac{Tr^2}{4\mathcal{K}_{M,y,z,T}}} \right) \frac{1}{Tr^2} \mathcal{K}_{M,y,z,T} \\ &\leq \left( 16 + \frac{4r}{\sqrt{m}} \right) \frac{1}{r^2} \mathcal{K}_{M,y,z,T}. \end{aligned}$$

Therefore, we have

$$\begin{aligned} e^{\frac{m}{4\lambda}} &\leq e^{(4m+r\sqrt{m})\frac{1}{r^2} \mathcal{K}_{M,y,z,T}} \\ &\leq e^{(4m+r\sqrt{m})\frac{1}{r^2} \frac{d}{2}} e^{(4m+r\sqrt{m})\frac{1}{r^2} (\frac{m}{T} + 8MT + 4M^2 T^2)} \\ &\quad \times \left( \frac{\int_{\Omega} (|y(x, 0)|^2 + |z(x, 0)|^2) dx}{\int_{\Omega} (|y(x, T)|^2 + |z(x, T)|^2) dx} \right)^{2(4m+r\sqrt{m})/r^2}. \end{aligned} \tag{3.2}$$

By Lemma 2.7, we get

$$\begin{aligned} &\int_{\Omega} |x - x_0|^2 (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \\ &\leq r^2 \int_{B_r} (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx. \end{aligned}$$

Then we derive that

$$\begin{aligned} &\int_{\Omega} (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{m}{4\lambda}} dx \\ &\leq \int_{\Omega} (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \\ &\leq \int_{\Omega \cap |x-x_0| \geq r} (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \\ &\quad + \int_{B_r} (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \\ &\leq \frac{1}{r^2} \int_{\Omega} |x - x_0|^2 (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \\ &\quad + \int_{B_r} (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \\ &\leq 2 \int_{B_r} (|y(x, T)|^2 + |z(x, T)|^2) e^{-\frac{|x-x_0|^2}{4\lambda}} dx \\ &\leq 2 \int_{B_r} (|y(x, T)|^2 + |z(x, T)|^2) dx. \end{aligned}$$

Thus,

$$\begin{aligned} &\int_{\Omega} (|y(x, T)|^2 + |z(x, T)|^2) dx \\ &\leq 2e^{\frac{m}{4\lambda}} \int_{B_r} (|y(x, T)|^2 + |z(x, T)|^2) dx \\ &\leq 2e^{(4m+r\sqrt{m})\frac{1}{r^2} \frac{d}{2}} e^{(4m+r\sqrt{m})\frac{1}{r^2} (\frac{m}{T} + 8MT + 4M^2 T^2)} \end{aligned}$$

$$\begin{aligned}
 & \times \left( \frac{\int_{\Omega} (|y(x, 0)|^2 + |z(x, 0)|^2) dx}{\int_{\Omega} (|y(x, T)|^2 + |z(x, T)|^2) dx} \right)^{2(4m+r\sqrt{m})/r^2} \\
 & \times \int_{B_r} (|y(x, T)|^2 + |z(x, T)|^2) dx. \tag{3.3}
 \end{aligned}$$

Then we have

$$\begin{aligned}
 & \int_{\Omega} (|y(x, T)|^2 + |z(x, T)|^2) dx \\
 & \leq C e^{\frac{C}{r^2}} e^{\frac{C}{T^2}} e^{\frac{C}{r^2}(MT+M^2T^2)} \left( \frac{\int_{\Omega} (|y(x, 0)|^2 + |z(x, 0)|^2) dx}{\int_{\Omega} (|y(x, T)|^2 + |z(x, T)|^2) dx} \right)^{C/r^2} \\
 & \times \int_{B_r} (|y(x, T)|^2 + |z(x, T)|^2) dx,
 \end{aligned}$$

which is equivalent to the following inequality:

$$\begin{aligned}
 & \int_{\Omega} (|y(x, T)|^2 + |z(x, T)|^2) dx \\
 & \leq C e^{C/T} e^{C(MT+M^2T^2)} \left( \int_{\Omega} (|y_0(x)|^2 + |z_0(x)|^2) dx \right)^{\frac{C}{r^2+C}} \\
 & \times \left( \int_{B_r} (|y(x, T)|^2 + |z(x, T)|^2) dx \right)^{\frac{r^2}{r^2+C}} \\
 & \leq C \exp\left(\frac{C}{T} + C(MT + M^2T^2)\right) \left( \int_{\Omega} (|y_0(x)|^2 + |z_0(x)|^2) dx \right)^{\frac{C}{r^2+C}} \\
 & \times \left( \int_{\omega} (|y(x, T)|^2 + |z(x, T)|^2) dx \right)^{\frac{r^2}{r^2+C}}.
 \end{aligned}$$

Let  $p = \frac{r^2}{r^2+C}$  and then the above inequality can be written as

$$\begin{aligned}
 & \int_{\Omega} (|y(x, T)|^2 + |z(x, T)|^2) dx \\
 & \leq C \exp\left(\frac{C}{T} + C(MT + M^2T^2)\right) \left( \int_{\Omega} (|y_0(x)|^2 + |z_0(x)|^2) dx \right)^{1-p} \\
 & \times \left( \int_{\omega} (|y(x, T)|^2 + |z(x, T)|^2) dx \right)^p. \tag{3.4}
 \end{aligned}$$

Thus, we can obtain (1.4). This completes the proof of this theorem. □

### 3.2 The proof of Theorem 1.2

*Proof* In order to give (1.5), we should first prove the following backward uniqueness estimate:

$$\frac{\|y(x, 0)\|_{L^2}^2 + \|z(x, 0)\|_{L^2}^2}{\|y(x, T)\|_{L^2}^2 + \|z(x, T)\|_{L^2}^2} \leq \exp\left(C(1 + 2M)Te^{CM^2T} \frac{\|y(x, 0)\|_{H_0^1}^2 + \|z(x, 0)\|_{H_0^1}^2}{\|y(x, 0)\|_{L^2}^2 + \|z(x, 0)\|_{L^2}^2}\right). \tag{3.5}$$

For  $y_0, z_0 \in H_0^1(\Omega)$ , the solutions of equation (1.1)  $y(x, t), z(x, t) \in L^2(0, T; H^2(\Omega) \cap H_0^1(\Omega))$ . Then we can define a function  $\Phi(t)$  as follows:

$$\Phi(t) = \frac{\|y(x, t)\|_{H_0^1}^2 + \|z(x, t)\|_{H_0^1}^2}{\|y(x, t)\|_{L^2}^2 + \|z(x, t)\|_{L^2}^2} \geq C.$$

Let  $f = ay + bz$ , and  $g = cy + dz$ . By direct computation, we obtain

$$\frac{d}{dt} \left( \frac{1}{2} \|y\|_{L^2}^2 \right) = -\|y\|_{H_0^1}^2 - \langle ay + bz, y \rangle, \tag{3.6}$$

$$\frac{d}{dt} \left( \frac{1}{2} \|z\|_{L^2}^2 \right) = -\|z\|_{H_0^1}^2 - \langle cy + dz, z \rangle, \tag{3.7}$$

$$\frac{d}{dt} \left( \frac{1}{2} \|y\|_{H_0^1}^2 \right) = -\|\Delta y\|_{L^2}^2 + \int_{\Omega} \Delta y (ay + bz) \, dx,$$

$$\frac{d}{dt} \left( \frac{1}{2} \|z\|_{H_0^1}^2 \right) = -\|\Delta z\|_{L^2}^2 + \int_{\Omega} \Delta z (cy + dz) \, dx.$$

