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New periodic solutions of singular Hamiltonian systems with fixed energies

Fengying Li^{1*}, Qingqing Hua² and Shiqing Zhang²

^{*}Correspondence: lify0308@163.com ¹School of Economic and Mathematics, Southwestern University of Finance and Economics, Chengdu, Sichuan 611130, P.R. China Full list of author information is available at the end of the article

Abstract

By using the variational minimizing method with a special constraint and the direct variational minimizing method without constraint, we study second-order Hamiltonian systems with a singular potential $V \in C^2(R^n \setminus O, R)$ and $V \in C^1(R^2 \setminus O, R)$, which may have an unbounded potential well, and prove the existence of non-trivial periodic solutions with a prescribed energy. Our results can be regarded as complements of the well-known theorems of Benci-Gluck-Ziller-Hayashi and Ambrosetti-Coti Zelati and so on. **MSC:** 35A15; 47J30

Keywords: second-order singular Hamiltonian systems; periodic solutions; variational methods

1 Introduction

Seifert [1] in 1948 and Rabinowitz [2, 3] in 1978 and 1979 studied classical second-order Hamiltonian systems without singularity, based on their work, Benci [4, 5] and Gluck and Ziller [6] and Hayashi [7] used a Jacobi metric and very complicated geodesic methods and algebraic topology to study the periodic solutions with a fixed energy of the following system:

$$\ddot{q} + V'(q) = O, \tag{1.1}$$

$$\frac{1}{2}|\dot{q}|^2 + V(q) = h. \tag{1.2}$$

They proved a very general theorem.

Theorem 1.1 Suppose $V \in C^2(\mathbb{R}^n, \mathbb{R})$, if

$$\left\{x \in \mathbb{R}^n | V(x) \le h\right\}$$

is bounded and non-empty, then (1.1)-(1.2) *has a periodic solution with energy h. Furthermore, if*

$$V'(x) \neq O$$
, $\forall x \in \{x \in \mathbb{R}^n | V(x) = h\}$,

then (1.1)-(1.2) has a nonconstant periodic solution with energy h.



©2014 Li et al.; licensee Springer. This is an Open Access article distributed under the terms of the Creative Commons Attribution License (http://creativecommons.org/licenses/by/2.0), which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited. For the existence of multiple periodic solutions for (1.1)-(1.2) with compact energy surfaces, we can refer to Groessen [8] and Long [9] and the references therein.

In the 1987 paper of Ambrosetti and Coti Zelati [10], Clark-Ekeland's dual action principle, Ambrosetti-Rabinowitz's mountain pass theorem *etc.* were used to study the existence of *T*-periodic solutions of the second-order equation

 $-\ddot{x} = \nabla U(x),$

where

$$U = V \in C^2(\Omega; \mathbf{R})$$

is such that

$$U(x) \to \infty$$
, $x \to \Gamma = \partial \Omega$;

here $\Omega \subset \mathbf{R}^n$ is a bounded and convex domain, and they got the following result.

Theorem 1.2 Suppose that

- (i) $U(O) = 0 = \min U$;
- (ii) $U(x) \le \theta(x, \nabla U(x))$ for some $\theta \in (0, \frac{1}{2})$ and for all x near Γ (superquadraticity near Γ);
- (iii) $(U''(x)y, y) \ge k|y|^2$ for some k > 0 and for all $(x, y) \in \Omega \times \mathbb{R}^N$.

Let ω_N be the greatest eigenvalue of U''(0) and $T_0 = (2/\omega_N)^{1/2}$. Then $-\ddot{x} = \nabla U(x)$ has for each $T \in (0, T_0)$ a periodic solution with minimal period T.

For C^r systems, a natural interesting problem is if

 $\left\{x \in \mathbb{R}^n | V(x) \le h\right\}$

is unbounded: can we get a nonconstant periodic solution for the system (1.1)-(1.2)?

In 1987, Offin [11] firstly generalized Theorem 1.1 to some non-compact cases under $V \in C^3(\mathbb{R}^n, \mathbb{R})$ and complicated geometrical assumptions on potential wells, but it seems to be difficult to verify this for concrete potentials under the geometrical conditions.

