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Research Article

Bargmann-Type Inequality for Half-Linear Differential Operators

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We consider the perturbed half-linear Euler differential equation $(\Phi(x'))' + [\gamma/t^p + c(t)]\Phi(x) = 0$, $\Phi(x) := |x|^{p-2}x$, p > 1, with the subcritical coefficient $\gamma < \gamma_p := ((p-1)/p)^p$. We establish a Bargmann-type necessary condition for the existence of a nontrivial solution of this equation with at least (n+1) zero points in $(0,\infty)$.

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1. Introduction

The classical Bargmann inequality [1] originates from the nonrelativistic quantum mechanics and gives an upper bound for the number of bound states produced by a radially symmetric potential in the two-body system. In the subsequent papers, various proofs and reformulations of this inequality have been presented, we refer to [2, Chapter XIII], and to [3–5] for some details.

In the language of singular differential operators, Bargmann's inequality concerns the one-dimensional Schrödinger operator

$$\tau(y) := y'' + \left[\frac{\gamma}{t^2} + c(t)\right] y, \quad \gamma < \frac{1}{4}, \ t \in (0, \infty). \tag{1.1}$$

It states that if the Friedrichs realization of τ has at least n negative eigenvalues below the essential spectrum (what is equivalent to the existence of a nontrivial solution of

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the equation $\tau(y) = 0$ having at least (n + 1) zeros in $(0, \infty)$, then

$$\int_0^\infty t c_+(t) dt > n \sqrt{1 - 4\gamma},\tag{1.2}$$

where $c_{+}(t) = \max\{c(t), 0\}.$

This inequality can be seen as follows. The Euler differential equation

$$x'' + \frac{\gamma}{t^2}x = 0\tag{1.3}$$

with the subcritical coefficient $\gamma < 1/4$ is disconjugate in $(0, \infty)$, that is, any nontrivial solution of (1.3) has at most one zero in this interval. Hence, if the equation $\tau(y) = 0$, with τ given by (1.1), has a solution with at least (n+1) positive zeros, the perturbation function c must be "sufficiently positive" in view of the Sturmian comparison theorem. Inequality (1.2) specifies exactly what "sufficient positiveness" means.

In this paper, we treat a similar problem in the scope of the theory of half-linear differential equations:

$$(r(t)\Phi(x'))' + c(t)\Phi(x) = 0, \qquad \Phi(x) := |x|^{p-2}x, \quad p > 1.$$
 (1.4)

In physical sciences, there are known phenomena which can be described by differential equations with the so-called *p*-Laplacian $\Delta_p u := \text{div } (\|\nabla u\|^{p-2} \nabla u)$, see, for example, [6]. If the potential in such an equation is radially symmetric, this equation can be reduced to a half-linear equation of the form (1.4).

There are many results of the linear oscillation theory, which concern the Sturm-Liouville differential equation:

$$(r(t)x')' + c(t)x = 0,$$
 (1.5)

which has been extended to (1.4). In particular, the linear Sturmian theory holds almost verbatim for (1.4), see, for example, [7, 8]. We will recall elements of the half-linear oscillation theory in the next section. Our main result concerns the perturbed half-linear Euler differential equation

$$\left(\Phi(x')\right) + \left[\frac{\gamma}{t^p} + c(t)\right]\Phi(x) = 0, \quad t \in (0, \infty), \tag{1.6}$$

where c is a continuous function, and shows that if γ is the so-called subcritical coefficient, that is, $\gamma < \gamma_p := (p/(p-1))^p$, and there exists a solution of (1.6) with at least (n+1) zeros in $(0,\infty)$, then the integral $\int_0^\infty t^{p-1}c_+(t)dt$ satisfies an inequality which reduces to (1.2) in the linear case p=2.

2. Preliminaries

In this short section, we present some elements of the half-linear oscillation theory which we need in the proof of our main result. As we have mentioned in the previous section, the linear

and half-linear oscillation theories are in many aspects very similar, so (1.4) can be classified as oscillatory or nonoscillatory as in the linear case.

If x is a solution of (1.4) such that $x(t) \neq 0$ is some interval I, then $w := r\Phi(x'/x)$ is a solution of the Riccati-type differential equation

$$w' + c(t) + (p-1)r^{1-q}|w|^q = 0, \quad q := \frac{p}{p-1}.$$
 (2.1)

If (1.4) is nonoscillatory, that is, (2.1) possesses a solution which exists on some interval $[T,\infty)$, among all such solutions of (2.1), there exists the *minimal* one \tilde{w} , minimal in the sense that any other solution w of (2.1) which exists on some interval $[t_w,\infty)$ satisfies $w(t) > \tilde{w}(t)$ in this interval, see [9,10] for details.