Then

$$\begin{aligned} \frac{d}{dt} \Phi(t) &= \frac{(\|y\|_{H_0^1}^2 + \|z\|_{H_0^1}^2)'(\|y\|_{L^2}^2 + \|z\|_{L^2}^2) - (\|y\|_{H_0^1}^2 + \|z\|_{H_0^1}^2)(\|y\|_{L^2}^2 + \|z\|_{L^2}^2)'}{(\|y\|_{L^2}^2 + \|z\|_{L^2}^2)^2} \\ &= \frac{2}{(\|y\|_{L^2}^2 + \|z\|_{L^2}^2)^2} \\ &\quad \times \left\{ \left( -\|\Delta y\|_{L^2}^2 - \|\Delta z\|_{L^2}^2 + \int_{\Omega} f \Delta y \, dx + \int_{\Omega} g \Delta z \, dx \right) (\|y\|_{L^2}^2 + \|z\|_{L^2}^2) \right. \\ &\quad \left. + \left( \|y\|_{H_0^1}^2 + \|z\|_{H_0^1}^2 + \int_{\Omega} fy \, dx + \int_{\Omega} gz \, dx \right) (\|y\|_{H_0^1}^2 + \|z\|_{H_0^1}^2) \right\} \\ &= \frac{2}{(\|y\|_{L^2}^2 + \|z\|_{L^2}^2)^2} \\ &\quad \times \left\{ \left( -\|\Delta y\|_{L^2}^2 - \|\Delta z\|_{L^2}^2 + \int_{\Omega} f \Delta y \, dx + \int_{\Omega} g \Delta z \, dx \right) (\|y\|_{L^2}^2 + \|z\|_{L^2}^2) \right. \\ &\quad \left. + \left( \|y\|_{H_0^1}^2 + \|z\|_{H_0^1}^2 + \frac{1}{2} \left( \int_{\Omega} fy \, dx + \int_{\Omega} gz \, dx \right) \right)^2 - \frac{1}{4} \left( \int_{\Omega} fy \, dx + \int_{\Omega} gz \, dx \right)^2 \right\}. \end{aligned}$$

Integrating by parts and using Cauchy-Schwarz inequality, we have

$$\begin{aligned} \frac{d}{dt} \Phi(t) &= \frac{2}{(\|y\|_{L^2}^2 + \|z\|_{L^2}^2)^2} \\ &\quad \times \left\{ \left( -\|\Delta y\|_{L^2}^2 - \|\Delta z\|_{L^2}^2 + \int_{\Omega} f \Delta y \, dx + \int_{\Omega} g \Delta z \, dx \right) (\|y\|_{L^2}^2 + \|z\|_{L^2}^2) \right. \\ &\quad \left. + \left( \int_{\Omega} \left( -\Delta y + \frac{f}{2} \right) y \, dx + \int_{\Omega} \left( -\Delta z + \frac{g}{2} \right) z \, dx \right)^2 - \frac{1}{4} \left( \int_{\Omega} fy \, dx + \int_{\Omega} gz \, dx \right)^2 \right\} \\ &\leq \frac{2}{(\|y\|_{L^2}^2 + \|z\|_{L^2}^2)^2} \\ &\quad \times \left\{ \left( -\|\Delta y\|_{L^2}^2 - \|\Delta z\|_{L^2}^2 + \int_{\Omega} f \Delta y \, dx + \int_{\Omega} g \Delta z \, dx \right) (\|y\|_{L^2}^2 + \|z\|_{L^2}^2) \right. \end{aligned}$$

$$\begin{aligned}
 & + \left( \left\| -\Delta y + \frac{f}{2} \right\|_{L^2}^2 + \left\| -\Delta z + \frac{g}{2} \right\|_{L^2}^2 \right)^2 - \frac{1}{4} \left( \int_{\Omega} f y \, dx + \int_{\Omega} g z \, dx \right)^2 \Big\} \\
 & \leq \frac{2}{(\|y\|_{L^2}^2 + \|z\|_{L^2}^2)^2} \\
 & \quad \times \left\{ \left( -\| \Delta y \|_{L^2}^2 - \| \Delta z \|_{L^2}^2 + \int_{\Omega} f \Delta y \, dx + \int_{\Omega} g \Delta z \, dx \right) (\|y\|_{L^2}^2 + \|z\|_{L^2}^2) \right. \\
 & \quad \left. + \left( \left\| -\Delta y + \frac{f}{2} \right\|_{L^2}^2 + \left\| -\Delta z + \frac{g}{2} \right\|_{L^2}^2 \right) (\|y\|_{L^2}^2 + \|z\|_{L^2}^2) \right\}.
 \end{aligned}$$