In 1988, Rabinowitz [12] studied multiple periodic solutions for classical Hamiltonian systems with potential $V \in C^1(R \times \mathbb{R}^n, \mathbb{R})$, where $V(q_1, \ldots, q_n; t)$ is T_i -periodic in positions $q_i \in \mathbb{R}$ and is T-periodic in t.

In 1990, using Clark-Ekeland's dual variational principle and Ambrosetti-Rabinowitz's mountain pass lemma, Coti Zelati *et al.* [13] studied Hamiltonian systems with convex potential wells, they got the following result.

Theorem 1.3 Let Ω be a convex open subset of \mathbb{R}^n containing the origin Ω . Let $V \in C^2(\Omega, \mathbb{R})$ be such that

- (V1) $V(q) \ge V(O) = 0, \forall q \in \Omega.$
- (V2) $\forall q \neq O, V''(q) > 0.$
- (V3) $\exists \omega > 0, s.t. V(q) \leq \frac{\omega}{2} ||q||^2, \forall ||q|| < \epsilon.$

(V4) $V''(q)^{-1} \to 0$, $||q|| \to 0$, or (V4)' $V''(q)^{-1} \to 0$, $q \to \partial \Omega$.

Then, for every $T < \frac{2\pi}{\sqrt{\omega}}$, (1.1) has a solution with minimal period T.

In Theorems 1.2 and 1.3, the authors assumed the convex conditions for potentials and potential wells so that they can apply Clark-Ekeland's dual variational principle; we notice that Theorems 1.1-1.3 essentially made the following assumption:

 $V(x) \to \infty$, $x \to \Gamma = \partial \Omega$.

So all the potential wells are bounded.

For singular Hamiltonian systems with a fixed energy $h \in R$, Ambrosetti and Coti Zelati in [14, 15] used Ljusternik-Schnirelmann theory on a C^1 manifold to get the following theorem.

Theorem 1.4 (Ambrosetti and Coti Zelati [14]) Suppose $V \in C^2(\mathbb{R}^n \setminus \{O\}, \mathbb{R})$ satisfies $V(q) \to -\infty, q \to 0$ and

- (A1) $3V'(u) \cdot u + (V''(u)u, u) \neq 0, \forall u \neq 0;$
- (A2) $V'(u) \cdot u > 0, \forall u \neq 0;$
- (A3) $\exists \alpha > 2, s.t. V'(u) \cdot u \leq -\alpha V(u), \forall u \neq 0;$
- (A4) $\exists \beta > 2, r > 0, s.t. V'(u) \cdot u \ge -\beta V(u), 0 < |u| < r;$
- (A5) $V(u) + \frac{1}{2}V'(u)u \le 0, \forall u \ne 0.$

Then (1.1)-(1.2) has at least one nonconstant periodic solution.

Besides Ambrosetti-Coti Zelati, many other mathematicians [16–34] studied singular Hamiltonian systems, here we only mention a related recent paper of Carminati, Sere and Tanaka [16]. They used complex variational and topological methods to generalize Pisani's results [17], and they got the following theorem.

Theorem 1.5 Suppose h > 0, $L_0 > 0$ and $V \in C^{\infty}(\mathbb{R}^n \setminus \{O\}, \mathbb{R})$ satisfies $V(q) \to -\infty$, $q \to 0$ and

- (B1) $V(q) \leq 0, \forall q \neq 0;$
- (B2) $V(q) + \frac{1}{2}V'(q)q \le h, \forall |q| \ge e^{L_0};$
- (B3) $V(q) + \frac{1}{2}V'(q)q \ge h, \forall |q| \le e^{-L_0};$
- (A4) $\exists \beta > 2, r > 0, s.t. V'(q) \cdot q \ge -\beta V(q), 0 < |q| < r.$

Then (1.1)-(1.2) has at least one periodic solution with the given energy h and whose action is at most $2\pi r_0$ with