In our treatment, the so-called half-linear Euler differential equation

$$\left(\Phi(x')\right)' + \frac{\gamma}{t^p}\Phi(x) = 0 \tag{2.2}$$

appears. If we look for a solution of this equation in the form $x(t) = t^{\lambda}$, then λ is a root of the algebraic equation

$$|\lambda|^p - \Phi(\lambda) + \frac{\gamma}{p-1} = 0. \tag{2.3}$$

By a simple calculation (see, e.g., [8, Section 1.3]), one finds that (2.3) has a real root if and only if γ is less than or equal to the so-called critical constant $\gamma_p := ((p-1)/p)^p$, and hence (2.2) is nonoscillatory if and only if $\gamma \leq \gamma_p$. In this case, the associated Riccati equation is of the form

$$w' + \frac{\gamma}{t^p} + (p-1)|w|^q = 0, \tag{2.4}$$

and its minimal solution is $\widetilde{w}(t) = \Phi(\lambda_1)t^{1-p}$, where λ_1 is the smaller of (the two real) roots of (2.3). If $v(t) = t^{p-1}w$, then v is a solution of the equation

$$v' = \frac{p-1}{t} - \frac{p-1}{t} |v|^q - \frac{\gamma}{t},\tag{2.5}$$

and $\tilde{v}(t) \equiv \Phi(\lambda_1)$ is the minimal solution of this equation. A detailed study of half-linear Euler equation and of its perturbations can be found in [11].

3. Bargmann's Type Inequality

In this section, we present our main results, the half-linear version of Bargmann's inequality. We are motivated by the work in [4] where a short proof of this inequality based on the Riccati technique is presented. Here we show that this method, properly modified, can also be applied to (1.6).

Theorem 3.1. Suppose that (1.6) with $\gamma < \gamma_p = ((p-1)/p)^p$ has a nontrivial solution with at least (n+1) zeros in $(0,\infty)$. Then

$$\int_0^\infty t^{p-1}c_+(t)dt > nk(\gamma, q),\tag{3.1}$$

where $k(\gamma, q)$ is the absolute value of the difference of the real roots of

$$F_{\gamma}(\lambda) := |\lambda|^{q} - \lambda + (q - 1)\gamma = 0 \tag{3.2}$$

and q = p/(p-1) is the conjugate number to p. Moreover, the constant $k(\gamma, q)$ is strict in the sense that for every $\varepsilon > 0$, there exists a continuous function c such that (1.6) possesses a solution with (n+1) zeros in $(0,\infty)$ and

$$\int_{0}^{\infty} t^{p-1} c_{+}(t) dt \le nk(\gamma, q) + \varepsilon. \tag{3.3}$$

Proof. Let x be a solution of (1.6) with (n + 1) zeros in $(0, \infty)$, denote these zeros by $t_0 < t_1 < \cdots < t_n$, and let $v(t) = t^{p-1}\Phi(x'/x)$. Then by a direct computation we see that v is a solution of the Riccati-type differential equation

$$v' = \frac{p-1}{t}v - \frac{\gamma}{t} - (p-1)|v|^{q} - t^{p-1}c(t)$$

$$= -(p-1)F_{\gamma}(v) - t^{p-1}c(t), \ t \in (t_{i}, t_{i+1}), \ i = 0, \dots, n-1,$$
(3.4)

$$v(t_{i-}) = -\infty, \qquad v(t_{i+}) = \infty. \tag{3.5}$$

Let $\lambda_1 < \lambda_2$ be the roots of (3.2). Such pair of roots exists and it is unique since the function $F_{\gamma}(\lambda)$ is convex, $F_{\gamma}(\pm \infty) = \infty$, $F'_{\gamma}(1/\Phi(q)) = 0$, and $F_{\gamma}(1/\Phi(q)) = (\gamma - \gamma_p)/(p-1) < 0$. According to (3.5), there exist ξ_i , $\eta_i \in (t_i, t_{i+1})$ such that $v(\xi_i) = \lambda_2$, $v(\eta_i) = \lambda_1$, and $\lambda_1 < v(t) < \lambda_2$ for $t \in (\xi_i, \eta_i)$, which means that $F_{\gamma}(v(t)) < 0$ for $t \in (\xi_i, \eta_i)$. Then, we have

$$\int_{0}^{\infty} t^{p-1} c_{+}(t) dt \geq \sum_{i=0}^{n} \int_{\xi_{i}}^{\eta_{i}} t^{p-1} c_{+}(t) dt \geq \sum_{i=0}^{n} \int_{\xi_{i}}^{\eta_{i}} t^{p-1} c(t) dt$$

$$= \sum_{i=1}^{n} \int_{\xi_{i}}^{\eta_{i}} \left[-v'(t) - (p-1) F_{\gamma}(v(t)) \right] dt > \sum_{i=1}^{n} v(t) \Big|_{\eta_{i}}^{\xi_{i}}$$

$$= \sum_{i=1}^{n} \left[v(\xi_{i}) - v(\eta_{i}) \right] = n(\lambda_{2} - \lambda_{1}) = nk(\gamma, q).$$
(3.6)