Thus,

$$\begin{aligned}
 \frac{d}{dt} \Phi(t) & \leq \frac{2}{(\|y\|_{L^2}^2 + \|z\|_{L^2}^2)^2} \left( \left\| \frac{f}{2} \right\|_{L^2}^2 + \left\| \frac{g}{2} \right\|_{L^2}^2 \right) (\|y\|_{L^2}^2 + \|z\|_{L^2}^2) \\
 & = \frac{\|f\|_{L^2}^2 + \|g\|_{L^2}^2}{2(\|y\|_{L^2}^2 + \|z\|_{L^2}^2)} = \frac{\|ay + bz\|_{L^2}^2 + \|cy + dz\|_{L^2}^2}{2(\|y\|_{L^2}^2 + \|z\|_{L^2}^2)} \\
 & \leq CM^2 \frac{\|y\|_{H_0^1}^2 + \|z\|_{H_0^1}^2}{\|y\|_{L^2}^2 + \|z\|_{L^2}^2} = CM^2 \Phi(t).
 \end{aligned}$$

With the aid of Gronwall’s inequality, we obtain, for  $t \in (0, T)$ ,

$$\Phi(t) \leq e^{CM^2 T} \Phi(0). \tag{3.8}$$

This, combining with (3.6) and (3.7), indicates that

$$\begin{aligned}
 0 & = \frac{1}{2} \frac{d}{dt} (\|y\|_{L^2}^2 + \|z\|_{L^2}^2) + \|y\|_{H_0^1}^2 + \|z\|_{H_0^1}^2 + \langle ay + bz, y \rangle + \langle cy + dz, z \rangle \\
 & \leq \frac{1}{2} \frac{d}{dt} (\|y\|_{L^2}^2 + \|z\|_{L^2}^2) + \Phi(t) (\|y\|_{L^2}^2 + \|z\|_{L^2}^2) + 2M \Phi(t) (\|y\|_{L^2}^2 + \|z\|_{L^2}^2) \\
 & \leq \frac{1}{2} \frac{d}{dt} (\|y\|_{L^2}^2 + \|z\|_{L^2}^2) + (1 + 2M) e^{CM^2 T} \Phi(0) (\|y\|_{L^2}^2 + \|z\|_{L^2}^2). \tag{3.9}
 \end{aligned}$$

Integrating (3.9) on  $(0, T)$ , we get the desired estimate

$$\|y(x, 0)\|_{L^2}^2 + \|z(x, 0)\|_{L^2}^2 \leq \exp(2(1 + 2M)Te^{CM^2 T} \Phi(0)) (\|y(x, T)\|_{L^2}^2 + \|z(x, T)\|_{L^2}^2).$$

Thus,

$$\frac{\|y(x, 0)\|_{L^2}^2 + \|z(x, 0)\|_{L^2}^2}{\|y(x, T)\|_{L^2}^2 + \|z(x, T)\|_{L^2}^2} \leq \exp \left( 2(1 + 2M)Te^{CM^2 T} \frac{\|y(x, 0)\|_{H_0^1}^2 + \|z(x, 0)\|_{H_0^1}^2}{\|y(x, 0)\|_{L^2}^2 + \|z(x, 0)\|_{L^2}^2} \right).$$

This is the proof of (3.5), and we have completed the proof of Theorem 1.2. □

### 4 Conclusions

In this work, we discuss the unique continuation for the linear coupled heat equations. Our results demonstrate that the value of the solution equation (1.1) can be determined uniquely by its value on an arbitrary open subset  $\omega$  of  $\Omega$  at any given positive time  $T$ . In



other words, if the solution of equation (1.1) satisfies  $y(\cdot, T) = z(\cdot, T) = 0$  on  $\omega$ , then  $y = z = 0$  on  $\Omega \times (0, T)$ . This is the quantitative version of the unique continuation property for equation (1.1).

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#### Competing interests

The authors declare that they have no competing interests.

#### Authors' contributions

GZ provided the question. GZ, KL and JL gave the proof for the main result together. All authors read and approved the final manuscript.

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