 $r_0 = \max\{[2(h - V(q))]^{\frac{1}{2}}; |q| = 1\}.$

Theorem 1.6 Suppose h > 0, $\rho_0 > 0$, and $V \in C^{\infty}(\mathbb{R}^n \setminus \{O\}, \mathbb{R})$ satisfies $V(q) \to -\infty, q \to 0$ and (B1), (A4) and

 $\begin{array}{ll} (\mathrm{B2})' & \lim_{|q| \to +\infty} V'(q) = O; \\ (\mathrm{B3})' & V(q) + \frac{1}{2}V'(q)q \geq h, \forall |q| \leq \rho_0. \end{array}$

Then (1.1)-(1.2) has at least one periodic solution with the given energy h whose action is at most $2\pi r_0$.

By using the variational minimizing method with a special constraint, we obtain the following result.

Theorem 1.7 Suppose $V \in C^2(\mathbb{R}^n \setminus \{O\}, \mathbb{R})$ and $V(q) \to -\infty, q \to 0$ and satisfies (A1)-(A3) and

Then for any h > 0, (1.1)-(1.2) has at least one nonconstant periodic solution with the given energy h.

Using the direct variational minimizing method, we get the following theorem.

Theorem 1.8 Suppose $V \in C^1(\mathbb{R}^2 \setminus \{O\}, \mathbb{R})$ and $V(q) \to -\infty, q \to 0$ and satisfies

- (B1)' $V(q) < h, \forall q \neq O;$
- $(\mathrm{P1})' \quad V'(u) \to O, \, \|u\| \to +\infty;$
- $(A3)' \exists \alpha > 2, \mu_2 > 0, s.t. V'(u) \cdot u \leq -\alpha V(u) + \mu_2, \forall u \neq 0;$
- (A4) $\exists \beta > 2, r > 0, s.t. V'(u) \cdot u \ge -\beta V(u), 0 < |u| < r.$

Then for any $h > \frac{\mu_2}{\alpha}$, (1.1)-(1.2) has at least one nonconstant periodic solution with the given energy h.

Corollary 1.9 *Suppose* $\alpha = \beta > 2$ *and*

 $V(x) = -|x|^{-\alpha}.$

Then for any h > 0, (1.1)-(1.2) has at least one nonconstant periodic solution with the given energy h.

Remark In Theorem 1.8, the assumption on regularity for potential V is weaker than Theorems 1.1-1.6. Comparing Theorem 1.5 with Theorem 1.8, our (B1)' is also weaker than (B1), and (A3)' is also different from (B2)-(B3) and (B3)'.

2 A few lemmas

Let

$$H^{1} = W^{1,2}(R/Z, R^{n}) = \{ u : R \to R^{n}, u \in L^{2}, \dot{u} \in L^{2}, u(t+1) = u(t) \}$$

Then the standard H^1 norm is equivalent to

$$\|u\| = \|u\|_{H^1} = \left(\int_0^1 |\dot{u}|^2 \, dt\right)^{1/2} + \big|u(0)\big|.$$

Let

$$\Lambda = \left\{ u \in H^1 | u(t) \neq O, \forall t \right\}.$$

Lemma 2.1 ([14]) Let

$$F = \left\{ u \in H^1 \middle| \int_0^1 \left(V(u) + \frac{1}{2} V'(u) u \right) dt = h \right\}.$$

If (A1) holds, then F is a C^1 manifold with codimension 1 in H^1 . Let

$$f(u) = \frac{1}{4} \int_0^1 |\dot{u}|^2 dt \int_0^1 V'(u) u dt$$

and let $\widetilde{u} \in F$ be such that $f'(\widetilde{u}) = O$ and $f(\widetilde{u}) > 0$. Set

$$\frac{1}{T^2} = \frac{\int_0^1 V'(\widetilde{u})\widetilde{u} \, dt}{\int_0^1 |\dot{\widetilde{u}}|^2 \, dt}.$$

If (A2) holds, then $\tilde{q}(t) = \tilde{u}(t/T)$ is a nonconstant *T*-periodic solution for (1.1)-(1.2). Moreover, if (A2) holds, then $f(u) \ge 0$ on *F* and f(u) = 0, $u \in F$ if and only if *u* is constant.