Now we prove that the constant $k(\gamma, q)$ is exact. Let $\varepsilon > 0$ be arbitrary and α_i, β_i, T_i be sequences of positive real numbers constructed in the following way. Let $t_0 \in (0, \infty)$ be arbitrary and consider the differential equation

$$\left(\Phi(x')\right)' + \frac{\gamma}{t^p}\Phi(x) = 0. \tag{3.7}$$

Denote by x_0 its nontrivial solution satisfying $x_0(t_0) = 0$, $x_0'(t_0) = 1$ (such solution exists and it is unique, see, e.g., [8, Section 1.1]) and let $v_0 := t^{p-1}\Phi(x_0'/x_0)$. Since $\lim_{t\to\infty}v_0(t) = v_2$, see [8, page 39], there exists $T_1 > t_0$ such that $v_0(T_1)$.

Now, let

$$\alpha_1 := \frac{\gamma_p - \gamma}{T_1}, \qquad \beta_1 := \frac{\varepsilon T_1}{4n(\gamma_p - \gamma)}, \tag{3.8}$$

and define for $t \in [T_1, T_1 + \beta_1]$ the function

$$\widehat{c}_1(t) := \frac{1}{\beta_1 t^{p-1}} \left[k(\gamma, q) + \frac{\varepsilon}{4n} + \alpha_1 \right]. \tag{3.9}$$

Consider the solution v of the equation

$$v' = -(p-1)\frac{|v|^q}{t} + (p-1)\frac{v}{t} - \frac{\gamma}{t} - t^{p-1}\hat{c}_1(t), \quad t \in [T_1, T_1 + \beta_1], \tag{3.10}$$

given by the initial conditions $v(T_1) = v_0(T_1)$. Then for $t \in [T_1, T_1 + \beta_1]$

$$v' = -\frac{p-1}{t} \left[|v|^q - v + \frac{\gamma_p}{p-1} \right] + \frac{\gamma_p - \gamma}{t} - t^{p-1} \widehat{c}_1(t)$$

$$\leq \frac{\gamma_p - \gamma}{t} - \frac{1}{\beta_i} \left(k(\gamma, q) + \frac{\varepsilon}{4n} \right) - \frac{\gamma_p - \gamma}{T_1}$$

$$\leq -\frac{1}{\beta_i} \left(k(\gamma, q) + \frac{\varepsilon}{4n} \right). \tag{3.11}$$

Hence,

$$v(T_1 + \beta_1) = v(T_1) + \int_{T_1}^{T_1 + \beta_1} v'(t)dt < v_2 + \frac{\varepsilon}{4n} - \left(k(\gamma, q) + \frac{\varepsilon}{4n}\right)$$

$$= v_2 - (v_2 - v_1) = v_1.$$
(3.12)

Now consider again (3.7) and the associated Riccati-type differential equation

$$v' = -\frac{\gamma}{t^p} + (p-1)v - (p-1)|v|^q \tag{3.13}$$

(which is related to (3.7) by the substitution $v=t^{p-1}\Phi(x'/x)$). This equation has a constant solution $v=v_1$ and this solution is the minimal one (see the end of Section 2). This means that any solution of (3.13) which starts with the initial condition $v(T_1+\beta_1)< v_1$ blows down to $-\infty$ at a finite time $t_1>T_1+\beta_1$, which is a zero point of the associated solution x of (3.7). Now, let

$$\widetilde{c}_{1}(t) = \begin{cases}
0, & t \in [t_{0}, T_{1}], \\
\widehat{c}_{1}(t), & t \in [T_{1}, T_{1} + \beta_{1}], \\
0, & t \in [T_{1} + \beta_{1}, t_{1}].
\end{cases}$$
(3.14)

In summary, we have constructed a solution of the equation

$$\left(\Phi(x')\right)' + \left[\frac{\gamma}{t^p} + \tilde{c}_1(t)\right]\Phi(x) = 0 \tag{3.15}$$

for which $x(t_0) = 0 = x(t_1)$ and

$$\int_{t_0}^{t_1} t^{p-1} \tilde{c}_1(t) dt = \int_{T_1}^{T_1 + \beta_1} t^{p-1} \hat{c}_1(t) dt$$

$$= k(\gamma, q) + \frac{\varepsilon}{4n} + \alpha_1 \beta_1$$

$$= k(\gamma, q) + \frac{\varepsilon}{4n} + \frac{\varepsilon}{4n}$$

$$= k(\gamma, q) + \frac{\varepsilon}{4n} + \frac{\varepsilon}{4n}$$

$$= k(\gamma, q) + \frac{\varepsilon}{2n}.$$
(3.16)