Lemma 2.2 ([8, 14]) $Let f(u) = \frac{1}{2} \int_0^1 |\dot{u}|^2 dt \int_0^1 (h - V(u)) dt and \tilde{u} \in \Lambda$ be such that $f'(\tilde{u}) = O$ and $f(\tilde{u}) > 0$. Set

$$\frac{1}{T^2} = \frac{\int_0^1 (h - V(\widetilde{u})) \, dt}{\frac{1}{2} \int_0^1 |\dot{\widetilde{u}}|^2 \, dt}.$$

Then $\tilde{q}(t) = \tilde{u}(t/T)$ is a nonconstant *T*-periodic solution for (1.1)-(1.2). Furthermore, if V(x) < h, $\forall x \neq O$, then $f(u) \ge 0$ on Λ and f(u) = 0, $u \in \Lambda$ if and only if u is a nonzero constant.

Lemma 2.3 (Sobolev-Rellich-Kondrachov [35, 36])

 $W^{1,2}(R/Z, \mathbb{R}^n) \subset C(\mathbb{R}/Z, \mathbb{R}^n)$

and the imbedding is compact.

Lemma 2.4 ([35, 36]) Let $q \in W^{1,2}(R/TZ, R^n)$.

(1) If q(0) = q(T) = O, then we have the Friedrics-Poincaré inequality:

$$\int_0^T \left| \dot{q}(t) \right|^2 dt \ge \left(\frac{\pi}{T}\right)^2 \int_0^T \left| q(t) \right|^2 dt$$

(2) If $\int_0^T q(t) dt = 0$, then we have Wirtinger's inequality:

$$\int_0^T \left| \dot{q}(t) \right|^2 dt \ge \left(\frac{2\pi}{T} \right)^2 \int_0^T \left| q(t) \right|^2 dt$$

and Sobolev's inequality:

$$\int_0^T \left| \dot{q}(t) \right|^2 dt \ge \frac{12}{T} \left| q(t) \right|_\infty^2.$$

Lemma 2.5 (Eberlein-Shmulyan [37]) A Banach space X is reflexive if and only if any bounded sequence in X has a weakly convergent subsequence.

Definition 2.6 (Tonelli [35]) Let *X* be a Banach space, $f : X \rightarrow R$.

(i) If for any $\{x_n\} \subset X$ strongly converges to $x_0 : x_n \to x_0$, we have

 $\liminf f(x_n) \ge f(x_0),$

then we call f(x) lower semi-continuous at x_0 .

(ii) If for any $\{x_n\} \subset X$ weakly converges to $x_0 : x_n \rightharpoonup x_0$, we have

 $\liminf f(x_n) \ge f(x_0),$

then we call f(x) weakly lower semi-continuous at x_0 .

Using the famous Ekeland variational principle, Ekeland proved the following.

Lemma 2.7 (Ekeland [38]) Let X be a Banach space, $F \subset X$ be a closed (weakly closed) subset, let $\delta(x_1, x_2)$ be the geodesic distance between two points x_1 and x_2 in X, $\delta(x, F)$ be the geodesic distance between x and the set F. Suppose that Φ defined on X is Gateauxdifferentiable and lower semi-continuous (or weakly lower semi-continuous) and assume $\Phi|_F$ restricted on F is bounded from below. Then there is a sequence $\{x_n\} \subset F$ such that

$$\delta(x_n, F) \to 0,$$

$$\Phi(x_n) \to \inf_F \Phi,$$

$$\left(1 + \|x_n\|\right) \left\| \Phi \right\|_F'(x_n) \right\| \to 0$$

Definition 2.8 ([38, 39]) Let *X* be a Banach space, $F \subset X$ be a closed subset. Suppose that Φ defined on *X* is Gateaux-differentiable, if sequence $\{x_n\} \subset F$ is such that

$$\begin{split} \delta(x_n, F) &\to 0, \\ \Phi(x_n) &\to c, \\ \left(1 + \|x_n\|\right) \left\| \Phi \|'_F(x_n) \right\| &\to 0, \end{split}$$

then $\{x_n\}$ has a strongly convergent subsequence.

Then we say that *f* satisfies the $(CPS)_{c,F}$ condition at the level *c* for the closed subset $F \subset X$.

We notice that if F = X, then the above condition is the classical Cerami-Palais-Smale condition [40].

We can give a weaker condition than the $(CPS)_{c,F}$ condition.

Definition 2.9 Let *X* be a Banach space, $F \subset X$ be a weakly closed subset. Suppose that Φ defined on *X* is Gateaux-differentiable, if sequence $\{x_n\} \subset F$ such that

$$\delta(x_n, F) \to 0$$
,

$$\Phi(x_n) \to c,$$

 $\|\Phi\|'_F(x_n)\| \to 0,$

then $\{x_n\}$ has a weakly convergent subsequence.

Then we say that f satisfies the $(WCPS)_{c,F}$ condition.

Lemma 2.10 (Gordon [18]) Let V satisfy the so-called Gordon strong force condition: There exists a neighborhood \mathcal{N} of O and a function $U \in C^1(\Omega, \mathbb{R})$ such that:

(i)
$$\lim_{s\to 0} U(x) = -\infty;$$

(ii) $-V(x) \ge |U'(x)|^2$ for every $x \in \mathcal{N} - \{O\}$.
Let

$$\partial \Lambda = \left\{ u \in H^1 = W^{1,2}(\mathbb{R}/\mathbb{Z},\mathbb{R}^n), \exists t_0, u(t_0) = O \right\}.$$

Then we have

$$\int_0^1 V(u)\,dt\to -\infty,\quad \forall u_n\rightharpoonup u\in\partial\Lambda.$$

Let

$$\partial \Lambda = \{ u \in H^1 = W^{1,2}(R/Z, R^n), \exists t_0, u(t_0) = 0 \}.$$

Then we have

$$\int_0^1 V(u)\,dt\to -\infty,\quad \forall u_n\rightharpoonup u\in\partial\Lambda.$$

By Lemmas 2.7 and 2.10, it is easy to prove the following.

Lemma 2.11 Let X be a Banach space, let $F \subset X$ be a weakly closed subset. Suppose that Φ defined on F is Gateaux-differentiable and weakly lower semi-continuous and bounded from below on F. If Φ satisfies the (CPS)_{inf Φ,F} condition or the (WCPS)_{inf Φ,F} condition, and suppose that

 $\Phi(u_n) \to +\infty, \quad u_n \rightharpoonup u \in \partial \Lambda,$

then Φ attains its infimum on *F*.

The next lemma is a variant on the classical Tonelli's theorem, whose proof is easy, so we omit its proof.

Lemma 2.12 Let X be a Banach space, $F \subset X$ be a weakly closed subset. Suppose that $\phi(u)$ is defined on an open subset $\Lambda \subset X$ and is Gateaux-differentiable on Λ and weakly lower semi-continuous and bounded from below on $\Lambda \cap F$, if ϕ is coercive, that is, $\phi(x) \to +\infty$ as $||x|| \to +\infty$, and suppose that

$$\phi(u_n) \to +\infty, \quad u_n \rightharpoonup u \in \partial \Lambda,$$

then ϕ attains its infimum on $\Lambda \cap F$.

3 The proof of Theorem 1.7

By the symmetrical condition (A5)', it is easy to prove that the critical point of the functional f on Λ_0 is also the critical point of the functional f on Λ .

Let

$$\partial \Lambda_0 = \left\{ u \in H^1 = W^{1,2}(R/Z, R^n), u(t+1/2) = -u(t), \exists t_0, u(t_0) = 0 \right\}.$$

Lemma 3.1 Assume (A4)' holds, then for any weakly convergent sequence $u_n \rightharpoonup u \in \partial \Lambda_0$, we have

$$f(u_n) \to +\infty.$$

Proof Similar to the proof of Zhang [19].