The construction of T_i , β_i , α_i , $\widehat{c}_i(t)$ and $\widetilde{c}_i(t)$, $i=2,\ldots,n$, is now analogical. As a result we obtain the function $\widetilde{c}:(0,\infty)\to[0,\infty)$ defined as $\widetilde{c}(t)=0$ for $t\in(0,t_0]$ and $t\in[t_n,\infty)$, and $\widetilde{c}(t)=\widetilde{c}_i(t)$ for $t\in[t_{i-1},t_i]$, for which

$$\int_{0}^{\infty} t^{p-1} \tilde{c}(t) dt = nk(\gamma, q) + \frac{\varepsilon}{2}, \tag{3.17}$$

and the equation

$$\left(\Phi(x')\right)' + \left[\frac{\gamma}{t^p} + \tilde{c}(t)\right]\Phi(x) = 0 \tag{3.18}$$

has a solution with zeros at $t = t_i$, i = 0, ..., n.

Finally, we change the discontinuous function $\tilde{c}(t)$ to a continuous one $c(t) \geq \tilde{c}(t)$ such that $\int_{t_0}^{t_n} t^{p-1} [c(t) - \tilde{c}(t)] dt < \varepsilon/2$. Such a modification is an easy technical construction which can be described explicitly, but for us is only important its existence. According to

the Sturmian comparison theorem, the equation $(\Phi(x'))' + [\gamma/t^p + c(t)]\Phi(x) = 0$ possesses a nontrivial solution with at least (n+1) zeros and

$$\int_{0}^{\infty} t^{p-1}c(t)dt \le nk(\gamma, q) + \varepsilon, \tag{3.19}$$

which we needed to prove.

Remark 3.2. If p = 2, then $F_{\gamma}(\lambda) = \lambda^2 - \lambda + \gamma$ and the roots of (3.2) are

$$\lambda_{1,2} = \frac{1}{2} \left(1 \pm \sqrt{1 - 4\gamma} \right). \tag{3.20}$$

Hence, $k(\gamma, 2) = |\lambda_1 - \lambda_2| = \sqrt{1 - 4\gamma}$ and (3.1) reduces to (1.2).

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References

- [1] V. Bargmann, "On the number of bound states in a central field of force," *Proceedings of the National Academy of Sciences of the United States of America*, vol. 38, pp. 961–966, 1952.
- [2] M. Reed and B. Simon, Methods of Modern Mathematical Physics, Vol. IV. Analysis of Operators, Academic Press, Boston, Mass, USA, 1978.
- [3] Ph. Blanchard and J. Stubbe, "Bound states for Schrödinger Hamiltonians: phase space methods and applications," *Reviews in Mathematical Physics*, vol. 8, no. 4, pp. 503–547, 1996.
- [4] K. M. Schmidt, "A short proof for Bargmann-type inequalities," *The Royal Society of London*, vol. 458, no. 2027, pp. 2829–2832, 2002.
- [5] N. Setô, "Bargmann's inequalities in spaces of arbitrary dimension," *Publications of the Research Institute for Mathematical Sciences. Kyoto University*, vol. 9, pp. 429–461, 1974.
- [6] J. I. Díaz, Nonlinear Partial Differential Equations and Free Boundaries. Vol. I: Elliptic Equations, vol. 106 of Research Notes in Mathematics, Pitman, Boston, Mass, USA, 1985.
- [7] R. P. Agarwal, S. R. Grace, and D. O'Regan, Oscillation Theory for Second Order Linear, Half-Linear, Superlinear and Sublinear Dynamic Equations, Kluwer Academic Publishers, Dordrecht, The Netherlands, 2002.
- [8] O. Doślý and P. Řehák, Half-Linear Differential Equations, vol. 202 of North-Holland Mathematics Studies, Elsevier, Amsterdam, The Netherlands, 2005.
- [9] Á. Elbert and T. Kusano, "Principal solutions of non-oscillatory half-linear differential equations," *Advances in Mathematical Sciences and Applications*, vol. 18, pp. 745–759, 1998.
- [10] J. D. Mirzov, "Principal and nonprincipal solutions of a nonlinear system," Tbilisskii Gosudarstvennyi Universitet. Institut Prikladnoi Matematiki. Trudy, vol. 31, pp. 100–117, 1988.
- [11] Á. Elbert and A. Schneider, "Perturbations of the half-linear Euler differential equation," Results in Mathematics, vol. 37, no. 1-2, pp. 56–83, 2000.