Lemma 3.2 $F \cap \Lambda$ is a weakly closed subset in H^1 .

Proof Let $\{u_n\} \subset F \cap \Lambda$ be a weakly convergent sequence, we use the embedding theorem to find which uniformly converges to $u \in H^1$.

Now we claim $u \in \Lambda$, and then it is obvious that $u \in F$. In fact, if $u \in \partial \Lambda$, by $V(q) \to -\infty$, $q \to 0$ and the condition (A4)' we have

$$-V(u) \ge C_1 |u|^{-\beta}, \quad 0 < |u| < r' < r.$$

So V(u) satisfies Gordon's strong force condition, and by his lemma, we have

$$\int_0^1 -V(u_n)\,dt\to +\infty,\quad \forall u_n\rightharpoonup u\in\partial\Lambda.$$

The condition (A4)' implies

$$V(u_n) + \frac{1}{2} \langle V'(u_n), u_n \rangle \geq \left(1 - \frac{\beta}{2}\right) V(u_n).$$

Hence

$$h = \int_0^1 \left[V(u_n) + \frac{1}{2} \langle V'(u_n), u_n \rangle \right] dt \to +\infty.$$

This is a contradiction.

Lemma 3.3 f(u) is weakly lower semi-continuous on $F \cap \Lambda_0$

Proof For any $\{u_n\} \subset F : u_n \rightharpoonup u$, then by Sobolev's embedding theorem and functional analysis, we have uniform convergence:

$$|u_n(t)-u(t)|_{\infty}\to 0.$$

(i) If
$$u \in \Lambda_0$$
, then by $V \in C^1(\mathbb{R}^n \setminus \{0\}, \mathbb{R})$, we have

$$\left|V(u_n(t))-V(u(t))\right|_{\infty}\to 0.$$

It's well known that the norm is weakly lower semi-continuous, we have

 $\liminf \|u_n\| \ge \|u\|.$

Hence

$$\liminf f(u_n) = \liminf \left(\frac{1}{2} \int_0^1 |\dot{u}_n|^2 dt\right) \int_0^1 (h - V(u_n)) dt,$$
$$\geq \frac{1}{2} \int_0^1 |\dot{u}|^2 dt \int_0^1 (h - V(u)) dt = f(u).$$

(ii) If $u \in \partial \Lambda_0$, then by our assumption on *V* which satisfies Gordon's strong force condition, we have

$$\int_0^1 -V(u_n)\,dt\to +\infty,\quad \forall u_n\rightharpoonup u\in\partial\Lambda_0.$$

(1) If $u \equiv 0$, then

 $|u_n|_{\infty} \to 0$, $n \to +\infty$.

Then similar to the proof in [19], we have

$$f(u_n) \ge 6|u_n|_{\infty}^{2-\beta} \to +\infty, \quad n \to +\infty.$$

So in this case we have

$$\liminf f(u_n) = +\infty \ge f(u).$$

(2) If $u \neq 0$, then by the weakly lower semi-continuity for norm, we have

$$\liminf ||u_n|| \ge ||u|| > 0.$$

So by Gordon's lemma, we have

$$\liminf f(u_n) = \liminf \left(\frac{1}{2} \int_0^1 |\dot{u}_n|^2 dt\right) \int_0^1 (h - V(u_n)) dt = +\infty$$
$$\geq \frac{1}{2} \int_0^1 |\dot{u}|^2 dt \int_0^1 (h - V(u)) dt = f(u).$$

Lemma 3.4 The functional f(u) has a positive lower bound on F.

Proof By the definitions of f(u) and F and the assumption (A2), we have

$$f(u) = \frac{1}{4} \int_0^1 |\dot{u}|^2 dt \int_0^1 (V'(u)u) dt \ge 0, \quad \forall u \in F.$$

By the definitions of the functional f(u) and its domain Λ_0 , and the conditions on the energy h > 0 and the potential V(u) < 0, it is easy to prove the following lemma.

Lemma 3.5 The functional f(u) is coercive.

Furthermore, we claim that

$$c = \inf_{F \cap \Lambda_0} f(u) > 0,$$

since otherwise, $u_0(t) = \text{const}$ attains the infimum 0, then by the symmetry of Λ_0 , we have $u_0(t) \equiv o$, which contradicts the definition of Λ_0 . Now by Lemmas 3.1-3.4 and Lemmas 2.11 and 2.12, we know f(u) attains the infimum on *F*, furthermore we know that the minimizer is nonconstant.

4 The proof of Theorem 1.8

In order to prove the Cerami-Palais-Smale type condition and get a nonconstant periodic solution in non-symmetrical case, we need to add a topological condition, we know that there are winding numbers (degrees) in the planar case, so we define

$$\Lambda_1 = \big\{ u \in \Lambda, \deg(u) \neq 0 \big\}.$$

Lemma 4.1 If $u_n \rightharpoonup u \in \partial \Lambda_1$, then $f(u_n) \rightarrow +\infty$.

Proof By V satisfying Gordon's strong force condition, we have

$$\int_0^1 -V(u_n)\,dt\to+\infty,\quad\forall u_n\rightharpoonup u\in\partial\Lambda_1.$$

(1) If $u \equiv 0$, then by Sobolev's embedding theorem, we have

$$|u_n|_{\infty} \to 0, \quad n \to +\infty.$$

Then by $deg(u_n) \neq 0$, we have c > 0 such that

$$c|u_n|_{\infty} \leq \|\dot{u}_n\|_{L^2}$$

and $\|\dot{u}_n\|_{L^2}$ is an equivalent norm of $W^{1,2}$ and

$$f(u_n) \ge c |u_n|_{\infty}^{2-\beta} \to +\infty, \quad n \to +\infty.$$

So in this case, we have

$$\liminf f(u_n) = +\infty \ge f(u).$$

(2) If $u \neq 0$, then by the weakly lower semi-continuity for the norm, we have

 $\liminf ||u_n|| \ge ||u|| > 0.$

So by Gordon's lemma, we have

$$\liminf f(u_n) = \liminf \left(\frac{1}{2} \int_0^1 |\dot{u}_n|^2 dt\right) \int_0^1 (h - V(u_n)) dt = +\infty$$
$$= \frac{1}{2} \int_0^1 |\dot{u}|^2 dt \int_0^1 (h - V(u)) dt = f(u).$$

Lemma 4.2 Under the assumptions of Theorem 1.8,

$$f(u) = \frac{1}{2} \int_0^1 |\dot{u}|^2 dt \int_0^1 (h - V(u)) dt$$

satisfies the (CPS)⁺ condition on Λ_1 , that is, if $\{u_n\} \subset \Lambda_1$ satisfies

$$f(u_n) \to c > 0, \quad (1 + ||u_n||) f'(u_n) \to O,$$
 (4.1)

then $\{u_n\}$ has a strongly convergent subsequence in Λ_1 .

Proof Since $f'(u_n)$ makes sense, we know

$$\{u_n\}\subset \Lambda_1.$$

We claim $\int_0^1 |\dot{u}_n|^2 dt$ is bounded. In fact, by $f(u_n) \to c$, we have

$$-\frac{1}{2}\|\dot{u}_n\|_{L^2}^2 \cdot \int_0^1 V(u_n) \, dt \to c - \frac{h}{2}\|\dot{u}_n\|_{L^2}^2. \tag{4.2}$$

By (A3)' we have

$$\langle f'(u_n), u_n \rangle = \|\dot{u}_n\|_{L^2}^2 \cdot \int_0^1 \left(h - V(u_n) - \frac{1}{2} \langle V'(u_n), u_n \rangle \right) dt$$

$$\geq \|\dot{u}_n\|_{L^2}^2 \int_0^1 \left[h - \frac{\mu_2}{2} - \left(1 - \frac{\alpha}{2} \right) V(u_n) \right] dt.$$
 (4.3)

By (4.2) and (4.3) we have

$$\langle f'(u_n), u_n \rangle \ge \left(h - \frac{\mu_2}{2}\right) \|\dot{u}_n\|_{L^2}^2 + \left(1 - \frac{\alpha}{2}\right) \left(2c - h\|\dot{u}_n\|_{L^2}^2\right)$$

$$= \left(\frac{\alpha}{2}h - \frac{\mu_2}{2}\right) \|\dot{u}_n\|_{L^2}^2 + C_1,$$
(4.4)

where $C_1 = 2(1 - \frac{\alpha}{2})c$, $\alpha > 2$, $h > \frac{\mu_2}{\alpha}$. So $\|\dot{u}_n\|_2 \le C_2$.

Then we claim $|u_n(0)|$ is bounded.

We notice that

$$f'(u_n) \cdot (u_n - u_n(0))$$

= $\int_0^1 |\dot{u}_n|^2 dt \int_0^1 (h - V(u_n)) dt$

$$-\frac{1}{2}\int_{0}^{1}|\dot{u}_{n}|^{2} dt \int_{0}^{1} \langle V'(u_{n}), u_{n} - u_{n}(0) \rangle dt$$

= $2f(u_{n}) - \frac{1}{2}\int_{0}^{1}|\dot{u}_{n}|^{2}\int_{0}^{1} \langle V'(u_{n}), u_{n} - u_{n}(0) \rangle dt.$ (4.5)

If $|u_n(0)|$ is unbounded, then there is a subsequence, still denoted by u_n s.t. $|u_n(0)| \to +\infty$. Since

$$\|\dot{u}_n\| \leq M_1,$$

we have

$$\min_{0\le t\le 1} \left| u_n(t) \right| \ge \left| u_n(0) \right| - \| \dot{u}_n \|_2 \to +\infty, \quad \text{as } n \to +\infty.$$

$$\tag{4.6}$$

By Friedrics-Poincaré's inequality and the condition (P1), we have

$$\int_{0}^{1} \left| \dot{u}_{n}(t) \right|^{2} dt \ge \pi^{2} \int_{0}^{1} \left| u_{n}(t) - u_{n}(0) \right|^{2} dt,$$
(4.7)

$$\int_0^1 V'(u_n) \big(u_n - u_n(0) \big) \, dt \to 0, \tag{4.8}$$

$$f'(u_n) \cdot (u_n - u_n(0)) \to 0. \tag{4.9}$$

So $f(u_n) \to 0$, which contradicts $f(u_n) \to c > 0$, hence $u_n(0)$ is bounded, and $||u_n|| = ||\dot{u}_n||_{L^2} + |u_n(0)|$ is bounded. Furthermore, similar to the proof of Ambrosetti and Coti Zelati [15], u_n strongly converges to $u \in \Lambda$.

It is easy to prove the following.

Lemma 4.3 Under the assumption $(B1)', f(u) \ge 0$ on Λ , that is, f has a lower bound.

Lemma 4.4 Under the assumptions of Theorem 1.8, f(u) is weakly lower semi-continuous on the closure $\overline{\Lambda}$ of Λ .

Now we can prove our Theorem 1.8, in fact, by Lemma 4.1, we know that the infimum of f on Λ_1 is equal to the infimum of f on the closure of Λ_1 . Furthermore, we can prove the infimum of f on Λ_1 is greater than zero, otherwise if it is zero, the corresponding minimizer must be constant, then the winding number is zero, which is a contradiction. Now by the above lemmas, especially Lemma 2.11, we know that f attains the positive infimum on Λ_1 and the corresponding minimizer must be nonconstant.

Competing interests

The authors declare that they have no competing interests.

Authors' contributions

The research and writing of this manuscript was a collaborative effort made by all the authors. All authors read and approved the final manuscript.

Author details

¹School of Economic and Mathematics, Southwestern University of Finance and Economics, Chengdu, Sichuan 611130, P.R. China. ²Department of Mathematics, Sichuan University, Chengdu, Sichuan 610068, P.R. China.

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